

A

Analysis of a Distributed Model for a MOS Capacitor

In this section, we will analyze the distributed model for the MOS capacitor, shown in Fig. A.1. The analysis here is modeled after [71].

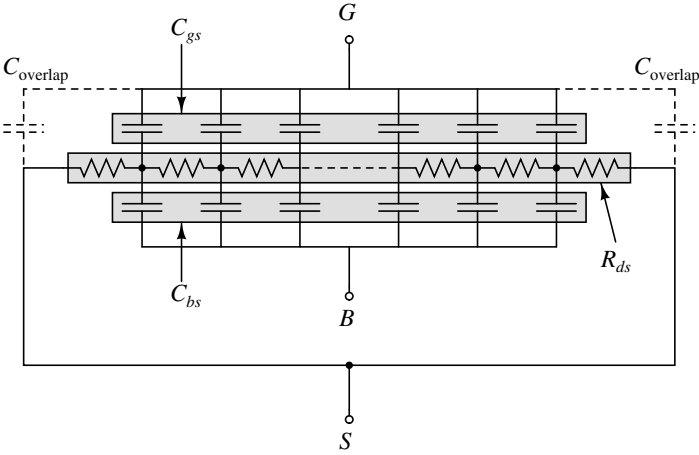


Fig. A.1: Distributed model for a MOS capacitor.

Of particular interest in this analysis, is the gate-source transadmittance, $\overline{Y_{gs}}(j\omega)$, which is defined as:

$$\overline{Y_{gs}} = \frac{\overline{I_g}}{\overline{V_s}} \Big|_{\overline{V_g}, \overline{V_b}=0} \quad (\text{A.1})$$

$\overline{I_g}$, $\overline{V_s}$ are shown in Fig. A.2. To evaluate $\overline{Y_{gs}}$, we can perform a thought experiment as shown in Fig. A.2. $\overline{I} = \overline{I_g} + \overline{I_b}$ is measured with an ammeter, when a signal $\overline{V_s}$ is applied at the drain-source. It can be seen that

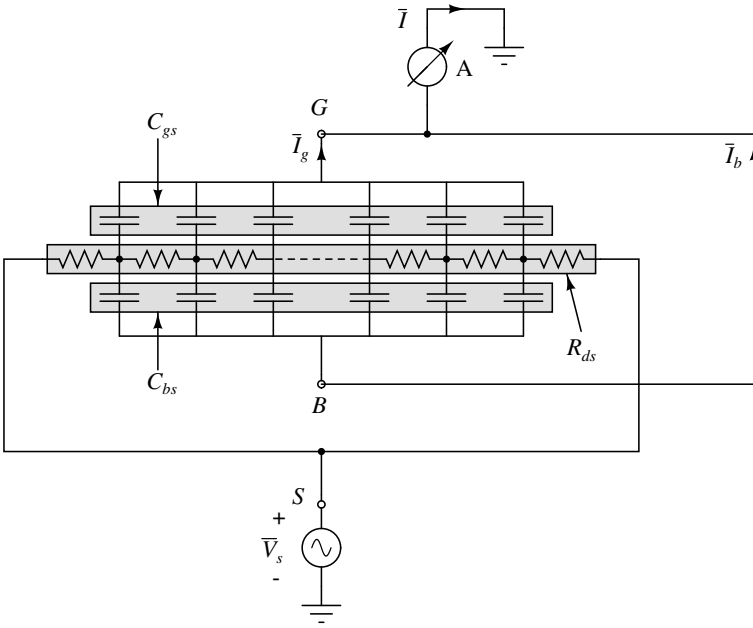


Fig. A.2: Thought experiment to measure \overline{V}_{gs} .

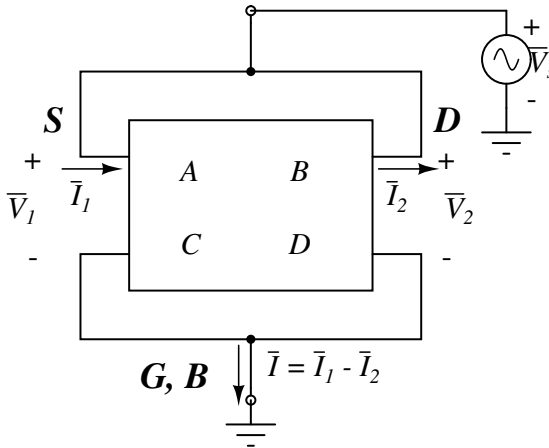


Fig. A.3: Replacing the MOS capacitor with its lumped transmission parameters.

$$\bar{I}_g = \frac{C_{gs}}{C_{gs} + C_{bs}} \bar{I} \tag{A.2}$$

The MOS capacitor in Fig. A.2 can be replaced with a circuit that has the same lumped transmission parameters as shown in Fig. A.3 with the correspondence in the terminals as shown. Then, \bar{Y}_{gs} can be represented in terms of the lumped transmission parameters in the following manner:

From the transmission matrix, we get:

$$\bar{V}_2 = A\bar{V}_1 + B\bar{I}_1 \tag{A.3}$$

$$\bar{I}_2 = C\bar{V}_1 + D\bar{I}_1 \tag{A.4}$$

Applying Kirchoff's equations to the circuit, we get:

$$\bar{V}_s = \bar{V}_1 = \bar{V}_2 \tag{A.5}$$

$$\bar{I} = \bar{I}_1 - \bar{I}_2 \tag{A.6}$$

From (A.3), (A.4) and (A.5), (A.6), we can deduce:

$$\bar{I}/\bar{V}_s = \frac{(1 - A)(1 - D)}{B} - C \tag{A.7}$$

\bar{Y}_{gs} can be obtained from (A.1), (A.2) and (A.7).

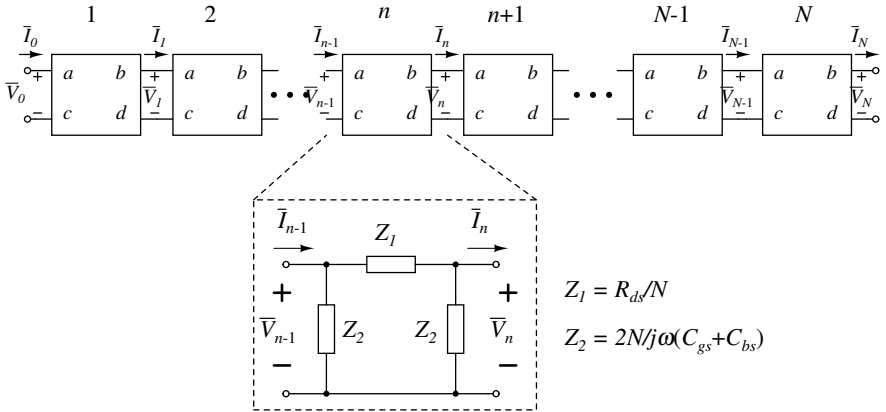


Fig. A.4: Expanding the distributed network into a cascade of N sections, $N \rightarrow \infty$.

The lumped transmission network in Fig. A.3 can be broken up into a cascade of N sections, where $N \rightarrow \infty$. This is shown in Fig. A.4. In each section, Z_1 and Z_2 are given by:

$$Z_1 = R_{ds}/N \tag{A.8}$$

$$Z_2 = 2N/j\omega(C_{gs} + C_{bs}) \tag{A.9}$$

For the n^{th} section,

$$\begin{bmatrix} \overline{V_{n+1}} \\ \overline{I_{n+1}} \end{bmatrix} = \begin{bmatrix} 1 + \frac{Z_1}{Z_2} & -Z_1 \\ -\frac{Z_1 + 2Z_2}{Z_2^2} & 1 + \frac{Z_1}{Z_2} \end{bmatrix} \begin{bmatrix} \overline{V_n} \\ \overline{I_n} \end{bmatrix} = [M] \begin{bmatrix} \overline{V_n} \\ \overline{I_n} \end{bmatrix}$$

And for the complete network,

$$\begin{bmatrix} \overline{V_N} \\ \overline{I_N} \end{bmatrix} = [M]^N \begin{bmatrix} \overline{V_0} \\ \overline{I_0} \end{bmatrix}$$

Note that $[M]^N$ is the transmission matrix for the entire network. Now let us identify the eigenvalues of the matrix $[M]$. If λ is an eigenvalue, then:

$$\det |[M] - \lambda[I]| \equiv 0$$

or,

$$\lambda^2 - 2\lambda(1 + Z_1/Z_2) + 1 = 0 \quad (\text{A.10})$$

(A.10) is the characteristic equation of the matrix $[M]$. If λ_1 and λ_2 are two solutions to this characteristic equation, and hence eigenvalues of $[M]$, then:

$$\lambda_1 \lambda_2 = 1 \quad (\text{A.11})$$

$$\lambda_1 + \lambda_2 = 2(1 + Z_1/Z_2) \quad (\text{A.12})$$

So, to satisfy (A.11),

$$\text{if } \lambda_1 = e^\tau \text{ and } \lambda_2 = e^{-\tau}$$

then to satisfy (A.12),

$$\cosh \tau = (1 + Z_1/Z_2) \quad (\text{A.13})$$

Now we can simplify the definition of $[M]$ using (A.13):

$$[M] = \begin{bmatrix} \cosh \tau & -Z_1 \\ -\frac{\sinh^2 \tau}{Z_1} & \cosh \tau \end{bmatrix} \quad (\text{A.14})$$

From the Cayley-Hamilton theorem,

$$[M]^N = C_0[I] + C_1[M] \quad (\text{A.15})$$

where C_0 and C_1 are constants, and

$$\lambda^N = C_0 + C_1 \lambda \quad (\text{A.16})$$

Inserting the specific eigenvalues in (A.16),

$$(e^\tau)^N = C_0 + C_1 e^\tau \quad (\text{A.17})$$

$$(e^{-\tau})^N = C_0 + C_1 e^{-\tau} \quad (\text{A.18})$$

From (A.17) and (A.18),

$$C_1 = \frac{\sinh N\tau}{\sinh \tau} \quad (\text{A.19})$$

$$C_0 = -\frac{\sinh(N-1)\tau}{\sinh \tau} \quad (\text{A.20})$$

From (A.15), (A.14) and (A.19), (A.20), we get:

$$[M]^N = \begin{bmatrix} \cosh(N\tau) & -Z_1 \frac{\sinh(N\tau)}{\sinh \tau} \\ -\frac{\sinh(N\tau) \sinh \tau}{Z_1} & \cosh(N\tau) \end{bmatrix} \quad (\text{A.21})$$

This gives us the transmission matrix for the complete network. In the limit where $N \rightarrow \infty$, (A.21) will be the transmission matrix for the distributed-RC network that we started with, as shown in Fig. A.3.

Now from (A.7) and (A.21), we obtain:

$$\bar{I}/\bar{V}_s = -\frac{[1 - \cosh(N\tau)]^2 \sinh \tau}{Z_1 \sinh(N\tau)} + \frac{\sinh \tau \sinh(N\tau)}{Z_1}$$

Simplifying further and using (A.8) and (A.9), we can deduce:

$$\bar{I}/\bar{V}_s = -\frac{2N}{R_{ds}} \cdot \frac{\sinh \tau}{\sinh(N\tau)} \cdot [1 - \cosh(N\tau)] \quad (\text{A.22})$$

where τ is given by (A.13), (A.8) and (A.9) as:

$$\tau = \cosh^{-1} \left[1 + \frac{j\omega R_{ds}(C_{gs} + C_{bs})}{2N^2} \right] \quad (\text{A.23})$$

As a sanity check, when $N = 1$, the network reduces to a simple lumped network, with the channel resistance, R_{ds} , shorted out. As a result, we expect the expressions (A.22), (A.13) to give us a purely capacitive result, i.e. $j\omega(C_{gs} + C_{bs})$.

From (A.23), in the limit when $N \rightarrow \infty$,

$$\lim_{N \rightarrow \infty} \tau = 0$$

From the definition of τ in (A.23),

$$N^2 \sinh^2(\tau/2) = j\omega R_{ds}(C_{gs} + C_{bs})/4$$

In the limit, as $N \rightarrow \infty$ and $\tau \rightarrow 0$,

$$\lim_{\substack{N \rightarrow \infty \\ \tau \rightarrow 0}} N^2 \sinh^2(\tau/2) = j\omega R_{ds}(C_{gs} + C_{bs})/4$$

Simplifying further, we obtain:

$$\lim_{\substack{N \rightarrow \infty \\ \tau \rightarrow 0}} N\tau = \sqrt{j\omega R_{ds}(C_{gs} + C_{bs})} = \alpha \quad (\text{A.24})$$

From (A.22), in the limit where $N \rightarrow \infty$ and also $\tau \rightarrow 0$:

$$\lim_{N \rightarrow \infty} \bar{I}/\bar{V}_s = \frac{2\alpha}{R_{ds} \sinh \alpha} (\cosh \alpha - 1) = \frac{\alpha^2}{R_{ds}} \cdot \frac{\tanh \alpha/2}{\alpha/2} \quad (\text{A.25})$$

This is rewritten, with the help of (A.24) to give:

$$\lim_{N \rightarrow \infty} \bar{I}/\bar{V}_s = j\omega(C_{gs} + C_{bs}) \cdot \frac{\tanh \alpha/2}{\alpha/2}$$

where α is given by (A.24).

Using (A.1) and (A.2), we can finally arrive at the desired expression for \bar{Y}_{gs} as:

$$\bar{Y}_{gs} = j\omega C_{gs} \cdot \frac{\tanh \alpha/2}{\alpha/2} \quad (\text{A.26})$$

where α is given by (A.24). As a second sanity check, when $\alpha \rightarrow 0$, (A.26) gives $j\omega C_{gs}$, which is expected.

For cases where N does not tend to ∞ , but is a large integer, note that $N\tau = \alpha$, as defined, as long as $\tau \rightarrow 0$. This means that the expression in (A.26) is approximately valid as long as N is large enough such that $\tau \rightarrow 0$. Looking at (A.13), this is true as long as:

$$\omega R_{ds}(C_{gs} + C_{bs}) \ll 2N^2 \quad (\text{A.27})$$

(A.27) can be used to determine the minimum number of discrete sections that will be required to perform a reasonably accurate simulation of this distributed network.

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