

Springer Theses

Recognizing Outstanding Ph.D. Research

Aims and Scope

The series “Springer Theses” brings together a selection of the very best Ph.D. theses from around the world and across the physical sciences. Nominated and endorsed by two recognized specialists, each published volume has been selected for its scientific excellence and the high impact of its contents for the pertinent field of research. For greater accessibility to non-specialists, the published versions include an extended introduction, as well as a foreword by the student’s supervisor explaining the special relevance of the work for the field. As a whole, the series will provide a valuable resource both for newcomers to the research fields described, and for other scientists seeking detailed background information on special questions. Finally, it provides an accredited documentation of the valuable contributions made by today’s younger generation of scientists.

Theses are accepted into the series by invited nomination only and must fulfill all of the following criteria

- They must be written in good English.
- The topic should fall within the confines of Chemistry, Physics, Earth Sciences, Engineering and related interdisciplinary fields such as Materials, Nanoscience, Chemical Engineering, Complex Systems and Biophysics.
- The work reported in the thesis must represent a significant scientific advance.
- If the thesis includes previously published material, permission to reproduce this must be gained from the respective copyright holder.
- They must have been examined and passed during the 12 months prior to nomination.
- Each thesis should include a foreword by the supervisor outlining the significance of its content.
- The theses should have a clearly defined structure including an introduction accessible to scientists not expert in that particular field.

More information about this series at <http://www.springer.com/series/8790>

James Matthews

Disc Winds Matter

Modelling Accretion and Outflows on All
Scales

Doctoral Thesis accepted by
University of Southampton, Southampton, UK

Author

Dr. James Matthews
Department of Physics
University of Oxford
Oxford
UK

Supervisor

Prof. Christian Knigge
Department of Physics and Astronomy
University of Southampton
Southampton
UK

ISSN 2190-5053

Springer Theses

ISBN 978-3-319-59182-7

DOI 10.1007/978-3-319-59183-4

ISSN 2190-5061 (electronic)

ISBN 978-3-319-59183-4 (eBook)

Library of Congress Control Number: 2017941543

© Springer International Publishing AG 2017

This work is subject to copyright. All rights are reserved by the Publisher, whether the whole or part of the material is concerned, specifically the rights of translation, reprinting, reuse of illustrations, recitation, broadcasting, reproduction on microfilms or in any other physical way, and transmission or information storage and retrieval, electronic adaptation, computer software, or by similar or dissimilar methodology now known or hereafter developed.

The use of general descriptive names, registered names, trademarks, service marks, etc. in this publication does not imply, even in the absence of a specific statement, that such names are exempt from the relevant protective laws and regulations and therefore free for general use.

The publisher, the authors and the editors are safe to assume that the advice and information in this book are believed to be true and accurate at the date of publication. Neither the publisher nor the authors or the editors give a warranty, express or implied, with respect to the material contained herein or for any errors or omissions that may have been made. The publisher remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Printed on acid-free paper

This Springer imprint is published by Springer Nature

The registered company is Springer International Publishing AG

The registered company address is: Gewerbestrasse 11, 6330 Cham, Switzerland

“Here, on the edge of what we know, in contact with the ocean of the unknown, shines the mystery and the beauty of the world.

And it’s breathtaking.”

Seven Brief Lessons on Physics,
Carlo Rovelli

“Good enough for government work.”

Christian Knigge

*Dedicated to my family—and in memory of
Grandpa.*

Supervisor's Foreword

The following Ph.D. thesis is the work of James Matthews, who obtained his Ph.D. in Astrophysics from the University of Southampton. Most Ph.D. theses in the UK are likely to be read in full by exactly three people (in addition to the author, hopefully): the supervisor, the internal examiner and the external examiner. As one of these three, I felt this would be a great shame and a terrible waste in this case.

Dr. Matthews' thesis is an outstanding piece of work on all fronts: in terms of its astrophysical significance, in the depth and breadth of understanding it demonstrates, and in the quality and clarity of the writing it contains. Briefly, the thesis describes the application of a state-of-the-art Monte Carlo radiative transfer code to two distinct classes of accreting compact objects: accreting white dwarfs in compact stellar binary systems and accreting supermassive black holes in quasars and active galactic nuclei (AGN). These two types of systems are separated from each other by nine orders of magnitude in mass and size, but they nevertheless share much common physics. More specifically, the thesis is focused on the accretion disc winds that are seen in both types of system, and particularly the imprint these outflows leave on our observations (particularly in spectroscopic data sets). In both settings, the results described in the thesis provide important new insights, while also highlighting the physical connection between the two classes.

In the accreting white dwarf setting, the presence of disc winds has long been inferred from the presence of ultraviolet absorption lines. Here, the thesis shows that the same winds can actually also produce optical emission lines (as well as a recombination continuum), whose origin has been a long-standing problem. In the AGN context, the thesis constructs a simple model of disc winds in quasars that is capable of explaining both the absorption and emission signatures seen in these systems. This is a crucial result that supports a disc-wind based unification scenario for quasars and AGN. Finally, the thesis includes a combined theoretical and observational investigation into the equivalent width distribution of the emission lines in AGN. This reveals a major challenge to all unification scenarios: the equivalent widths of quasars that show broad absorption lines should be different from those that do not show them—but the observed equivalent width distributions are actually indistinguishable. The most likely resolution to this problem is that

the dominant emission from quasars and AGN does not arise in a standard, geometrically thin, optically thick accretion disc, as is commonly assumed.

What makes these results particularly remarkable is the depth and breadth of understanding that is necessary to achieve them. The thesis demonstrates this understanding very well, with excellent introductory material on topics ranging from accretion physics and Monte Carlo radiative transfer to compact binaries and quasars/AGN. For example, on the radiative transfer side, the overall Monte Carlo algorithm is outlined clearly and concisely, and key techniques (such as the micro-clumping approximation and the macro-atom approach) are described carefully and in detail. The author's excellent grasp of both the observational and theoretical literature on both accreting white dwarfs and AGN is also evident throughout.

In summary, this thesis is well worth reading not just for researchers interested in its primary results, but for anybody looking for a clear introduction to any of the topics it touches on. Beginning graduate students, in particular, should find it extremely useful, both for its content and as an example of what a great thesis can look like.

Southampton, UK
May 2017

Prof. Christian Knigge
Thesis Supervisor

Abstract

The release of gravitational potential energy as mass falls towards a potential well is the most efficient energetic phenomenon in the universe. This *accretion* process is thought to power the radiative central engines of active galaxies, known as active galactic nuclei (AGN) and is also responsible for the emission from compact binary systems such as X-ray binaries and cataclysmic variables (CVs). Perhaps remarkably, accretion invariably produces *outflow* of material; such outflows are ubiquitous in accreting systems across 10 orders of magnitude in mass, and there is good evidence that mass-loaded winds are launched from the accretion discs of quasars and CVs. Perhaps the most spectacular evidence for accretion disc winds is the blue-shifted, broad absorption lines (BALs) in UV resonance lines, seen in CVs and BAL quasars. As well as imprinting absorption features, disc winds may affect the line and continuum emission from accreting objects. They thus offer a natural way to *unify* much of the phenomenology of CVs and AGN. When one considers that disc winds are also a possible mechanism for AGN feedback, it becomes clear that understanding the physics and true spectral imprint of these winds is of wide-ranging astrophysical significance.

In this thesis, I use the state-of-the-art Monte Carlo radiative transfer (MCRT) code, Python, to conduct a series of simulations designed to test simple biconical disc wind models. I initially give a broad introduction to the field, starting with the basic physics of accretion before describing the observational appearance of accreting systems, from CVs to quasars. Throughout the introduction, I focus particularly on themes of universality and unification, which are central ideas to the work presented here. These themes are explored further in Chap. 2, in which I describe the theoretical and observational status of the current understanding of accretion disc winds. Chapter 3 contains a detailed description of MCRT techniques with particular focus on Leon Lucy's macro-atom scheme and its implementation in Python. This chapter is in some senses the least accessible as it provides a thorough documenting of the Monte Carlo estimators used in the simulations. Such rigour is important, however, and thus I also go on to describe a series of code validation exercises.

Having tested my methods thoroughly, I then explore whether the winds that are responsible for the UV BALs in high-state CVs could also have an effect on the optical spectrum. I find that the wind produces strong emission in the Balmer series, He II 4686 Å and a series of He I lines. The model shows the observed trends with inclination and in some cases produces sufficient recombination continuum emission to fill in the Balmer photoabsorption edge intrinsic to disc atmospheres. The results suggest that disc winds could have a significant impact on the optical spectra of high-state CVs.

The next step was to apply the techniques to quasar winds in a test of disc wind unification models. In previous efforts, the outflow tended to become ‘over-ionized’, and BAL features were only present if the X-ray luminosity was limited to around 10^{43} erg s⁻¹. The outflow also failed to produce significant line emission. Motivated by these problems, I introduce a simple treatment of clumping and find that it allows BAL features to form in the rest-frame UV at more realistic X-ray luminosities. The fiducial model shows good agreement with AGN X-ray properties and the wind produces strong line emission. Despite these successes, the model cannot reproduce all emission lines seen in quasar spectra with the correct equivalent width (EW) ratios, and I find the emission line EWs have a strong dependence on inclination.

Informed by the quasar wind modelling, I use the Sloan Digital Sky Survey to examine the emission line EW distributions of quasars in the context of geometric unification. I find that the observed distributions are not consistent with a model in which an equatorial BAL outflow rises from a foreshortened accretion disc. I discuss this finding in the context of other observational orientation indicators.

Finally, I summarize my findings and suggest avenues for future work. Overall, the work presented here suggests that *disc winds matter*. They not only act as a spectral ‘filter’ for the underlying accretion continuum, but may actually dominate the emergent spectrum from accreting objects. As a result, unveiling their driving mechanisms, mass-loss rates and ionization structure is an important goal for the astronomical community.

Acknowledgements

First and foremost, I would like to thank Christian Knigge, for being such an enthusiastic, helpful and stimulating supervisor throughout my Ph.D. Christian, I *always* left our meetings more positive than before—that speaks volumes—and I greatly enjoyed our conversations about science and Explosions in the Sky. I am also extremely grateful to Knox Long for all his assistance, writing the majority of the code and sharing some of his astronomy knowledge with me, and Stuart Sim for being immensely helpful throughout, especially when it comes to radiative transfer. I would also like to thank Nick Higginbottom for hours upon hours of assistance and friendship, and Sam Mangham for his help and input nearer the end of my Ph.D. To all of the above people; thank you for being part of a collaboration that does some great science, but also knew when to discuss a disastrous code bug with a knowing smile and ironic joke. I would also like to thank Daniel Proga, Omer Blaes, Ivan Hubeny, Shane Davis, Mike Brotherton, Mike DiPompeo, Frederic Marin, Daniel Capellupo, Dirk Grupe, Simo Scaringi, Adam Foster, Randall Smith, Chris Done, Anna Pala, Boris Gaensicke, Patrick Woudt and countless others for useful correspondence or stimulating scientific conversations. I am also very grateful to my masters supervisor Rosanne Di Stefano for inspiring me to continue doing research beyond my year in Boston, and to Lance Miller and Poshak Gandhi for being fair, friendly and helpful Ph.D. examiners.

Apart from those who I worked with, I am grateful to everyone who helped make Southampton a happy place to be over 3+ years. There are too many to name, but I will indulge a few. To Cat, thank you for being generally lovely, being so patient with me and having great taste in comedy. To Sam Connolly, you have been ever-present and someone who I can always rely on for good beer, music and knowing looks. Sadie Jones, you were there when I needed you most for a cwtch or wanted to hear you pronounce year or ear. To Rob, Arran, Juan, Georgios and Poppy, thanks for being true friends throughout the process. Big thanks to everyone else in the department for making it a super cool place to work, my office mates Mari, Judith, Stew and CFro, and my housemates Rich, Stu, Luke, Ollie, Sophie and Szymon. Thanks also go to all the staff for being so friendly and approachable,

particularly Poshak, Diego, Francesco, Mark and Seb who have all helped me at various points through the Ph.D., and to Phil Charles for his sage advice since the early days.

There are also many people, outside of Southampton, who I must thank. Most of all, Mum, Dad, Beth and Nick, and the rest of my family—you have always supported me throughout my studies and just been all round lovely people. Cheers to Nick, Jack and the Hannahs for being great mates and helping me avoid Lemony Snicketts. Thank you Josh, James and Alex, and our now manager John, for allowing me to play music that I love in Waking Aida. Eschaton and Full Heal were both released during my Ph.D., and in many ways they are each a thesis within themselves.

This work was supported by the Science and Technology Facilities Council via a 3.5 year Ph.D. studentship. This work made use of the Sloan Digital Sky Survey. Funding for the Sloan Digital Sky Survey has been provided by the Alfred P. Sloan Foundation, the U.S. Department of Energy Office of Science, and the Participating Institutions. I acknowledge the use of the IRIDIS High Performance Computing Facility, and associated support services at the University of Southampton, in the completion of this work. Figures were produced using the matplotlib plotting library. I would also like to thank Padma and Ramon from Springer for helping me through the publication process.

Declaration of Authorship

I, James Matthews, declare that this thesis titled, ‘Disc Winds Matter: Modelling Accretion and Outflow on All Scales’ and the work presented in it are my own.

The first two chapters of this thesis provide a general introduction to the field, and thus are based on the relevant literature. Where a figure is not produced by me, I have acknowledged this clearly with a credit to the relevant publication. Chapter 3 contains a description of the methods used. This is partly a description of the radiative transfer code Python, which was originally developed by Knox Long and Christian Knigge (Long and knigge 2002), but also includes substantial detail on the ‘macro-atom’ technique, which was proposed by Leon Lucy and incorporated into Python by Stuart Sim for a study on young stellar objects (Sim et al. 2005). Although I have put significant effort into testing, fixing and developing this scheme, I did not write the original code to deal with macro-atoms in PYTHON.

Chapters 4–6 were studies I led under the guidance of my supervisor. For these chapters I conducted all simulations and data analysis, produced all the figures and wrote the text. These chapters are adapted from the following three papers with permission from MNRAS:

- Chapter 4: Matthews J. H., Knigge C., Long K. S., Sim S. A., Higginbottom N., ‘The impact of accretion disc winds on the optical spectra of cataclysmic variables’, 2015, MNRAS, 450, 3331.
- Chapter 5: Matthews J. H., Knigge C., Long K. S., Sim S. A., Higginbottom N., Mangham S. W., ‘Testing quasar unification: radiative transfer in clumpy winds’, 2016, MNRAS, 458, 293.
- Chapter 6: Matthews J. H., Knigge C., Long K. S., ‘Quasar emission lines as probes of orientation: implications for disc wind geometries and unification’, MNRAS, DOI:10.1093/mnras/stx231.

The following additional publications are not included in this thesis, although some of the work presented here did contribute towards the respective results:

- Higginbottom N., Knigge C., Long K. S., Sim S. A., Matthews J. H., ‘A simple disc wind model for broad absorption line quasars’, 2013, MNRAS, 436, 2.3.

- Higginbottom, N., Proga D., Knigge C., Long K. S., Matthews J. H., Sim S. A., ‘Line-driven disc Winds in Active Galactic Nuclei: The Critical Importance of Ionization’ and Radiative Transfer, 2014, ApJ, 789, 1.
- Shankar F., Calderone G., Knigge C., Matthews J. H., et al., ‘The Optical–UV Emissivity of Quasars: Dependence on Black Hole Mass and Radio Loudness’, 2016, ApJ Letters, 818, 1.

References

- Long KS, Knigge C (2002) Modeling the Spectral Signatures of Accretion Disk Winds: A New Monte Carlo Approach. APJ 579:725–740. doi:[10.1086/342879](https://doi.org/10.1086/342879). [arXiv:astro-ph/0208011](https://arxiv.org/abs/astro-ph/0208011)
- Sim SA, Drew JE, Long KS (2005) Two-dimensional Monte Carlo simulations of HI line formation in massive young stellar object disc winds. MNRAS 363:615–627, [10.1111/j.1365-2966.2005.09472.x](https://doi.org/10.1111/j.1365-2966.2005.09472.x), [arXiv:astro-ph/0508103](https://arxiv.org/abs/astro-ph/0508103)

Contents

1	Introduction	1
1.1	The Physics of Accretion	2
1.1.1	Spherical Accretion and the Eddington Limit	4
1.1.2	Accretion Discs	5
1.1.3	Boundary Layers, Black Hole Spin and the ISCO	8
1.1.4	The Emergent Spectrum	9
1.2	Accreting Compact Binaries	10
1.2.1	Roche Lobe-Overflow	11
1.2.2	Cataclysmic Variables	13
1.2.3	Low Mass X-Ray Binaries	18
1.3	Quasars and Active Galactic Nuclei	18
1.3.1	AGN Unification and the Dusty Torus	19
1.3.2	X-Ray Properties of AGN	21
1.3.3	The Broad Line Region: Connection to Winds and Unification	24
1.4	The Current Understanding of the Disc Continuum	24
1.4.1	The Spectral Shape of CV Discs	25
1.4.2	The Big Blue Bump in AGN	26
1.5	The Universality of Accretion	27
1.5.1	The RMS-Flux Relation	27
1.5.2	Accretion States and Disc-Jet Coupling	27
1.5.3	A Global Picture	29
	References	29
2	Accretion Disc Winds	39
2.1	Observational Evidence	39
2.1.1	Cataclysmic Variables	39
2.1.2	X-Ray Binaries	42
2.1.3	AGN and Quasars	44
2.1.4	Stellar Winds	49
2.1.5	Outflow Physics	50

- 2.2 Driving Mechanisms. 51
 - 2.2.1 Thermal Winds 52
 - 2.2.2 Radiatively Driven Winds 52
 - 2.2.3 Line-Driven Winds 52
 - 2.2.4 Magnetic Winds 54
- 2.3 Accretion Disc Wind Models 55
 - 2.3.1 MCGV95: A Line-Driven Wind Model for AGN 55
 - 2.3.2 De Kool and Begelman: A Radiatively Driven,
Magnetically Confined Wind. 56
 - 2.3.3 Elvis 2000: A Structure for Quasars 57
 - 2.3.4 Proga et al.: Line-Driven Hydrodynamic Models
for AGN and CVs. 58
- 2.4 A Kinematic Prescription for a Biconical Wind 60
- 2.5 The Big Picture: AGN Feedback 62
 - 2.5.1 Observational Evidence for Feedback 63
 - 2.5.2 Alternative Explanations 65
- References. 65
- 3 Monte Carlo Radiative Transfer and Ionization 77**
 - 3.1 Fundamentals of Radiative Transfer 77
 - 3.1.1 Spectral Line Formation 79
 - 3.1.2 Local Thermodynamic Equilibrium 79
 - 3.1.3 The Two Level Atom 81
 - 3.1.4 The Sobolev Approximation 82
 - 3.1.5 Monte Carlo Approaches 84
 - 3.2 PYTHON: A Monte Carlo Ionization
and Radiative Transfer Code 84
 - 3.2.1 Basics 85
 - 3.2.2 Radiation Packets 86
 - 3.2.3 Radiative Transfer Procedure 87
 - 3.3 Macro-Atoms 90
 - 3.3.1 Transition Probabilities 91
 - 3.3.2 Rate Equations 92
 - 3.3.3 Macro-Atom Estimators 94
 - 3.3.4 *k*-Packets 97
 - 3.3.5 Putting It All Together 98
 - 3.3.6 Ionization Fractions and Level Populations 98
 - 3.3.7 Numerical Issues and Population Inversions 100
 - 3.4 A Hybrid Line Transfer Scheme: Including Simple-Atoms. 101
 - 3.4.1 Line Transfer 101
 - 3.4.2 Heating and Cooling Estimators 102
 - 3.5 Heating and Cooling Balance 105
 - 3.5.1 Convergence 105
 - 3.6 Spectral Cycles. 106
 - 3.6.1 Macro-Atom Emissivity Calculation 107

3.7	Atomic Data	108
3.7.1	Macro-Atom Level and Line Data	108
3.7.2	Photoionization Cross-Sections	109
3.8	Code Validation	110
3.8.1	Testing Against Cloudy.	110
3.8.2	Macro-Atom Testing Against TARDIS and Theory	114
3.8.3	Testing Line Transfer Modes	116
3.9	Code Maintenance and Version Control	117
3.9.1	Parallelisation	117
	References.	118
4	The Impact of Accretion Disc Winds on the Optical Spectra of Cataclysmic Variables.	121
4.1	Introduction	121
4.1.1	Sources and Sinks of Radiation.	122
4.2	A Benchmark Disc Wind Model	124
4.2.1	Physical Structure and Ionization State	124
4.2.2	Synthetic Spectra.	126
4.3	A Revised Model Optimized for Optical Wavelengths	130
4.3.1	Synthetic Spectra.	131
4.3.2	Continuum Shape and the Balmer Jump	134
4.3.3	Line Profile Shapes: Producing Single-peaked Emission	135
4.3.4	Sensitivity to Model Parameters	137
4.3.5	Comparison to RW Tri	137
4.3.6	A Note on Collision Strengths	139
4.4	Conclusions	140
	References.	141
5	Testing Quasar Unification: Radiative Transfer in Clumpy Winds	143
5.1	Introduction	143
5.2	A Clumpy Biconical Disk Wind Model for Quasars.	144
5.2.1	Photon Sources	145
5.2.2	A Simple Approximation for Clumping	145
5.2.3	The Simulation Grid	147
5.3	Results and Analysis from a Fiducial Model	149
5.3.1	Physical Conditions and Ionization State.	150
5.3.2	Synthetic Spectra: Comparison to Observations	152
5.3.3	X-Ray Properties.	157
5.3.4	LoBALs and Ionization Stratification	158

5.4	Discussion	159
5.4.1	Parameter Sensitivity	159
5.4.2	Inclination Trends: FWHM and EW	161
5.5	Summary and Conclusions	164
	References	165
6	Quasar Emission Lines as Probes of Orientation and Unification	169
6.1	Introduction	169
6.2	Data Sample	171
6.3	The Angular Distribution of Emission from an Accretion Disc	174
6.3.1	Standard Thin Disc Models	174
6.3.2	Including GR and Opacity Effects	175
6.4	Predicted EW Distributions Compared to Observations	175
6.4.1	Fitting the Quasar Distribution	175
6.4.2	Comparing Non-BAL and LoBAL Distributions: Sample A	177
6.4.3	Alternative Shapes for the Intrinsic EW Distribution	181
6.4.4	Caveats and Selection Effects	183
6.5	Discussion	183
6.5.1	Eigenvector 1	184
6.5.2	Polarisation	185
6.5.3	The Effect of Obscuration	188
6.5.4	Line Anisotropy	188
6.6	Conclusions	192
	References	193
7	Conclusions and Future Work	197
7.1	Suggestions for Future Work	200
7.1.1	CVs as Accretion and Outflow Laboratories	200
7.1.2	Improving the Treatment of Clumping	200
7.1.3	Improving Atomic Models	201
7.1.4	Using Radiative Transfer to Make Reverberation and Microlensing Predictions	202
7.1.5	Placing BAL Quasars on the Eigenvector 1 Parameter Space	203
7.2	Closing Remarks	203
	References	204
	Index	207

Abbreviations

ADAF	Advection dominated accretion flow
AGN	Active galactic nuclei/nucleus
BAL	Broad absorption line
BALQSO	Broad absorption line quasar
BBB	Big blue bump
BEL	Broad emission line
BH	Black hole
BI	Balnicity index
BL	Boundary layer
BLR	Broad line region
CDF	Cumulative distribution function
CGS	Centimetre-gram-second
CV	Cataclysmic variable
DN(e)	Dwarf nova(e)
EW	Equivalent width
HERG	High excitation radio galaxy
HMXB	High-mass X-ray binary
HSE	Hydrostatic equilibrium
IP	Intermediate polar
IR	Infra-red
ISCO	Innermost stable circular orbit
LHS	Left-hand side
LINER	Low ionization nuclear emission line region
LMXB	Low-mass X-ray binary
LTE	Local thermodynamic equilibrium
MCRT	Monte Carlo radiative transfer
MRI	Magneto-rotational instability
NAL	Narrow absorption line
NEL	Narrow emission line
NL	Nova-like variable

NLR	Narrow line region
NS	Neutron star
QSO	Quasi-stellar object/Quasar
RHS	Right-hand side
RIAF	Radiatively inefficient accretion flow
RL	Radio-loud
RLOF	Roche lobe overflow
SA	Sobolev approximation
SED	Spectral energy distribution
SFR	Star formation rate
SMBH	Supermassive black hole
sSFR	Specific star formation rate
SXXS	Soft X-ray excess
UFO	Ultra-fast outflow
UV	Ultraviolet
WA	Warm absorber
WD	White dwarf
WHIM	Warm, highly ionized medium
XRB	X-ray binary
YSO	Young stellar object

Physical Constants

Speed of light, c	$2.997\,924\,58 \times 10^{10} \text{ cm s}^{-1}$
Boltzmann constant, k	$1.380\,658 \times 10^{-16} \text{ erg K}^{-1}$
Gravitational constant, G	$6.672\,599 \times 10^{-8} \text{ cm}^3 \text{ g}^{-1} \text{ s}^{-2}$
Solar mass, M_{\odot}	$1.988\,550 \times 10^{33} \text{ g}$
Solar radius, L_{\odot}	$6.957\,000 \times 10^{10} \text{ cm}$
Thomson cross-section, σ_T	$6.652\,458 \times 10^{-25} \text{ cm}^2$
Planck constant, h	$6.626\,076 \times 10^{-27} \text{ cm}^2 \text{ g s}^{-1}$
Stefan–Boltzmann constant, σ	$5.670\,367 \times 10^{-5} \text{ erg cm}^{-2} \text{ K}^{-4} \text{ s}^{-1}$
Parsec, pc	$3.085\,678 \times 10^{18} \text{ cm}$
Proton mass, m_p	$1.672\,621 \times 10^{-24} \text{ g}$
Electron mass, m_e	$9.109\,390 \times 10^{-28} \text{ g}$
Electron volt, eV	$1.602\,177 \times 10^{-12} \text{ erg}$
Electron charge, q_e	$4.803\,207 \times 10^{-10} \text{ esu}$
π	3.141 593
e	2.718 282

All constants are given to 6 decimal places, except for the speed of light which is exact. Throughout this thesis I use the centimetre-gram-second (CGS) system of units unless otherwise stated.

List of Figures

Fig. 1.1 The temperature profile of an accretion disc for three different classes of compact object 7

Fig. 1.2 Angular velocity as a function of radius in an accretion disc around a rotating compact object with angular velocity Ω_* . Ω_K is the Keplerian angular velocity. This *graph* also helps explain why there is a turnover in the temperature–radius relation, as $D(R)$ is proportional to the square of the *derivative* of this quantity. The boundary condition responsible for this turnover is briefly discussed in the text 8

Fig. 1.3 The radius of the ISCO, R_{ISCO} , and the horizon, R_H , is as a function of the BH spin parameter, a_* . $a_* = 0$ corresponds to a Schwarzschild BH, and $a_* = 1$ and $a_* = -1$ to prograde and retrograde Kerr BHs, respectively. Note that this figure ignores the counteracting torque of photons swallowed by the BH, which actually limits a_* to a value of around 0.998 (Thorne 1974) 9

Fig. 1.4 Accretion disc SEDs for three different compact objects, corresponding roughly to a quasar, an XRB and a CV. The SEDs are computed via an area-weighted sum of blackbodies with effective temperatures governed by Eq. 1.19, and the $\nu^{1/3}$ shape in the middle of the spectra can be seen 10

Fig. 1.5 A comparison between a disc spectrum computed with stellar atmosphere models and standard multi-temperature blackbody disc. The major photoabsorption *edges* are marked. Flux is reprocessed into different wavelengths by bound-free opacities, and line blanketing also has a big effect on the spectrum. The Hydrogen and Helium lines also experience a degree of pressure broadening 11

Fig. 1.6	Artist's impression of a low-mass X-ray binary (<i>top</i>) and cataclysmic variable (<i>bottom</i>). The key components are marked, and the clear similarity in overall structure is apparent. <i>Credit Rob Hynes</i>	12
Fig. 1.7	Schematic showing the Roche lobes in a binary system, where $M_1 > M_2$. The first Lagrangian point is marked, as are the locations of the individual and system centres of mass	13
Fig. 1.8	<i>Data AAVSO</i> . A year in the life of SS Cyg, showing the characteristic repeated outbursts and periods of quiescence typical of a DN. SS Cyg has been undergoing this activity since it was first observed in 1896	14
Fig. 1.9	Spectra of SS Cyg during an outburst cycle, showing the evolution from minimum to maximum light. The rise is characterised by the appearance of an optically thick accretion disc spectrum. The flux scale is approximate. <i>Credit Hessman et al. 1984/Dhillon 1996</i>	15
Fig. 1.10	UV spectrum of RW Tri in and out of eclipse, showing strong lines in C IV 1550 Å and Ly α , among others <i>Credit Noebauer et al. 2010</i>	16
Fig. 1.11	Optical spectra of three nova-like variables: RW Sex (<i>top</i> ; Beuermann et al. 1992), IX Vel (<i>top middle</i> A.F. Pala and B.T. Gaensicke, private communication) and RW Tri in and out of eclipse (<i>bottom two panels</i> Groot et al. 2004). The data for RW Sex and RW Tri were digitized from the respective publications, and the IX Vel spectrum was obtained using the XSHOOTER spectrograph on the Very Large Telescope on 2014 October 10. These systems have approximate inclinations of 30°, 60° and 80° respectively. The trend of increasing Balmer line emission with inclination can be seen. In RWTri strong single-peaked emission in the Balmer lines is seen even in eclipse, indicating that the lines may be formed in a spatially extensive disc wind, and there is even a suggestion of a (potentially wind-formed) recombination continuum in the eclipsed spectrum. I have attempted to show each spectrum over a similar dynamic range	17
Fig. 1.12	Template spectra, from the AGN atlas, for four common types of AGN. Obtained from http://www.stsci.edu/hst/observatory/crds/cdb_s_agn.html	20
Fig. 1.13	A unified scheme for AGN. The broad line region is compact and can thus be obscured from view by a dusty torus, whereas the larger narrow line region is always visible. The radio-loud/radio-quiet dichotomy is explained by the presence or absence of a radio jet. <i>Credit Modified from Urry and Padovani (1995). Copyright PASP, reprinted by permission of the authors</i>	22

Fig. 1.14 Approximate average broadband SEDs for a few types of AGN. The series of characteristic bumps can be clearly seen. The Soft-X-ray excess is also visible (see Sect. 1.3.2.1). *Credit* Koratkar and Blaes (1999). *Copyright* PASP, *reprinted by permission of the authors* 23

Fig. 1.15 Occam’s quasar. How far can this general picture take us when trying to explain the behaviour of quasars and other accreting compact objects? 25

Fig. 1.16 Hardness-intensity diagrams for three types of accreting compact objects: the dwarf nova SS Cyg (*left*), in addition to neutron star (*middle*) and black hole (*right*) LMXBs. The dotted line shows the ‘jet line’ in LMXBs and the *arrows* show the expected temporal evolution of the system. For SS Cyg the *x*-axis shows the fraction of the hard X-ray power law component in the system compared to the luminosity of the combined luminosity of disc and power law. This quantity is analogous to the X-ray hardness in LMXBs. *Credit* Fig. 1.1. K rding et al. 2008, *Science*, 320, 1318. *Reprinted with permission from AAAS*. 28

Fig. 2.1 Diagram showing how an expanding envelope or wind presenting significant line opacity around a continuum source leads to the formation of P-Cygni profiles. The *black arrows* denote the outflow direction and the *blue arrows* typical scattering interactions 40

Fig. 2.2 UV spectrum of the DN TW Vir during outburst. The P-Cygni profiles can be seen clearly, demonstrating that a strong, fast outflow is present in the system. *Credit* Cordova and Mason (1982). *Copyright* AAS. *Reproduced with permission*. 41

Fig. 2.3 UV spectrum of Z Cam (*black*), compared to a synthetic spectrum from MCRT simulations (*grey*). *Credit* Long and Knigge (2002). *Copyright* AAS. *Reproduced with permission* 41

Fig. 2.4 A comparison between a line profile, normalised to have peak intensity of 1, produced from a Keplerian disk (*solid line*) and the same model with an additional disc wind (*dashed line*). The radial velocity component of the disc wind modifies the escape probabilities across the disc, causing a single-peaked line to form. *Credit* Murray and Chiang (1997) 42

Fig. 2.5 A cartoon illustrating the expected geometry of the disc and outflows in LMXBs in the soft and hard states. *Credit* Figure 3, Ponti et al. 2012, “Ubiquitous equatorial accretion disc winds in black hole soft states”, MNRAS letters, 422, 11 43

Fig. 2.6	Hardness-luminosity diagram for four dipping LMXBs, demonstrating that winds appear only in the soft state. The points are colour-coded by system, and are shown against the background <i>grey</i> points of all LMXBs studied by Ponti et al. (2012). None of the low inclination sources in their sample show Fe xxvi absorption detections. <i>Credit</i> Figure 2, Ponti et al. 2012, “Ubiquitous equatorial accretion disc winds in black hole soft states”, MNRAS letters, 422, 11	44
Fig. 2.7	<i>Top panel</i> A comparison between the SDSS non-BAL, HiBAL and LoBAL composite spectra as presented in Reichard et al. (2003). <i>Bottom two panels</i> Two individual examples of a HiBAL and LoBAL quasar spectrum, respectively. In all panels some of the more prominent lines are labeled and shaded, and the object name is given in the <i>bottom two plots</i>	45
Fig. 2.8	X-ray spectrum of PDS 456 fitted with a P-Cygni profile from a spherical outflow model. XMM-Newton data is shown in <i>black</i> with two combined NuStar observations in <i>blue</i> . <i>Credit</i> Figure 3, Nardini et al. (2015), Science, 347, 860. Reprinted with permission from AAAS.	49
Fig. 2.9	UV spectrum of one of the brightest massive O stars, the O4 supergiant ζ Puppis. The spectrum is obtained from IUE UV observations.	50
Fig. 2.10	Cartoon showing the geometry of the MCGV95 model. <i>Credit</i> Murray et al. 1995. Copyright AAS. Reproduced with permission	55
Fig. 2.11	A cartoon showing the components in the De Kool and Begelman model. <i>Credit</i> De Kool and Begelman 1995. Copyright AAS. Reproduced with permission.	56
Fig. 2.12	A schematic showing the main features of the Elvis model. A biconical wind rises from an accretion disc, and the observed spectrum is determined purely by the viewing angle of the observer. <i>Credit</i> Martin Elvis	57
Fig. 2.13	A snapshot of the PK04 model. <i>Colours</i> shows the density and arrows show the velocity of the flow. The radial lines separate three areas described by H14. <i>Credit</i> Higginbottom et al. (2014). Copyright AAS. Reproduced with permission	59
Fig. 2.14	A schematic showing the geometry and kinematics of the SV93 model	60
Fig. 2.15	The SV93 velocity law for various values of the acceleration exponent, α	61
Fig. 2.16	The $M_{BH} - \sigma_*$ correlation. <i>Credit</i> Gültekin et al. (2009). Copyright AAS. Reproduced with permission	63

Fig. 2.17 The $M_{BH} - M_{bulge}$ correlation. *Credit* McConnell and Ma (2013). Copyright AAS. Reproduced with permission. 64

Fig. 3.1 A schematic showing a ray obliquely incident on a surface of area dA . The labeled quantities are used in the definition of specific intensity 78

Fig. 3.2 A flowchart showing the basic operation of PYTHON 85

Fig. 3.3 The process of choosing a scattering location in a cell. *Credit* Mazzali and Lucy 1993. 88

Fig. 3.4 The decision tree traversed by an energy packet in macro-atom mode, depicting the interaction between radiation (r -packets), the thermal pool (k -packets), and ionization and ionization/excitation energy (macro-atoms). The probabilities at each decision point are marked, and are defined in the text. The *red, blue* and *green coloured arrows* represent radiant, kinetic and ionization/excitation energy respectively. The *symbols* are defined in the text, except C_{cont} and C_{ex} which refer to cooling contributions to radiative and excitation energy respectively 99

Fig. 3.5 An example of a modeled spectrum in PYTHON compared to the recorded MC spectrum, from an individual cell in an AGN model. 104

Fig. 3.6 The average temperature and fraction of converged cells in a typical CV model, shown as a function of the number of ionization cycles completed. 106

Fig. 3.7 A synthetic spectrum after 30 spectral cycles with 100,000 photons from simple CV wind model at a 60° viewing angle. Spectra produced with both the extract and live or die modes are shown. The effectiveness of the extract variance reduction technique can be clearly seen, and we can see that the spectral shape is unaltered 107

Fig. 3.8 A comparison of the soft X-ray regime of the H13 model, with two different datasets. *standard73* is the dataset with old, unextrapolated cross-sections and *standard77* instead includes extrapolated cross-sections as described in the text 110

Fig. 3.9 Relative ion fractions as a function of ionization parameter from the hybrid macro-atom scheme, with Hydrogen and Helium treated as a full macro-atom (*dotted lines*), compared to both CLOUDY (*solid lines*) and PYTHON in simple-atom only mode (H13, crosses). The *colours* correspond to the different ionic stages and are labeled for the CLOUDY case. 111

Fig. 3.10	As Fig. 3.9, but for Helium	111
Fig. 3.11	As Fig. 3.9, but for carbon	112
Fig. 3.12	As Fig. 3.9, but for nitrogen	112
Fig. 3.13	As Fig. 3.9, but for oxygen	113
Fig. 3.14	As Fig. 3.9, but for the first 11 ionization stages of iron.	113
Fig. 3.15	<i>Top Panel</i> ‘Case B’ Balmer decrements computed with PYTHON compared to analytic calculations by Seaton (1959). Both calculations are calculated at $T_e = 10,000\text{K}$. <i>Bottom Panel</i> a comparison of He I level populations (the most complex ion we currently treat as a macro-atom) between PYTHON and TARDIS models. The calculation is conducted in thin shell mode with physical parameters of $n_e = 5.96 \times 10^4 \text{ cm}^{-3}$, $T_e = 30,600\text{K}$, $T_R = 43,482\text{K}$ and $W = 9.65 \times 10^{-5}$	114
Fig. 3.16	Comparison between TARDIS and PYTHON synthetic spectra from a simple 1D supernova model. A bug in Doppler shifting of photons was discovered around PYTHON 76, meaning that the code now gives even better agreement than presented in Kerzendorf and Sim (2014).	115
Fig. 3.17	A comparison between weight reduction and line transfer mode. The model is the H13 fiducial BALQSO model—agreement is generally even better for CV models	116
Fig. 3.18	Commit history from Feb 3, 2013 to Feb 29, 2016, showing the regular code development that makes version control such a necessity to a collaborative code project. Produced using the Github API and plotting capability	117
Fig. 3.19	Total runtime per cycle for an AGN run as a function of the number of processors.	118
Fig. 4.1	Cartoon illustrating the geometry and kinematics of the benchmark CV wind model.	122
Fig. 4.2	The spectral energy distribution of the accretion disc and white dwarf used in the ionization cycles for the CV modelling. The important ionization edges for hydrogen and helium are marked	123
Fig. 4.3	The physical properties of the wind—note the logarithmic scale. Near the disc plane the wind is dense, with low poloidal velocities. As the wind accelerates it becomes less dense and more highly ionized. The dominant He ion is almost always He III, apart from in a small portion of the wind at the base, which is partially shielded from the inner disc	126

Fig. 4.4 UV (*left*) and optical (*right*) synthetic spectra for model A, the benchmark model, computed at sightlines of 10, 27.5, 45, 62.5 and 80°. The *inset plots* show zoomed-in line profiles for He II 1640 Å and H α . Double-peaked line emission can be seen in He II 1640 Å, He II 4686 Å, H α and some He I lines, but the line emission is not always sufficient to overcome the absorption cores from the stellar atmosphere models. The model also produces a prominent He II 3202 Å line at high inclinations 127

Fig. 4.5 Synthetic optical spectra from model A computed for sightlines of 10, 27.5, 45, 62.5 and 80°. In these plots the flux is divided by a polynomial fit to the underlying continuum redward of the Balmer edge, so that line-to-continuum ratios and the true depth of the Balmer jump can be shown 128

Fig. 4.6 Total packet-binned spectra across all viewing angles, in units of monochromatic luminosity. The *thick black line* shows the total integrated escaping spectrum, while the *green line* shows disc photons which escape without being reprocessed by the wind. The *red line* show the contributions from reprocessed photons. Recombination continuum emission blueward of the Balmer edge is already prominent relative to other wind continuum processes, but is not sufficient to fill in the Balmer jump in this specific model 129

Fig. 4.7 UV (*left*) and optical (*right*) synthetic spectra for model B computed at sightlines of 10, 27.5, 45, 62.5 and 80°. Model A is shown in *grey* for comparison. The *inset plots* show zoomed-in line profiles for He II 1640 Å and H α . The Balmer and He are double-peaked, albeit with narrower profiles. Strong He II 4686 Å emission can be seen, as well as a trend of a deeper Balmer jump with decreasing inclination 132

Fig. 4.8 Synthetic optical spectra from model B computed for sightlines of 10, 27.5, 45, 62.5 and 80°. Model A is shown in *grey* for comparison. In these plots the flux is divided by a polynomial fit to the underlying continuum redward of the Balmer edge, so that line-to-continuum ratios and the true depth of the Balmer jump can be shown 133

Fig. 4.9 Total packet-binned spectra across all viewing angles, in units of monochromatic luminosity. The *thick black line* shows the total integrated escaping spectrum, while the *green line* shows disc photons which escape without being reprocessed by the wind. The *red line* show the contributions from reprocessed photons. In this denser model the reprocessed contribution is significant compared to the escaping disc spectrum. The Balmer continuum emission is prominent, and the wind has a clear effect on the overall spectral shape 134

Fig. 4.10	<p>$H\alpha$ line profiles, normalized to 1, plotted in velocity space for three models with varying kinematic properties, computed at an inclination of 80°. The benchmark model and the improved optical model described in Sect. 4.3 are labeled as A and B respectively, and a third model (X) which has an increased acceleration length of $R_v = 283.8 R_{WD}$, and $\alpha = 4$ is also shown. The x-axis limits correspond to the Keplerian velocity at $4 R_{WD}$, the inner edge of the wind. There is a narrowing of the lines, and a single-peaked line in model X. This is not due to radial velocity shear (see Sect. 4.3.3)</p>	136
Fig. 4.11	<p><i>Top Panel</i> In and out of eclipse spectra of the high inclination NL RW Tri. <i>Bottom Panel</i> In and out of eclipse synthetic spectra from model B. The artificial ‘absorption’ feature just redward of the Balmer jump is due to the reasons described in Sect. 5.2</p>	138
Fig. 5.1	<p>A cartoon showing the geometry and some key parameters of the biconical quasar wind model.</p>	144
Fig. 5.2	<p>The input spectrum used for the quasar modelling. Note that although the X-ray power law is allowed to extend to very high frequencies, the radiation beyond $\sim 10^{-18}$ Hz is unimportant due to the rapid decline in the photoionization cross-sections at high energies. The results here are thus insensitive to the maximum frequency adopted for this power law component.</p>	146
Fig. 5.3	<p>Some typical length scales for the fiducial model. This places a formal limit of $\sim 10^{12}$ cm on clump sizes in the microclumping framework, and confirms that the cells are sufficiently larger than the Sobolev length in almost all cases.</p>	148
Fig. 5.4	<p>Contour plots showing the logarithm of some important physical properties of the outflow. The spatial scales are logarithmic and the x and z scales are not the same. The panels show the electron temperature, T_e, electron density, n_e, the ionization parameter, U_H, ion fractions, F, for three different ions and line luminosities, L, for two different lines. The <i>line</i> luminosities represent the luminosity of photons escaping the Sobolev region for each line. These photons do not necessarily escape to infinity. The <i>solid black line</i> marks a sphere at $1000 r_G$. The <i>dotted lines</i> show the 72° and 78° sightlines to the centre of the system, and illustrate that different <i>sightlines</i> intersect material of different ionization states.</p>	151
Fig. 5.5	<p>The electron density in the model, this time on linear axes in order to illustrate the density contrasts and scale of the system. The plot is on the scale of the acceleration length, whereas the <i>inset</i> is a box of $2700 \times 800 r_G$, where the <i>bottom left corner</i> corresponds to the base of the innermost streamline</p>	152

Fig. 5.6 Synthetic spectra at four viewing angles for the fiducial model. At 20° and 60° I show a comparison to an SDSS quasar composite from Reichard et al. (2003). At 73° and 76° I show a comparison to an *HST* STIS spectrum of the high BALnicity BALQSO PG0946+301 (Arav et al. 2000), and an SDSS spectrum of the LoBAL quasar SDSS J162901.63+453406.0, respectively. The *dotted line* shows a disc only continuum to show the effect of the outflow on the continuum level. All the spectra are scaled to the model flux at 2000 \AA , except for the *HST* STIS spectrum of PG0946+301, which is scaled to 1350 \AA due to the incomplete wavelength coverage 153

Fig. 5.7 A cartoon describing the broad classes of sightline in the fiducial model, illustrating how geometric effects lead to the different emergent spectra. The *colour gradient* is approximate, but indicates the stratified ionization structure, from highly ionized (*yellow*) to low ionization (*purple*) material 154

Fig. 5.8 Synthetic spectra at two viewing angles, this time in frequency space and including the optical band, compared to the non-BAL SDSS quasar composite. The spectra are normalised to the flux at 2000 \AA , then an offset of 2 is applied per spectrum for clarity—the *dotted lines* show the zero point of F_ν/F_{2000} in each case. 156

Fig. 5.9 X-ray (2 keV) luminosity of the clumped model (*red squares*) and the H13 model (*blue squares*), plotted against monochromatic luminosity at 2500 \AA . The *points* are labeled according to inclination; angles $> 70^\circ$ correspond to BALs in this scheme (see Fig. 5.7). Also plotted are measurements from the COMBI-7 AGN and the BQS samples (Steffen et al. 2006) and the Saez et al. (2012) sample of BALQSOs. The *dotted line* shows the best fit relation for non-BALQSOs from Steffen et al. (2006) 158

Fig. 5.10 C IV, Mg II, Al III and Fe II line profiles for viewing angles from 72 to 79° . The profiles are plotted relative to the local continuum with an offset applied for clarity. Lower ionization profiles appear at a subset of high inclinations, compared to the ubiquitous C IV profile. 159

Fig. 5.11 The EW of the C IV 1550 \AA line at 20° plotted against **a** the *BI* of C IV 1550 \AA at 75° , **b** the EW of the Mg II 2800 \AA line at 20° and **c** the EW of $\text{Ly}\alpha$ at 20° . The *circles* correspond to the simulation grid for two different values of f_ν , and the fiducial model is marked with an *orange star*. I also show a higher X-ray luminosity model and a higher mass loss rate with a *red star*. 160

Fig. 5.12	EW_{RW} as a function of inclination in the fiducial model, compared to 1/2 EW from the quasar sample of Di Pompeo et al. (2012; D12). The <i>shaded region</i> corresponds to the standard deviation of the D12 sample	162
Fig. 5.13	f_{FWHM} as a function of inclination from the fiducial model for three different lines. The Yong et al. (2016) predictions from two models, and the Pancoast et al. (2014a, b) modelling results for five Seyfert I galaxies are also shown. At inclination angles $\gtrsim 70^\circ$ radiative transfer effects can become important as the sightline looks into the flow; here I show only moderate to low inclination angles where the system would be observed as a non-BAL quasar.	163
Fig. 6.1	BH mass and Eddington fraction measurements from S11 for Sample A (<i>top</i>) and Sample B (<i>bottom</i>). The LoBALQSOs in sample A and BALQSOs in Sample B are also plotted in each case.	171
Fig. 6.2	Histograms of equivalent widths for three emission lines from the two different samples. The mean of the BAL and non-BAL quasar distributions are labeled in each case, and the histograms are normalised.	173
Fig. 6.3	Angular variation of continuum luminosity from AGNSPEC and classical thin disc models. The monochromatic continuum luminosities is divided by the monochromatic continuum luminosity at 10° , from AGNSPEC and classical thin disc models, at three different wavelengths. The models are computed for an Eddington fraction of 0.2 and $M_{BH} = 10^9 M_\odot$. In each panel I show both Kerr and Schwarzschild AGNSPEC models, and the classical models are for both pure foreshortened discs and foreshortened and limb darkened (LD) discs.	176
Fig. 6.4	Theoretical EW distributions from the numerical experiment described in Sect. 6.4.1 for a few different maximum angles. The results in the <i>top panel</i> use the same intrinsic distribution as Model 1 from R11 (shown in <i>black</i>), whereas the <i>bottom panel</i> shows the distributions obtained from a narrower intrinsic Gaussian. By the time the maximum angle is limited to around 70° the cutoff is clear even at moderate values of EW.	178
Fig. 6.5	The $EW[O\ III]$ distribution of quasars in the R11 sample (<i>black points</i>), with \sqrt{N} errorbars, and the best fit model with a maximum viewing angle of 84° . The intrinsic Gaussian distribution is shown with a <i>dashed line</i> . The plotted data is equivalent to the non-BAL histogram in the <i>top left panel</i> of Fig. 6.2, except that it uses linear binning and adopts the R11 sample criteria rather than using sample A	179

Fig. 6.6 $\Delta\chi^2$ as a function of maximum angle, θ_1 , calculated in steps of 0.1° . The choice for μ_* and σ_* is left free in each case. The *dotted line* marks the 3σ confidence interval. 179

Fig. 6.7 The geometry of the toy model used to carry out the Monte Carlo simulations. 180

Fig. 6.8 Heat map showing the results of the MC simulation described in Sect. 6.4.2. The quantities shown are discussed further in the text, but correspond to (clockwise from *top left*): the p_{KS} value from a comparison between the mock BAL dataset and the observed BAL dataset, the reduced χ^2 from the fit to the non-BAL EW distribution, the difference in mean EW between the mock BAL and mock non-BAL datasets, and the BAL fraction expected for the geometry in question 181

Fig. 6.9 Heat map showing the results of the MC simulation described in Sect. 6.4.2, but this time conducted with a Log-normal intrinsic distribution. The quantities shown are the same as in Fig. 6.8. Note the different scale for the *colour coding* of the χ^2/dof panel (*upper right*) 182

Fig. 6.10 Eigenvector 1 for LoBAL and non-BAL quasars. FWHM of the $H\beta$ line plotted against the relative iron strength, R_{FeII} . The *colour coding* corresponds to the EW of $[O III] 5007 \text{ \AA}$. The *dots* mark all quasars from sample A, while the *squares* mark those with Mg II LoBALs. A few of the Mg II LoBALQSOs are missing due to their lack of FWHM $[H\beta]$ measurements. The *arrows* show the approximate direction of the expected inclination (θ) trend under both the SH14 and R11 interpretations, and the expected trend in Eddington fraction (L/L_{Edd}) from SH14 only. 185

Fig. 6.11 LoBAL fraction compared to global LoBAL fraction in Eigenvector 1 space, in bins of $\Delta R_{FeII} = 1$ and $\Delta FWHM [H\beta] = 3000 \text{ km s}^{-1}$. The contour shows the outermost contour from Fig. 6.10 for reference. The text shows $N_{LoBAL}/N_{non-BAL}$, where N_{LoBAL} is the number of LoBALQSOs in the bin and $N_{non-BAL}$ in the number of non-BAL quasars in the bin 186

Fig. 6.12 Histograms of polarisation percentages for BAL quasars from Schmidt and Hines (1999) together with the Marin (2014) AGN sample 187

Fig. 6.13 Cumulative distribution functions of the histograms shown in Fig. 6.12 for BAL quasars from Schmidt and Hines (1999) together with the Marin (2014) AGN sample. The *colour-coding* and *x-axis scale* are the same as Fig. 6.12. The translucent *vertical lines* mark the median value in each sample. 187

- Fig. 6.14 The line angular emissivity function from a Keplerian disc as a function of inclination angle, θ , for a few different azimuthal angles, ϕ . The azimuthally-averaged case, $\bar{\epsilon}_{k,line}$ (*thick black line*), and the zero Keplerian velocity shear case, $\epsilon_{0,line}(\theta) = \cos \theta$ (*dotted line*), are also shown 190
- Fig. 6.15 The azimuthally-averaged line angular emissivity function, $\bar{\epsilon}_{k,line}$, from a Keplerian disc as a function of inclination angle for the four models shown in Table 6.1. The model parameters are the values of H / R , v_k and v_{th} . The zero Keplerian velocity shear case, $\epsilon_{0,line}(\theta) = \cos \theta$ (*dotted line*), is also shown. Unless the disc is very thin ($H/R \sim 0.001$), $\bar{\epsilon}_{k,line}$ shows large deviations from $\cos \theta$ 191
- Fig. 7.1 Ion fractions as a function of ionization parameter when Carbon is treated as a simple-atom (*dotted lines*) or macro-atom (*solid lines*) Although serious, as yet unsolved, numerical problems appear at high ionization parameters, the fundamental machinery for treating Carbon as a macro-atom is now in place 201
- Fig. 7.2 Velocity-resolved transfer function for $H\alpha$ from the fiducial quasar model presented in Chap. 5. The transfer function is for a viewing angle of 40° . *Credit Sam Mangham* 202

List of Tables

Table 1.1 Approximate values of compactness and accretion efficiency for four different compact objects with typical parameters. Here, R_* for black holes is set to the event horizon for a Schwarzschild black hole, whereas values of $\eta \sim 0.1$ in discs are more realistic due to the presence of an innermost stable orbit (see Sect. 1.1.3). SMBH stands for supermassive black hole. 3

Table 4.1 Parameters used for the geometry and kinematics of the benchmark CV model (model A), which is optimized for the UV band, and a model which is optimized for the optical band and described in Sect. 4.3 (model B). For model B, only parameters which are altered are given—otherwise the model A parameter is used. P_{orb} is the orbital period (the value for RW Tri from Walker 1963 is adopted, see Sect. 4.3.5) and R_2 is the radius of a sphere with the volume of the secondary’s Roche lobe. Other quantities are defined in the text or Fig. 4.1. Secondary star parameters are only quoted for model B as I do not show eclipses with the benchmark model (see Sect. 4.3.5) 125

Table 5.1 The grid points used in the parameter search. The sensitivity to some of these parameters is discussed further in Sect. 5.4.1 148

Table 5.2 Wind geometry parameters used in the fiducial model, as defined in the text and Fig. 5.1. Parameters differing from the benchmark model of H13 are highlighted with an asterisk 149

Table 5.3	Some derived spectral properties of the fiducial model, at 20° , compared to observations. The observed values are taken from the Shen et al. (2011) SDSS DR7 Quasar catalog, and correspond to mean values with standard deviations in log space from a subsample with $8.5 > \log(M_{BH}) < 9.5$ and $-1.5 < \log(L_{bol}/L_{Edd}) < 0$, where the BH mass is a C IV virial estimate. Units are logarithms of values in erg s^{-1}	155
Table 6.1	The values of the Keplerian velocity, v_k , thermal velocity, v_{th} , and ratio of disc scale height to radius, H/R , for four models. These values are used as inputs to calculate $\bar{\epsilon}_{k,line}$ as shown in Fig. 6.15, and model B is also used in Fig. 6.14	191