

Phenomenology of $U(1)_F$ extension of inert-doublet model with exotic scalars and leptons

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Received: 11 December 2017 / Accepted: 10 February 2018 / Published online: 21 February 2018
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Abstract In this work we will extend the inert-doublet model (IDM) by adding a new $U(1)_F$ gauge symmetry to it, under which, a Z_2 even scalar (ϕ_2) and Z_2 odd right handed component of two exotic charged leptons (F_{eR} , $F_{\mu R}$), are charged. We also add one Z_2 even real scalar (ϕ_1) and one complex scalar (ϕ), three neutral Majorana right handed fermions (N_1 , N_2 , N_3), two left handed components of the exotic charged leptons (F_{eL} , $F_{\mu L}$) as well as F_τ are all odd under the Z_2 , all of which are not charged under the $U(1)_F$. With these new particles added to the IDM, we have a model which can give two scalar DM candidates, together they can explain the present DM relic density as well as the muon (g-2) anomaly simultaneously. Also in this model the neutrino masses are generated at one loop level. One of the most peculiar feature of this model is that non-trivial solution to the axial gauge anomaly free conditions lead to the prediction of a stable very heavy partner to the electron (F_e), whose present collider limit (13 TeV LHC) on its mass should be around $m_{F_e} \geq \text{few TeV}$.

1 Introduction

The self consistency of the standard-model (SM) has been established (tentatively at least) with LHC discovery of the scalar behaving like the Higgs scalar of the SM. And about years ago since the SM has been formulated, it has went through numerous experimental test and no major disagreement with its prediction has been discovered. But still our universe turn out to be at least little more complicated than the SM can anticipate. One major discovery that point towards incompleteness of SM is the presence of dark-matter (DM). Another was the discovery of neutrino oscillation which means that neutrinos has small masses where as in SM neutrinos are massless. Then also it turn out that CP violation in SM due to Kobayashi-Maskawa (KM) [1] phase turn out to

be too small to explain the observed excess of matter over the anti-matter. So the most important questions that theoretical and experimental efforts in years to come will be related to nature of DM, mechanism behind the neutrino oscillation phenomena and search for new sources of CP violations.

In this work we will extend the SM by adding one more $U(1)_F$ gauge group to it with introducing only new exotic scalars and leptons. This simple extension turn out to explain the DM relic density, loop generated neutrino masses, Baryon genesis and muon (g-2) anomaly. One of the peculiar features of this particular realization of $U(1)_F$ extension of SM is that, due to solutions to the axial anomaly free conditions, if the electric charge of F_τ is taken as a free parameter, then the electric charge of F_e will have opposite sign to that of the F_μ , so if F_μ to explain the observed deviation in muon (g-2) then F_e require to be stable.

This paper is organized as follows. In the next section we introduce various particles in our model and write down Lagrangian invariant under all the symmetries imposed on the particles. Then in section III and IV we dwell into the consequences of the model related to muon (g-2), neutrino mass, DM, Baryon-Genesis and collider signature of the exotic charged leptons. We conclude in section V.

2 Model details

In this section we give details of the model. Our model introduces one additional $U(1)_F$ local gauge to the SM gauge group. So the full gauge group of our model is $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_F$. Only new exotic leptons and extra-scalars, which are singlet under the SM gauge group $SU(3)_c \times SU(2)_L$ are introduced to the IDM particle content. We introduce three charged leptons $F_{e,\mu,\tau}$ with each new leptons carrying respective lepton number 1 i.e for instance, lepton number of F_e is $(\mathcal{N}_e, \mathcal{N}_\mu, \mathcal{N}_\tau) = (1, 0, 0)$ etc. We require only the right handed component of the new

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Table 1 The charge assignments of new particles under the full gauge group $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_F$. Our choice of $U(1)_F$ charge of ϕ_2 is related to our choice of particular solution of the axial anomaly free equations. In this particular choice, ϕ_2 gives masses to F_e and F_μ and ϕ_1 gives mass to F_τ . Where index $i = e, \mu, \tau$ and our Y of SM $U(1)_Y$ is same as $\frac{Y}{2}$ in the usual convention

Particles	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)_F$	Z_2
F_{iR}	1	1	Y_i	n_i	-1
F_{iL}	1	1	Y_i	0	-1
N_{iR}	1	1	0	0	-1
ϕ	1	1	0	0	-1
ϕ_1	1	1	0	0	+1
ϕ_2	1	1	0	$n_{\phi_2} = n_\mu = -n_e$	+1
η	1	2	1/2	0	-1

leptons to be charged under the $U(1)_F$. These new leptons are odd under a discrete Z_2 symmetry. Also three more additional neutral right handed Majorana fermions N_{iR} are added which are also odd under the Z_2 but these neutral fermions does not carry $U(1)_F$ charges. In addition to those new fermions, we have the one inert doublet η which carries the same charges under the SM gauge group as the SM Higgs except that this new scalar doublet is odd under the Z_2 (Inert-doublet model (IDM)) whose VEV is zero. The new neutral fermions can have very large Majorana masses M_i given as $\bar{N}_{iR} M_i N_{iR}^c$, which, along with η_0 , can help generate neutrino masses at one loop [2]. We introduce another scalar ϕ , singlet under all the gauge symmetries except it is odd under the discrete Z_2 . We also require one real scalar, ϕ_1 , singlet under the full gauge group of the model and even under Z_2 which develops a non-zero VEV v_1 . Also we introduce one complex scalar, ϕ_2 , which is charged under the $U(1)_F$ and develops a non-zero VEV v_2 that breaks the $U(1)_F$ gauge symmetry spontaneously and give mass to the new gauge boson Z_F . We require the real scalar and the complex scalar just introduced because non-trivial solution to the anomaly free condition lead to one exotic lepton uncharged and other two exotic leptons charged under $U(1)_F$. So we require two scalars to generate masses for the fermions carrying different $U(1)_F$ charges including zero. The new particles and their charges under various symmetries are tabulated in the Table 1.

To make the theory free of axial anomaly the Y_i s and n_i s must solve the following equations,

$$\begin{aligned}
 \sum_{i=1}^3 Y_i^2 n_i &= 0 \\
 \sum_{i=1}^3 n_i^2 Y_i &= 0 \\
 \sum_{i=1}^3 n_i^3 &= 0
 \end{aligned}
 \tag{1}$$

which are the anomaly free conditions coming from $U(1)_Y^2 U(1)_F$, $U(1)_F^2 U(1)_Y$ and $U(1)_F^3$ respectively. But there is one more anomaly free condition due to gravity as $Gravity^2 U(1)_F$ which gives

$$\sum_{i=1}^3 n_i = 0.
 \tag{2}$$

One simplest solutions of Eqs. (1) and (2) for Y_i s and n_i s is trivial solution i.e $n_e = n_\mu = n_\tau = 0$. The trivial solution is same as not introducing the $U(1)_F$ gauge symmetry at all as all the new leptons decouples from the Z_F gauge boson. As it will be clear in the following paragraphs, the advantage of introducing the $U(1)_F$ is that non-trivial solutions to the axial anomaly free conditions given above require that one of the exotic leptons is forbidden to have Yukawa coupling with its corresponding SM lepton, and therefore should not show any deviation in the (g-2) of the corresponding SM lepton. Since (g-2) of electron has shown no deviation from the SM prediction, it is clear that $y_e \bar{e}_R \phi F_e$ should be the forbidden one.

A non-trivial solution of Eqs. (1) and (2) for Y_i s and n_i s that is interesting is given when we set either n_e or n_μ or n_τ equal to zero. In this work we set $n_\tau = 0$ which make the electric charge of F_τ a free parameter. Then the four equations are solved for Y_e, Y_μ, n_e and n_μ if we set $Y_e = -Y_\mu$ and $n_e = -n_\mu$, but for F_μ to explain the observed anomaly in δa_μ [3], we require $Y_\mu = Q_{F_\mu} = -1$ which implies that $Y_e = Q_{F_e} = +1$ and so charge conservation and Lorenz invariance forbid F_e to contribute to δa_e , which is consistent with the experimental findings as no deviation in δa_e from the SM prediction has been reported. If we require Yukawa terms such as $y_\mu \bar{\mu}_R F_{\mu L} \phi$, to explain the observed anomaly in muon (g-2), then term such $y_e \bar{e}_R F_{eL} \phi$ are forbidden by charge conservation and term such as $y_e \bar{e}_R F_{eL}^c \phi$ is forbidden by Lorenz invariance, so F_e will be a stable heavy charged particle. Although Y_τ is a free parameter, if we also set $Y_\tau = Q_{F_\tau} = -1$, then we can expect a deviation in δa_τ in future measurements coming from $y_\tau \bar{\tau}_R \phi F_\tau$ term.

2.1 Scalar sector

The general Lagrangian invariant under all the symmetries of the model in the scalar sector can be written as

$$\begin{aligned}
 \mathcal{L}_{scalar} &= |D_\mu^F \phi_2|^2 + |D_\mu H|^2 + |D_\mu \eta|^2 + |\partial_\mu \phi|^2 + \partial_\mu \phi_1|^2 \\
 &\quad - V(H, \eta, \phi, \phi_1, \phi_2)
 \end{aligned}
 \tag{3}$$

where $D_\mu^F = \partial_\mu - i n_{\phi_2} g_F Z_{F\mu}$ and $D_\mu = \partial_\mu - i g \sigma \cdot W_\mu - i Y g' B_\mu$ with H being the SM Higgs scalar, η being the inert-doublet and $Z_{F\mu}$ and B_μ being gauge fields of $U(1)_F$ and

$U(1)_Y$ respectively with ϕ, ϕ_1 and ϕ_2 being new SM singlet scalars. The scalar potential can be expressed as

$$\begin{aligned}
 V(H, \phi, \phi_1, \phi_2) = & m^2 H^\dagger H + m_\eta^2 \eta^\dagger \eta + m_\phi^2 \phi^\dagger \phi + m_{\phi_1}^2 \phi_1^\dagger \phi_1 \\
 & + m_\phi^2 \phi_2^\dagger \phi_2 + \lambda_1 (H^\dagger H)^2 + \lambda_2 (\eta^\dagger \eta)^2 \\
 & + m_{H\eta\phi} (H^\dagger \eta + \eta^\dagger H) \phi + \lambda_3 (H^\dagger H) (\eta^\dagger \eta) \\
 & + \lambda_4 |H^\dagger \eta|^2 + \frac{\lambda_5}{2} \{ (H^\dagger \eta)^2 + h.c \} \\
 & + \lambda_\phi (\phi^\dagger \phi)^2 + \lambda_{\phi_1} (\phi_1^\dagger \phi_1)^2 + \lambda_{\phi_2} (\phi_2^\dagger \phi_2)^2 \\
 & + \lambda_{H\phi} (H^\dagger H) (\phi^\dagger \phi) \\
 & + \lambda_{H1} (H^\dagger H) (\phi_1^\dagger \phi_1) + \lambda_{H2} (H^\dagger H) (\phi_2^\dagger \phi_2) \\
 & + \lambda_{\eta\phi} (\eta^\dagger \eta) (\phi^\dagger \phi) + \lambda_{\eta\phi_1} (\eta^\dagger \eta) (\phi_1^\dagger \phi_1) \\
 & + \lambda_{\eta\phi_2} (\eta^\dagger \eta) (\phi_2^\dagger \phi_2) + \lambda_{1\phi} (\phi_1^\dagger \phi_1) (\phi^\dagger \phi) \\
 & + \lambda_{2\phi} (\phi_2^\dagger \phi_2) (\phi^\dagger \phi) + \lambda_{12} (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2) \quad (4)
 \end{aligned}$$

where we have $H^T = (G^+, \frac{v_0+h+iG^0}{\sqrt{2}}), \eta^T = (H^+, \frac{\eta_R+i\eta_I}{\sqrt{2}}), \phi = \frac{\phi_R+i\phi_I}{\sqrt{2}}, \phi_1 = v_1 + \sigma$ and $\phi_2 = \frac{v_2+S+iG^s}{\sqrt{2}}$ with G^\pm, G^0 and G^s being the Nambu-Goldstone bosons. And $v_0 \approx 246$ is the SM Higgs VEV and similarly v_1 and v_2 are the VEV of real field ϕ_1 and the complex field ϕ_2 respectively. We take the couplings in the limit such that the SM Higgs H and the inert-doublet η decouples from the other new singlet scalars i.e $(\lambda_{H1}, \lambda_{H2}, \lambda_{H\phi}) \approx (\lambda_{\eta\phi_1}, \lambda_{\eta\phi_2}, \lambda_{\eta\phi}) \approx 0$. Also with hindsight, to explain the muon (g-2) anomaly and DM, we require the mixing between η and ϕ to be very small, and since terms containing the $m_{H\eta\phi}$ mix the η and ϕ after Higgs develops a nonzero VEV, we set $m_{H\eta\phi} = 0$ in this work so that no mixing between η and ϕ is introduced. Then the mass of the Higgs particle is

$$M_h^2 = -2m^2 = 2\lambda_1 v_0^2 \quad (5)$$

where as the masses of the charged Higgs H^\pm and two neutral scalars η_R^0 and η_I^0 are given by

$$\begin{aligned}
 M_{H^\pm}^2 &= m_\eta^2 + \lambda_3 v_0^2/2, \\
 M_{\eta_R^0}^2 &= m_\eta^2 + (\lambda_3 + \lambda_4 + \lambda_5) v_0^2/2, \\
 M_{\eta_I^0}^2 &= m_\eta^2 + (\lambda_3 + \lambda_4 - \lambda_5) v_0^2/2. \quad (6)
 \end{aligned}$$

As we can see from the above equation, which of the two (η_R^0 or η_I^0) is the lighter depends on the sign of the λ_5 , if $\lambda_5 < 0$ then η_R is the lighter one and so DM candidate and vice versa.

2.2 Fermionic sector

As mentioned before, the symmetries of the model allows the neutral fermions N_{iR} to have very large Majorana masses given as $\bar{N}_{iR} M_i N_{iR}$ where M_i is a 3×3 Majorana mass matrix. We do not entertain in this work how such a very

heavy Majorana masses can be generated which could come from unified theories such as supergravity etc., see [4–7]. The symmetries of the model also allows Yukawa terms such as $h_{ij} \bar{L}_i i \sigma_2 \eta N_{iR}$, which along with the Majorana mass terms for the N_{iR} can generate neutrino masses at one loop [2]. We can also have Yukawa terms such as $y_i \bar{l}_{iR} \phi F_{iL}$ where l_{iR} 's are SM right handed charged leptons. But non trivial solutions to the $U(1)_F$ anomaly free conditions restrict at maximum only two of the three leptons is allowed to have such Yukawa interactions. Since no deviation has been reported in the electron (g-2) other than the SM predicted value, in the following sections we will assume that $F_{\mu L}$ and $F_{\tau L}$ have electric charge given by $Q_{F_{\mu,\tau}} = Y_{F_{\mu,\tau}} = -1$ and then the axial gauge anomaly free conditions sets $Q_{F_e} = Y_{F_e} = +1$, so then only muon and tau (g-2) has contributions from respective exotic leptons in our model. The general fermionic sector Lagrangian invariant under the full symmetries of the model can be written as

$$\begin{aligned}
 \mathcal{L}_{fermion} = & \sum_{i=e}^{\tau} \bar{F}_i \gamma^\mu (i D_{i\mu}) F_i + \sum_{i=e}^{\tau} \bar{F}_i \gamma^\mu (i \partial_\mu P_L) F_i \\
 & + \sum_{i=e}^{\tau} \bar{N}_{Ri} M_i N_{Ri}^c \\
 & + \sum_{i,j=e}^{\tau} h_{ij} \bar{L}_i i \sigma_2 \eta^* N_{jR} + \sum_{i=\mu}^{\tau} y_i \bar{l}_{iR} \phi F_{iL} \\
 & + \sum_{i=e}^{\mu} g_{F_i} \bar{F}_{iR} \phi_i F_{iL} + g_{F_\tau} \bar{F}_{\tau R} \phi_1 F_{\tau L} + h.c \quad (7)
 \end{aligned}$$

where $D_{i\mu} = \partial_\mu P_R - i Y_i g' B_\mu - i n_i g_{F_i} Z_{F\mu} P_R$ with $P_{R/L} = \frac{1}{2}(1 \pm \gamma_5)$ and $\phi_e^\dagger = \phi_\mu = \phi_2$ as $n_e = -n_\mu$ from the solutions of the anomaly free conditions. The B_μ is the SM $U(1)_Y$ gauge boson and $Z_{F\mu}$ being the new $U(1)_F$ gauge boson. We can express the B_μ in terms of SM Z_μ and A_μ as

$$B_\mu = -\sin \theta_W Z_\mu + \cos \theta_W A_\mu, \quad (8)$$

so then we can express the interaction of the new exotic leptons with the SM weak gauge boson Z_μ and photon A_μ as

$$\sum_{i=e}^{\tau} \bar{F}_i \gamma^\mu Y_i g' B_\mu F_i = e \sum_{i=e}^{\tau} Y_i \bar{F}_i (-\tan \theta_W Z_\mu + A_\mu) \gamma^\mu F_i \quad (9)$$

where $g' = \frac{e}{\cos \theta_W}$. From the above Eq. (9), we can see that the production of these new leptons at LHC will be via the process $pp \rightarrow Z^*/\gamma^* \rightarrow F_i^+ F_i^-$. Even though the production of the three exotic leptons are similar in nature, their signature mainly coming from their decays differs drastically. This is because, since $y_e \bar{e}_R \phi F_{eL}$ term is forbidden by the axial gauge anomaly free conditions, in this model, the exotic

lepton F_e will be a stable charged particle just like its SM counterpart, the electron. The other two exotic leptons, F_μ and F_τ , are not stable as they can decay via $F_{\mu/\tau} \rightarrow \mu/\tau + \phi$. Due to unbroken Z_2 symmetry, ϕ will be a stable particle if $m_{F_\mu} > m_\phi + m_\mu$ is satisfied and so ϕ can be a DM candidate. But we will see in Sect. 3, if $y_\mu \bar{\mu}_R \phi F_{\mu L}$ to explain the muon (g-2) within 1σ of the experimental value which require large y_μ , then contribution of ϕ to the present relic density of DM will be negligibly tiny. In early universe, most of the F_e would have been depleted via $F_e^+ F_e^- \rightarrow Z_F/Z^*/\gamma^* \rightarrow F_\mu^+ F_\mu^- (F_\tau^+ F_\tau^-) \rightarrow \mu^+ \mu^- (\tau^+ \tau^-) + \text{missing energy } (\phi\phi)$. The collider signature of F_e will be basically the collider signature of a very heavy stable charged particle, and given LHC reaching 13 TeV reported no such signatures of a stable heavy charged particle, we expect the mass of the F_e to be larger than few TeV. However for the other two exotic leptons, the main collider signatures will be

$$e^+ e^- / pp \rightarrow Z^*/\gamma^* \rightarrow F_\mu^+ F_\mu^- (F_\tau^+ F_\tau^-) \rightarrow \mu^+ \mu^- (\tau^+ \tau^-) + \text{missing energy } (\phi\phi), \tag{10}$$

and other key collider signatures will be

$$e^+ e^- / pp \rightarrow Z^*/\gamma^* \rightarrow \gamma + \phi_1/\phi_2 (\text{via triangle loop}) \rightarrow \gamma + \text{missing energy } (\phi_1/\phi_2 \rightarrow \phi\phi) \tag{11}$$

and

$$e^+ e^- / pp \rightarrow Z^*/\gamma^* \rightarrow \gamma + \phi_1/\phi_2 (\text{via triangle loop}) \rightarrow \gamma + (\gamma + \gamma (\text{or } Z)) (\text{via triangle loop}). \tag{12}$$

In this sense we think $e^+ e^-$ colliders such as ILC (being a precision machine) will be more sensitive in detecting the missing energies in the final states of above reactions than Hadron colliders such as LHC.

3 Muon (g-2) anomaly

The contribution from the term $y_\mu \bar{\mu}_R \phi F_{\mu L}$ from Eq. (7) to the muon (g-2) anomaly can be expressed as [8]

$$\delta a_\mu = \frac{m_\mu^2 y_\mu^2}{16\pi^2} \int_0^1 dx \frac{x^2 - x^3}{m_\mu^2 x^2 + (m_{F_\mu}^2 - m_\mu^2)x + m_\phi^2(1-x)}. \tag{13}$$

And as pointed out in [3], in this type of models where the scalar in the loop does not develops VEV, the constrains from the heavy charged lepton searches, which has rule out at 90% C.L for $m_{F_{Heavy}} \leq 100.8$ GeV, does not apply because the heavy charged leptons in our model can not decay into the final state those experiments looked for such as $F^\pm \rightarrow W^\pm \nu$ and $F^\pm \rightarrow l^\pm Z$. Therefore in our model the new charged lepton (F_μ) can have mass less the 100 GeV, but the mass of

F_μ must satisfy $m_{F_\mu} > m_Z/2$ so that no contribution from F_μ to the Z decay width (via $Z \rightarrow \bar{F}_\mu F_\mu$), which is very precisely measured and no deviation from SM prediction has been found. Similarly if $m_\phi < m_Z/2$ then also contribution to the invisible Z decay width ($Z \rightarrow \phi\phi$ via triangle loop) is expected to be large and so we require $m_\phi > m_Z/2$. In agreement with the experimental non-observation of long lived heavy charged particles, we require $m_{F_\mu} \geq m_\phi + m_\mu$ with large Yukawa coupling, i.e $y_\mu \approx 1$ so that F_μ decays quickly after production as $F_\mu \rightarrow \phi + \mu$.

All these constrains on the mass can be satisfied along with explaining the muon (g-2) anomaly within 1σ of the δa_μ^{Exp} if we have

$$\frac{y_\mu}{m_{F_\mu}} \approx 0.0188 \text{ GeV}^{-1}, \tag{14}$$

where we have taken the approximation of $m_\phi \approx m_{F_\mu} \gg m_\mu$ in the Eq. (13). Then for $y_\mu \approx 1$, we have $m_\phi \approx m_{F_\mu} \approx 53$ GeV, and for $y_\mu \approx 2\sqrt{\pi}$, we have $m_\phi \approx m_{F_\mu} \approx 189$ GeV, assuming the condition $m_{F_\mu} > m_\phi + m_\mu$ is satisfied. Due to unbroken Z_2 symmetry, ϕ will be a DM candidate, but as pointed out in [9], for such large Yukawa couplings the contributions to the present relic density of DM by ϕ will be negligibly tiny. In case of fermion F_μ affecting the Z decay $Z \rightarrow \bar{\mu}\mu$ via the triangle loop, by adapting the formula II.39 of [10], we have from the triangle loop contribution to $Z \rightarrow \bar{\mu}\mu$ given as

$$Br(Z \rightarrow \bar{\mu}\mu)_{triangle} = \frac{G_F}{3\sqrt{2}\pi} \frac{m_Z^3}{(16\pi^2)^2 \Gamma_Z} s_w^2 t_w^2 y_\mu^4 |F_2(F_\mu, \phi) + F_3(F_\mu, \phi)|^2 \tag{15}$$

where $t_w = \frac{s_w}{\sqrt{1-s_w^2}}$ and $s_w = \sin \theta_w$ with θ_w being the Wienberg angle and

$$F_2(F_\mu, \phi) = \int_0^1 dx (1-x) \ln[(1-x)m_{F_\mu}^2 + xm_\phi^2],$$

$$F_3(F_\mu, \phi) = \int_0^1 dx \int_0^{1-x} dy \times \frac{(xy-1)m_Z^2 + (m_{F_\mu}^2 - m_\phi^2)(1-x-y) - \Delta \ln \Delta}{\Delta} \tag{16}$$

with $\Delta = -xym_Z^2 + (x+y)(m_{F_\mu}^2 - m_\phi^2) + m_\phi^2$ and total Z width $\Gamma_Z = 2.4952 \pm 0.0023$ GeV. For $m_{F_\mu} \approx 53 \approx m_\phi$ and $y_\mu = 1$, we have $Br(Z \rightarrow \bar{\mu}\mu)_{triangle} = 2.19 \times 10^{-6}$ compared to the $Br(Z \rightarrow \bar{\mu}\mu)_{Exp} = (3.366 \pm 0.007)\%$ [11], so the triangle loop contribution from the new particles to $Br(Z \rightarrow \bar{\mu}\mu)$ is an order of magnitude smaller than the error in the experimental value.

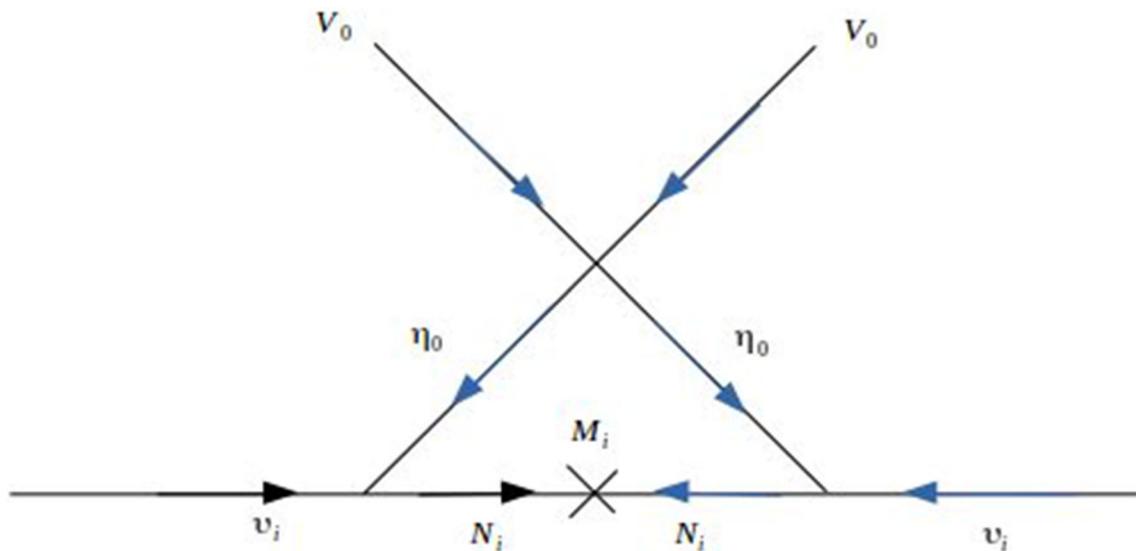


Fig. 1 Neutrino mass generation at one loop level. Where here V_0 is the VEV of SM Higgs and η_0 denoting either η_R or η_I

4 Loop generation of neutrino masses and dark-matter

It has been shown more than 10 years ago in [12], that with addition of three neutral Majoran fermions which are also odd under the Z_2 symmetry to the inert-doublet model (IDM), the origins of small neutrino mass, DM and Baryogenesis can all be linked and explained. The small Majorana neutrino masses for the left handed neutrinos can arise at one loop level via the scotogenic mechanism [2] as shown in Fig. 1, where the neutral components of the inert-doublet and N_i propagating in the loop inducing indirect interaction of the left handed neutrinos with the Higgs vacuum and M_i and generate a mass term $M_{\alpha\beta} \bar{\nu}_\alpha^c \nu_\beta$. In the scotogenic model the loop generated neutrino mass matrix $M_{\alpha\beta}$ can be expressed as [2]

$$M_{\alpha\beta} = \sum_i \frac{h_{\alpha i} h_{\beta i} M_i}{16\pi^2} \times \left[\frac{m_R^2}{m_R^2 - M_i^2} \ln \frac{m_R^2}{M_i^2} - \frac{m_I^2}{m_I^2 - M_i^2} \ln \frac{m_I^2}{M_i^2} \right], \quad (17)$$

where denoting the inert-doublet as $\eta^T = (\eta^+, \eta^0)$ with $\eta^0 = \eta_R^0 + i\eta_I^0$ and $m_{R,I}$ being the masses of the the $\eta_{R,I}^0$ respectively. The lightest one among the N_i and $\eta_{R,I}^0$ is the DM. In this work we have taken the η_R^0 as the lightest and therefore DM candidate. However as shown in [12], if one of the $\eta_{R,I}^0$ is the DM then there is also the possibility of generating Baryon asymmetry by converting the lepton-asymmetry into baryon asymmetry via sphalerons [13]. As a numerical estimate, for $|h_{11}|^2 \approx 1.14 \times 10^{-5}$ and the the term in the bracket in Eq. (17) being equal to $\frac{v_0^2}{M_1^2}$, where N_{R1} being the lightest of the heavy neutrinos and $v_0 \approx 185$ GeV, then for

$M_1 \approx 2.6 \times 10^7$ GeV we get the masses of the light neutrinos as $m_{11} = M_{11} \approx O(0.1)$ eV and also these parameters will be able to generate required baryon excess [7]. Also in this model constrains from small neutrino masses has little effect on determining the DM annihilation cross section and therefore determining the DM relic density.

The neutral Lightest Stable Particle (LSP) of the IDM can be a DM candidate has been first pointed out in [2]. We take $\lambda_5 < 0$ and so of the two neutral scalars η_R and η_I , η_R is the DM candidate in our model. In [14–16], before the Higgs discovery at LHC, it has been shown that in detail analysis η_R can account the present relic density of DM if the masses of the scalar DM lies in two regions, $m_{\eta_R} \leq m_W \approx 80$ Gev or $m_{\eta_R} \geq 500$ Gev assuming the Higgs mass as 120 GeV or 200 GeV. However, in [17, 18] the authors did a thorough analysis of IDM DM after discovery of the SM like Higgs at $m_h \approx 126$ GeV, combine with direct searches and LHC Run One data, they have ruled out the low and intermediate mass region of IDM DM except in a very narrow strip near the $m_{\eta_R} = m_h/2$, corresponding to very small values of the $h - \eta_R - \eta_R$ coupling (the funnel regions). The recent LUX [19] and PandaX [20] data reaffirm their conclusions more strongly. Therefore we will focus on the high mass region where the DM mass is $m_{\eta_R} \geq 500$ GeV. In this region if $m_\eta < m_{\eta_R}$, then the annihilation into two gauge bosons dominate, and if $m_{\eta_R} > m_\eta$, then annihilation into two Higgs bosons dominate. It turn out coannihilation plays little role in this region. In the limit of small mass splitting and $\lambda_L \approx 0$, the relic density can be accounted for $m_{\eta_R} \geq 500$ GeV [16] where $\lambda_L = (\lambda_3 + \lambda_4 + \lambda_5)/2$. The contribution to relic density tends to increase as m_{η_R} increase but this can be controlled by increasing the mass splitting as annihilation cross section tends to increase as mass splitting increases

[14–16]. However, in DM theories (including ours) where it is assumed that in some point in the early universe DM particles are in thermal equilibrium with the rest of the matter, their annihilation cross section is bound by the partial waves unitarity of the S-matrix, which in turn constrains the relic density and mass of the particle. The unitarity bounds on total annihilation cross section for a scalar DM particle is given as [21]

$$\langle \sigma v \rangle_{v \rightarrow 0}^{\text{unitary}} \approx \frac{4\pi}{m_{\eta_R}^2} \sqrt{\frac{x_f}{\pi}}, \quad (18)$$

where $x_f = \frac{m_{\eta_R}}{T_f}$ with T_f being the freeze out temperature. It has been shown in [21] that both unitarity and WMAP constrains can be satisfied for the scalar DM mass $m_{\eta_R} \leq 130$ TeV.

5 Conclusions

In this work we have proposed a simple extension of the SM by only adding three Z_2 odd exotic charged leptons (F_e, F_μ, F_τ) whose right handed component are charged under a new $U(1)_F$ gauge symmetry, three Z_2 odd neutral fermions (N_1, N_2, N_3) singlet under the entire gauge symmetry, one $SU(2)_L$ doublet (η) odd under a discrete Z_2 symmetry whose VEV is zero (the inert-doublet model), one Z_2 odd scalar (ϕ) singlet under the entire gauge group and whose VEV is zero, one Z_2 even scalar (ϕ_1) also singlet under the entire gauge group which develops a non zero VEV v_1 , one more Z_2 even scalar (ϕ_2) charged under the $U(1)_F$ whose non-zero VEV breaks the $U(1)_F$ gauge symmetry spontaneously and give mass to the gauge boson Z_F . With addition of the above new particles, we have been able to build a model which can give two DM candidate in terms of lightest neutral component of the inert-doublet (η_R in our case) and ϕ but if ϕ to explain the muon (g-2) anomaly than ϕ will contribute only a tiny fraction to the present DM relic density, so most of the present relic density of DM will consist of η_R whose mass can be in the range 500 GeV to 130 TeV. In this model neutrino masses are generated at one loop level as well as baryon-genesis via lepto-genesis is also possible. We have also given key signatures of these new exotic leptons at LHC and future e^+e^- colliders. We find the the key signature will be in the form of $e^+e^-/pp \rightarrow Z^*\gamma^* \rightarrow F_\mu^+ F_\mu^- (F_\tau^+ F_\tau^-) \rightarrow \mu^+ \mu^- (\tau^+ \tau^-) + \text{missing energy} (\phi\phi)$

and $e^+e^-/pp \rightarrow Z^*\gamma^* \rightarrow \gamma + \phi_2$ (via triangle loop) $\rightarrow \gamma + \text{missing energy} (\phi_2 \rightarrow \phi\phi)$. One of the most peculiar signature of this model is the existence of a stable very heavy (about $m_{F_e} \geq \text{few TeV}$) charged lepton partner to the electron.

Acknowledgements This work is supported and funded by the Department of Atomic Energy of the Government of India and by the Government of U.P.

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