

Large deviations for the Sine_{β} and Sch_{τ} processes

Diane Holcomb · Benedek Valkó

Received: 18 November 2013 / Revised: 3 November 2014 / Published online: 22 November 2014 © Springer-Verlag Berlin Heidelberg 2014

Abstract The $\operatorname{Sine}_{\beta}$ process is the bulk point process limit of the Gaussian β -ensemble. For $\beta=1,2$ and 4 this process gives the limit of the GOE, GUE and GSE random matrix models. The $\operatorname{Sch}_{\tau}$ process is obtained similarly as the bulk scaling limit of the spectrum of certain discrete one-dimensional random Schrödinger operators. Both point processes have asymptotically constant average density, in our chosen normalization one expects close to $\frac{1}{2\pi}\lambda$ points in a large interval of length λ . We prove large deviation principles for the average densities of the processes, identifying the rate function in both cases. Our approach is based on the representation of the counting functions using coupled systems of stochastic differential equations. Our techniques work for the full range of parameter values. The results are novel even in the classical $\beta=1,2$ and 4 cases for the $\operatorname{Sine}_{\beta}$ process. They are consistent with the existing rigorous results on large gap probabilities and confirm the physical predictions made using log-gas arguments.

Mathematics Subject Classification 60B20 · 60F10 · 15B52

1 Introduction

The Gaussian orthogonal, unitary and symplectic ensembles (GOE, GUE, GSE) are some of the most studied random matrix models. These are symmetric (resp. Hermitian or symplectic) matrices with i.i.d. standard real (resp. complex or quaternion) normal

Department of Mathematics, University of Arizona, Tucson, AZ, USA e-mail: dianeholcomb@math.arizona.edu

B. Valkó (⊠)

Department of Mathematics, University of Wisconsin - Madison, Madison, WI, USA e-mail: valko@math.wisc.edu



D. Holcomb

entries modulo the appropriate symmetry. It has been know from the classical results of Gaudin and Mehta [20] that if we appropriately scale the spectrum in the bulk (e.g. near zero) then we obtain a limiting point process. The point process can be described via its n-point correlation functions. These are given by determinantal formulas in the GUE case and Pfaffian formulas in the GOE, GSE cases. (See [2,13,20] for more details and the precise description.)

The GOE, GUE, GSE models can be naturally included in a one-parameter family of distributions. The joint eigenvalue distribution for these classical models is known to be

$$p(\lambda_1, \dots, \lambda_n) = \frac{1}{Z_{n,\beta}} \prod_{1 < i < j < n} |\lambda_i - \lambda_j|^{\beta} e^{-\frac{\beta}{4} \sum_{i=1}^n \lambda_i^2}$$
(1)

where β is equal to 1, 2 and 4 in the three cases. Note, that the constant $\beta/4$ in the exponential can be easily changed via linear scaling. It is natural to consider the density (1) for general $\beta > 0$, this is the Gaussian (or Hermitian) β -ensemble. In [23] the authors show the existence of the bulk scaling limit for general β . In particular, if $\Lambda_{n,\beta}$ is distributed according to (1) then $2\sqrt{n}\Lambda_{n,\beta}$ converges to a random point process, denoted by Sine $_{\beta}$. For $\beta = 1, 2, 4$ this gives the bulk limit process for the GOE, GUE, GSE ensembles.

The $Sine_{\beta}$ process can be described through its counting function using a system of stochastic differential equations. Consider the system

$$d\alpha_{\lambda} = \lambda \frac{\beta}{4} e^{-\frac{\beta}{4}t} dt + \text{Re}\left[(e^{-i\alpha_{\lambda}} - 1)(dB_1 + idB_2) \right], \quad \alpha_{\lambda}(0) = 0, \quad t \in [0, \infty)$$
 (2)

where B_1 , B_2 are independent standard Brownian motions. Note, that this is a one-parameter family of SDEs driven by the same complex Brownian motion. In [23] it was shown that $N_{\beta}(\lambda) = \lim_{t \to \infty} \frac{1}{2\pi} \alpha_{\lambda}(t)$ exists almost surely and it is an integer valued monotone increasing function in λ . Moreover, the function $\lambda \to N_{\beta}(\lambda)$ has the same distribution as the counting function of the Sine_{\beta} process, i.e. the distribution of the number of points in $[0, \lambda]$ for $\lambda > 0$ is given by that of $N_{\beta}(\lambda)$.

Note, that for any fixed λ the process α_{λ} satisfies the SDE

$$d\alpha_{\lambda} = \lambda \frac{\beta}{4} e^{-\frac{\beta}{4}t} dt + 2\sin(\alpha_{\lambda}/2) dB_{t}, \qquad \alpha_{\lambda}(0) = 0, \quad t \in [0, \infty)$$
 (3)

where $B_t = B_t^{(\lambda)} = \int_0^t \text{Re} \left[-ie^{-\frac{1}{2}i\alpha_{\lambda}(s)}d(B_1 + idB_2) \right]$ is a standard Brownian motion which depends on λ . Thus, if we are interested in the number of points in a given interval $[0, \lambda]$ then it is enough to study the SDE (3) instead of the system (2).

Using the SDE characterization of the $\operatorname{Sine}_{\beta}$ process one can show that it is translation invariant with density $(2\pi)^{-1}$ (see [23]). In particular, in a large interval $[0, \lambda]$ one expects roughly $(2\pi)^{-1}\lambda$ points. In [19] the authors refined this by showing that $N_{\beta}(\lambda)$ satisfies a central limit theorem, it is asymptotically normal with mean $\frac{\lambda}{2\pi}$ and variance $\frac{2}{\beta\pi^2}\log\lambda$.



The main goal of the current paper is to characterize the large deviation behavior of $N_{\beta}(\lambda)$. We will find the asymptotic probability of seeing an average density different from $(2\pi)^{-1}$ on a large interval. Our main theorem will show that $\lambda^{-1}N_{\beta}(\lambda)$ satisfies a large deviation principle with a good rate function.

Before stating the exact form of the theorem we need to introduce a couple of notations. We will use

$$K(a) = \int_0^{\pi/2} \frac{dx}{\sqrt{1 - a\sin^2 x}}, \quad E(a) = \int_0^{\pi/2} \sqrt{1 - a\sin^2 x} dx, \tag{4}$$

for the complete elliptic integrals of the first and second kind, respectively. Note that there are several conventions denoting these functions, we use the one in [1]. We also introduce the following function for a < 1:

$$\mathcal{H}(a) = (1 - a)K(a) - E(a). \tag{5}$$

Now we are ready to state our main theorem.

Theorem 1 Fix $\beta > 0$. The sequence of random variables $\frac{1}{\lambda}N_{\beta}(\lambda)$ satisfies a large deviation principle with scale λ^2 and good rate function $\beta I_{Sine}(\rho)$ with

$$I_{Sine}(\rho) = \frac{1}{8} \left[\frac{\nu}{8} + \rho \mathcal{H}(\nu) \right], \qquad \nu = \gamma^{(-1)}(\rho), \tag{6}$$

where $\gamma^{(-1)}$ denotes the inverse of the continuous, strictly decreasing function given by

$$\gamma(\nu) = \begin{cases} \frac{\mathcal{H}(\nu)}{8} \int_{-\infty}^{\nu} \mathcal{H}^{-2}(x) dx, & \text{if } \nu < 0, \\ \frac{1}{2\pi}, & \text{if } \nu = 0, \\ \frac{\mathcal{H}(\nu)}{8} \int_{1}^{\nu} \mathcal{H}^{-2}(x) dx, & \text{if } 0 < \nu < 1, \\ 0, & \text{if } \nu = 1. \end{cases}$$
 (7)

Roughly speaking, this means that the probability of seeing close to $\rho\lambda$ points in $[0, \lambda]$ for a large λ is asymptotically $e^{-\lambda^2 \beta I_{\text{Sine}}(\rho)}$. The precise statement is that if G is an open, and F is a closed subset of $[0, \infty)$ then

$$\lim_{\lambda \to \infty} \inf \frac{1}{\lambda^2} P(\frac{1}{\lambda} N_{\beta}(\lambda) \in G) \ge -\inf_{x \in G} \beta I_{\text{Sine}}(x),$$

$$\lim_{\lambda \to \infty} \sup \frac{1}{\lambda^2} P(\frac{1}{\lambda} N_{\beta}(\lambda) \in F) \le -\inf_{x \in F} \beta I_{\text{Sine}}(x).$$

The function γ may also be defined as the solution to the equation $4x(1-x)\gamma''(x) = \gamma(x)$ on the intervals $(-\infty, 0]$ and [0, 1] with boundary conditions $\lim_{x\to 0^{\pm}} \gamma(x) = \frac{1}{2\pi}$, $\gamma(1) = 0$ and $\lim_{x\to -\infty} \frac{\gamma(x)}{\sqrt{|x|}} = \frac{1}{4}$. The rate function $I_{\text{Sine}}(\rho)$ is strictly convex



and non-negative with $I_{\rm Sine}(\frac{1}{2\pi})=0$ and $I_{\rm Sine}(0)=\frac{1}{64}$. The function $I_{\rm Sine}(\frac{1}{2\pi}+x)$ behaves like $-\frac{\pi^2 x^2}{4\log(1/|x|)}$ for small |x|, and $I_{\rm Sine}(\rho)$ grows like $\frac{1}{2}\rho^2\log\rho$ as $\rho\to\infty$. These statements will be proved in Proposition 20.

We note that the behavior of $I_{\rm Sine}(\rho)$ near $\rho=\frac{1}{2\pi}$ is formally consistent with the already mentioned central limit theorem of $N_{\beta}(\lambda)$. For $\rho=\frac{1}{2\pi}+x$ with a small, but fixed |x| the probability of seeing close to $\frac{1}{2\pi}\lambda+x\lambda$ points in $[0,\lambda]$ is approximately $\exp\left(-\frac{\beta\pi^2\lambda^2x^2}{4\log(1/|x|)}\right)$. Now let us assume, that this is true even if x decays with λ , even though this regime is not covered in our theorem. If we substitute $\lambda x=\sqrt{\frac{2}{\beta\pi^2}\log\lambda}\cdot y$ (with a fixed y), then this probability would asymptotically equal to $e^{-y^2/2}$. This is in agreement with the fact that $N_{\beta}(\lambda)$ is asymptotically normal with mean $\frac{1}{2\pi}\lambda$ and variance $\frac{2}{8\pi^2}\log\lambda$.

Before moving on, a couple of historical notes are in order. In [23] the authors also show another large deviation statement for the $\operatorname{Sine}_{\beta}$ process regarding large intervals, namely that the asymptotic probability of not seeing any points in $[0, \lambda]$ is approximately $e^{-\frac{\beta}{64}\lambda^2}$. In [24] this result was sharpened by providing the more precise asymptotics of

$$P(N_{\beta}(\lambda) = 0) = (\kappa_{\beta} + o(1))\lambda^{\upsilon_{\beta}} \exp\left\{-\frac{\beta}{64}\lambda^{2} + \left(\frac{\beta}{8} - \frac{1}{4}\right)\lambda\right\}, \text{ as } \lambda \to \infty$$
 (8)

with $\upsilon_{\beta} = \frac{1}{4} \left(\frac{\beta}{2} - \frac{2}{\beta} - 3 \right)$ and a positive constant κ_{β} whose value was not determined. Similar results have been proven before for the classical cases $\beta = 1, 2, 4$, see e.g. [3, 8,22,25]. Moreover, the value of κ_{β} and higher order asymptotics were also established for these specific cases by [7,11,18]. Further extension in the classical cases include the exact asymptotics of $P(N_{\beta}(\lambda) = n)$ for fixed n and also for $n = o(\lambda)$. (See [22] and [13] for details.) In all of these results the main term of the asymptotic probability is $e^{-\frac{\beta}{64}\lambda^2}$. This is consistent with our result, as Theorem 1 and $I_{\text{Sine}}(0) = \frac{1}{64}$ implies

$$\lim_{\varepsilon \to 0} \lim_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(N_{\beta}(\lambda) \le \varepsilon \lambda) = -\frac{\beta}{64}.$$

The large deviation rate function (6) has been predicted using non-rigorous scaling and log-gas arguments in [10] and [12]. (See Section 14.6 of [13] for an overview.) Using the same techniques [14] treats the corresponding problem for the soft edge and hard edge limit processes of β -ensembles.

One can also study the large deviation behavior of the empirical distribution of the β -ensembles on a macroscopic level. It is known that after scaling with \sqrt{n} the empirical measure of the distribution (1) converges to the Wigner semicircle law. In [4] the authors prove a large deviation principle for the scaled empirical measure, this describes the asymptotic probability of seeing a different density profile than the semicircle. One could consider our theorem a microscopic analogue of that result.

We will study the asymptotic behavior of the diffusions (2) and (3) by comparing them to similar diffusions with piecewise constant drifts. This connects us to another,



related symmetric random matrix ensemble. Let $H_{n,\sigma}$ be a random symmetric tridiagonal matrix with entries equal to 1 above and below the diagonal and i.i.d. normals with mean zero and variance $\frac{\sigma^2}{n}$ on the diagonal.

$$H_{n,\sigma} = \begin{pmatrix} \omega_1 & 1 & & & \\ 1 & \omega_2 & 1 & & & \\ & 1 & \ddots & & \\ & & & \ddots & 1 \\ & & & 1 & \omega_n \end{pmatrix}, \quad \omega_i \sim N(0, \sigma^2 n^{-1}). \tag{9}$$

The matrix $H_{n,\sigma}$ can be viewed as a one-dimensional discrete random Schrödinger operator. In [19] it was shown that the bulk scaling limit of the spectrum of $H_{n,\sigma}$ (along appropriate subsequences) is a point process with density $(2\pi)^{-1}$ denoted by Sch_{τ} . (The parameter $\tau > 0$ depends on σ and the point in the spectrum where we zoom in to take the limit.) The process Sch_{τ} can be characterized via its counting function in a similar way to the Sine_{β} process. Consider the following one-parameter family of SDEs:

$$d\phi_{\lambda} = \lambda dt + dB_0 + \text{Re}\left[e^{-i\phi\lambda}(dB_1 + idB_2)\right], \quad \phi_{\lambda}(0) = 0, \quad t \in [0, \infty) \quad (10)$$

where B_0 , B_1 , B_2 are independent standard Brownian motions. Then the random set

$$\Lambda_{\tau} := \{\lambda : \phi_{\lambda/\tau}(\tau) \in 2\pi \mathbb{Z}\}$$

has the same distribution as $\operatorname{Sch}_{\tau}$. Denote the counting function of the process by \widetilde{N}_{τ} , i.e. for $\lambda>0$ let $\widetilde{N}_{\tau}(\lambda)=\#(\operatorname{Sch}_{\tau}\cap[0,\lambda])$. In [19] it was shown that $\widetilde{N}_{\tau}(\lambda)$ is close to a normal with mean $\frac{\lambda}{2\pi}$ and a constant variance $\frac{\tau}{4\pi^2}$. In our next result we derive the large deviation behavior of $\widetilde{N}_{\tau}(\lambda)$, this is the analogue of Theorem 1 for the $\operatorname{Sch}_{\tau}$ processes.

Theorem 2 Fix $\tau > 0$. The sequence of random variables $\frac{1}{\lambda} \widetilde{N}_{\tau}(\lambda)$ satisfies a large deviation principle with scale λ^2 and rate function $\frac{1}{\tau} I_{Sch}(\cdot)$ where $I_{Sch}(\rho) = \mathcal{I}(2\pi\rho)$ and for q > 0

$$\mathcal{I}(q) = \frac{2-a}{8} - \frac{E(a)}{4K(a)}, \text{ with } a = a(q) = K^{-1}(\pi/(2q))$$
 (11)

and $\mathcal{I}(0) = 1/8$.

The rate function $I_{\text{Sch}}(\rho)$ is strictly convex and locally quadratic at the absolute minimum point $\rho=\frac{1}{2\pi}$. (See Proposition 16.) The local behavior of $I_{\text{Sch}}(\rho)$ at $\rho=\frac{1}{2\pi}$ is formally consistent with the fact that $N_{\tau}(\lambda)-\frac{\lambda}{2\pi}$ is close to a normal random variable with a constant variance $\frac{\tau}{4\pi^2}$.

The proofs of Theorems 1 and 2 will rely on path level large deviation principles on the corresponding stochastic differential equations. These in turn will follow by



analyzing the hitting time of 2π for the diffusion

$$d\widetilde{\alpha}_{\lambda} = \lambda dt + 2\sin(\widetilde{\alpha}_{\lambda}/2)dB, \qquad \widetilde{\alpha}_{\lambda}(0) = 0, \quad t \in [0, \infty)$$
 (12)

for a fixed large λ . Note, that for a fixed λ the process $\widetilde{\alpha}_{\lambda}(t)$ is equal in distribution to $\phi_{\lambda}(t) - \phi_{0}(t)$ from (10).

In the next section we summarize some of the important properties of the SDEs we work with, and state the needed path level large deviation results. In Sect. 3 we study the diffusion $\widetilde{\alpha}_{\lambda}$ from (12) using the Cameron–Martin–Girsanov change of measure technique. In Sect. 4 we will derive a path level large deviation principle for the diffusion $\widetilde{\alpha}_{\lambda}$ from (12). Using this result and comparing α_{λ} from (3) to a similar diffusion with piecewise constant drift allows us to derive a path level large deviation principle for α_{λ} . This is done in Sect. 5. In Sect. 6 we analyze the rate functions for the path level large deviations and in Sect. 7 we complete the proofs of Theorems 1 and 2. In the "Appendix" we will discuss various properties and asymptotics of the used special functions.

2 Properties of the diffusions corresponding to Sine₈ and Sch₇

Our starting point is the observation that if $\lambda > 0$ is fixed, then if the diffusion $\widetilde{\alpha}_{\lambda}$ (defined in (12)) hits $2n\pi$ for $n \in \mathbb{Z}$, it will stay above it. This can be seen from the fact that when $\widetilde{\alpha}_{\lambda}$ hits $2n\pi$ the noise term vanishes, but the drift term is always positive. Introduce the notations

$$\lfloor y \rfloor_{2\pi} = \max\{2\pi k : 2\pi k \le y\}, \qquad \lceil y \rceil_{2\pi} = \min\{2\pi k : 2\pi k \ge y\}.$$

From the strong Markov property we immediately get the following proposition.

Proposition 3 Fix $\lambda > 0$. Then the process $\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi}$ is non-decreasing in t. Moreover, the waiting times between the jump times of this process are i.i.d. with the same distribution as the hitting time

$$\tau_{\lambda} = \inf\{t : \widetilde{\alpha}_{\lambda}(t) > 2\pi\}. \tag{13}$$

Consider the diffusions $\widetilde{\alpha}_{\lambda}^{(1)}$ and $\widetilde{\alpha}_{\lambda}^{(2)}$ which are strong solutions of the SDE (12), but with initial conditions $\widetilde{\alpha}_{\lambda}^{(1)}(0) = c_1 \leq \widetilde{\alpha}_{\lambda}^{(2)} = c_2$. Then a simple coupling argument shows that $\widetilde{\alpha}_{\lambda}^{(1)}(t) \leq \widetilde{\alpha}_{\lambda}^{(2)}(t)$ for all $t \geq 0$. Our next proposition will build on this statement using the strong Markov property.

Proposition 4 Let $0 = t_0 < t_1 < t_2 < \cdots < t_n = T$ and fix a $\lambda > 0$. Consider the solution $\widetilde{\alpha}_{\lambda}(t)$ of (12) on [0, T]. Then there exists independent random variables $\xi_1, \xi_2, \ldots, \xi_n$ so that

$$|\xi_i|_{2\pi} < |\widetilde{\alpha}_{\lambda}(t_i)|_{2\pi} - |\widetilde{\alpha}_{\lambda}(t_{i-1})|_{2\pi} < |\xi_i|_{2\pi} + 2\pi, \quad 1 < i < n,$$
 (14)

and ξ_i is distributed as $\widetilde{\alpha}_{\lambda}(t_i - t_{i-1})$.



Proof Let $\hat{\alpha}_i(s)$ be defined as the strong solution of (12) on $[t_{i-1}, t_i]$ with initial condition $\hat{\alpha}_i(t_{i-1}) = 0$ and let $\xi_i = \hat{\alpha}_i(t_i)$. Clearly, $\xi_i, 1 \le i \le n$ are independent random variables and $\xi_i \stackrel{d}{=} \widetilde{\alpha}_{\lambda}(t_i - t_{i-1})$, we just have to show that (14) holds. Fix an integer $1 \le i \le n$ and for $s \in [t_{i-1}, t_i]$ define

$$\widetilde{\alpha}_{\lambda}^{(1)}(s) = \hat{\alpha}_i(s) + \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi}, \qquad \widetilde{\alpha}_{\lambda}^{(2)}(s) = \hat{\alpha}_i(s) + \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi} + 2\pi.$$

Then $\widetilde{\alpha}_{\lambda}$, $\widetilde{\alpha}_{\lambda}^{(1)}$, $\widetilde{\alpha}_{\lambda}^{(2)}$ are all strong solutions of (12) on $[t_{i-1}, t_i]$ with initial conditions

$$\widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) \leq \widetilde{\alpha}_{\lambda}(t_{i-1}) \leq \widetilde{\alpha}_{\lambda}^{(2)}(t_{i-1}) = \widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) + 2\pi.$$

The ordering is preserved by the coupling so we have

$$\widetilde{\alpha}_{\lambda}^{(1)}(t_i) \leq \widetilde{\alpha}_{\lambda}(t_i) \leq \widetilde{\alpha}_{\lambda}^{(2)}(t_i) = \widetilde{\alpha}_{\lambda}^{(1)}(t_i) + 2\pi.$$

(See Fig. 1 for an illustration.) From this we get

$$\lfloor \widetilde{\alpha}_{\lambda}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi} = \lfloor \widetilde{\alpha}_{\lambda}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) \rfloor_{2\pi}$$

$$\geq \lfloor \widetilde{\alpha}_{\lambda}^{(1)}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) \rfloor_{2\pi} = \lfloor \xi_{i} \rfloor_{2\pi},$$

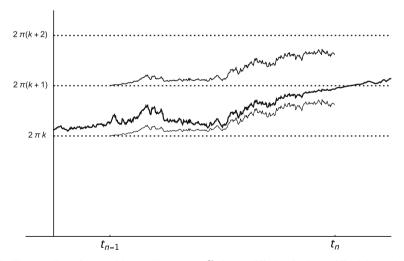


Fig. 1 The coupling of Proposition 4. The process $\widetilde{\alpha}_{\lambda}$ is the diffusion in the *middle*, it is sandwiched between $\widetilde{\alpha}_{\lambda}^{(1)}$ and $\widetilde{\alpha}_{\lambda}^{(2)} = \widetilde{\alpha}_{\lambda}^{(1)} + 2\pi$ which start at integer multiples of 2π at the beginning of the coupling interval



and

$$\lfloor \widetilde{\alpha}_{\lambda}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi} = \lfloor \widetilde{\alpha}_{\lambda}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) \rfloor_{2\pi}$$

$$\leq \lfloor \widetilde{\alpha}_{\lambda}^{(2)}(t_{i}) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}_{\lambda}^{(1)}(t_{i-1}) \rfloor_{2\pi} = \lfloor \xi_{i} \rfloor_{2\pi} + 2\pi.$$

We will also need another type of coupling for a slightly more general family of diffusions. Consider the SDE

$$d\xi_{f,c} = fdt + \text{Re}((e^{-i\xi_{f,c}} - 1)(dB_1 + idB_2)), \quad \xi_{f,c}(0) = c, \quad t \in [0, \infty)$$
 (15)

where f is an integrable non-negative function. Note, that for fixed f, c this process has the same distribution as

$$d\widetilde{\xi}_{f,c} = f dt + 2\sin(\widetilde{\xi}_{f,c}/2)dB, \quad \widetilde{\xi}_{f,c}(0) = c, \quad t \in [0, \infty)$$
 (16)

The following properties of $\xi_{f,c}$ follow from the basic theory of diffusions and standard coupling arguments.

- **Proposition 5** (i) Let $\tau_{2\pi n}$ be the hitting time of $2\pi n$, where $2\pi n > c$ and n is an integer. Then for any $t > \tau_{2\pi n}$ we have $\xi_{f,c} \geq 2\pi n$. In particular, if $c \geq 0$ then $\xi_{f,c}(t)$ stays non-negative for all t > 0.
- (ii) If $f \ge g$ and $\xi_{f,a}$ and $\xi_{g,b}$ are driven by the same Brownian motions then $\xi_{f,a} \xi_{g,b}$ has the same distribution as $\xi_{f-g,a-b}$. If $a \ge b$ then $\xi_{f,a} \xi_{g,b}$ stays a.s. nonnegative for all t.
- (iii) For any finite T we have the following exponential tail bound

$$P(\xi_{f,0}(T) \ge ka) \le 2\left(\frac{\int_0^T f(t)dt}{2\pi a}\right)^k, \quad k \in \mathbb{N}.$$
 (17)

If $\int_0^\infty f(t)dt < \infty$ then $\xi_{f,c}(\infty) = \lim_{t \to \infty} \xi_{f,c}(t)$ exists a.s. and the previous bound holds for $T = \infty$ as well.

Sketch of the proof The first statement follows from the strong Markov property and the fact that in (16) the noise term vanishes if $\tilde{\xi}_{f,c} \in 2\pi\mathbb{Z}$, but the drift is always non-negative. The first part of (ii) follows by considering the difference of the SDEs for $\xi_{f,a}, \xi_{g,b}$ and noting that $(e^{-i\xi_{f,a}} - e^{-i\xi_{g,b}})(dB_1 + idB_2)$ has the same distribution as $(e^{-i(\xi_{f,a} - \xi_{g,b})} - 1)(dB_1 + idB_2)$. The second part of (ii) follows from the first statement.

Finally, (17) follows from the Markov inequality and the strong Markov property. The existence of the limit is proved in Proposition 9 of [23]. (See that proposition for more details on the proof.)

Our main theorems will follow from the following path level large deviations.



Theorem 6 Fix $\beta > 0$ and let $\alpha_{\lambda}(t)$ be the process defined in (2) or (3). Then the sequence of rescaled processes $(\frac{\alpha_{\lambda}(t)}{\lambda}, t \in [0, \infty))$ satisfies a large deviation principle on $C[0, \infty)$ with the uniform topology with scale λ^2 and good rate function $\mathcal{J}_{Sine_{\beta}}$. The rate function $\mathcal{J}_{Sine_{\beta}}$ is defined as

$$\mathcal{J}_{\mathrm{Sine}_{\beta}}(g) = \int_{0}^{\infty} f^{2}(t) \mathcal{I}\left(g'(t)/\mathfrak{f}(t)\right) dt$$
, with $\mathfrak{f}(t) = \mathfrak{f}_{\beta}(t) = \frac{\beta}{4}e^{-\frac{\beta}{4}t}$

in the case where g(0) = 0 and g is absolutely continuous with non-negative derivative g'. In all other cases $\mathcal{J}_{Sine_{\beta}}(g)$ is defined as ∞ .

Theorem 7 Fix T > 0 and let $\widetilde{\alpha}_{\lambda}(t)$ be the process defined in (12). Then the sequence of rescaled processes $(\frac{\widetilde{\alpha}_{\lambda}(t)}{\lambda}, t \in [0, T])$ satisfies a large deviation principle on C[0, T] with the uniform topology with scale λ^2 and good rate function $\mathcal{J}_{Sch,T}$. The rate function is defined as

$$\mathcal{J}_{\mathrm{Sch},T}(g) = \int_0^T \mathcal{I}\left(g'(t)\right) dt$$

in the case where g(0) = 0 and g is absolutely continuous with non-negative derivative g', and $\mathcal{J}_{Sch,T}(g) = \infty$ in all other cases.

In order to prove Theorem 7 we observe that $\frac{\widetilde{\alpha}_{\lambda}(t)}{\lambda}$ is close to $\frac{\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi}}{\lambda}$ for large λ and by Proposition 3 we only need to analyze the hitting time τ_{λ} to understand the evolution of $\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi}$. The proof of Theorem 6 will follow along similar lines after approximating the drift in (2) with a piecewise constant function.

3 Analysis of the hitting time τ_{λ}

The following proposition summarizes our bounds on the relevant hitting times.

Proposition 8 Let $\tau_{\lambda} = \inf\{t : \widetilde{\alpha}_{\lambda}(t) \geq 2\pi\}$ where $\widetilde{\alpha}_{\lambda}$ is the solution of (12) and fix a < 1. Then we have

$$Ee^{\frac{\lambda^2 a}{8}\tau_{\lambda} - \frac{\lambda(|a| \wedge \sqrt{|a|})}{4}\tau_{\lambda}} \le e^{-\lambda \mathcal{H}(a)}.$$
 (18)

Let $t_a = 4K(a)$ and fix $0 < \varepsilon < |t_a - 2\pi|$. Then we have

$$P(\lambda \tau_{\lambda} \in [t_a - \varepsilon, t_a + \varepsilon]) \ge A(\varepsilon, \lambda, a) e^{-\lambda (\mathcal{H}(a) + \frac{at_a}{8}) - \lambda \frac{|a|}{8} - \lambda \frac{|a|}{2} (t_a + \varepsilon)}$$
(19)

where $\lim_{\lambda\to\infty} A(\varepsilon,\lambda,a) = 1$ for fixed a,ε .

Our first step is a change of variables in (12). We introduce $X_{\lambda}(t) = \log(\tan(\widetilde{\alpha}_{\lambda}(t)/4))$, by Itô's formula this satisfies the SDE

$$dX_{\lambda} = \frac{\lambda}{2} \cosh X_{\lambda} dt + \frac{1}{2} \tanh X_{\lambda} dt + dB_{t}, \qquad X_{\lambda}(0) = -\infty.$$
 (20)



The distribution of the hitting time of 2π for $\widetilde{\alpha}_{\lambda}(t)$ is the same as that of the hitting time of ∞ for X_{λ} . With a small abuse of notation from now on we will use the notation τ_{λ} for the blow-up time of $X_{\lambda}(t)$, i.e. $\tau_{\lambda} = \sup\{t : X_{\lambda}(t) < \infty\}$. In order to study τ_{λ} we will introduce a similar diffusion with a modified drift. Let a < 1 and consider

$$dY_{\lambda,a} = \frac{\lambda}{2} \sqrt{\cosh^2 Y_{\lambda,a} - a} \, dt + \frac{1}{2} \tanh Y_{\lambda,a} dt + dB_t, \qquad Y_{\lambda,a}(0) = -\infty. \tag{21}$$

To prove Proposition 8 we will choose an appropriate a and compare X_{λ} with the diffusion $Y_{\lambda,a}$ using the Cameron–Martin–Girsanov formula. Introduce the following notations for the drifts:

$$f_{\lambda}(x) = \frac{\lambda}{2} \cosh x + \frac{1}{2} \tanh x, \qquad h_{\lambda,a}(y) = \frac{\lambda}{2} \sqrt{\cosh^2 y - a} + \frac{1}{2} \tanh y.$$

Note, that we have the uniform bound

$$\left| f_{\lambda}(x) - h_{\lambda,a}(x) \right| \le \frac{1}{2} \lambda |a|. \tag{22}$$

The following proposition will be our main tool for our estimates.

Proposition 9 Fix a < 1 and consider $X = X_{\lambda}$ and $Y = Y_{\lambda,a}$. Denote by τ_{λ} and $\tau_{Y,\lambda}$ the blowup times of X and Y to ∞ . Then for any s > 0, we have

$$P(\lambda \tau_{\lambda} > s) = E\left[\mathbf{1}(\lambda \tau_{Y,\lambda} > s)e^{-G_{s/\lambda}(Y)}\right],\tag{23}$$

and

$$1 = Ee^{-G_{\tau \wedge (s/\lambda)}(Y)} = Ee^{G_{\tau \wedge (s/\lambda)}(X)}, \tag{24}$$

where

$$G_s(X) = \int_0^s h_{\lambda,a}(X(t)) - f_{\lambda}(X(t))dX - \frac{1}{2} \int_0^s (h_{\lambda,a}^2(X) - f_{\lambda}^2(X))dt.$$

Proof This is just the Cameron–Martin–Girsanov formula for diffusions with explosion. Note, that because of (22) the process $e^{G_{\tau \wedge s}(X)}$ satisfies the Novikov criterion and it is a positive martingale. From this the usual steps of the proof can be completed (see e.g. [16,17]).

Proof of Proposition 8 We first estimate the Girsanov exponent

$$G_s(X) = \frac{\lambda}{2} \int_0^s (\sqrt{\cosh^2 X - a} - \cosh X) dX$$
$$-\frac{1}{2} \int_0^s \left(-\frac{\lambda^2}{4} a + \frac{\lambda}{2} (\sqrt{\cosh^2 X - a} - \cosh X) \tanh X \right) dt.$$



Applying Itô's formula for $\theta(X) = h_{\lambda,a}(X) - f_{\lambda}(X)$ we have that $\int_0^t \theta(X) dX = \int_{X_0}^{X_t} \theta(x) dx - \frac{1}{2} \int_0^s \theta'(X) dt$. This gives us

$$G_s(X) = \frac{\lambda^2 a}{8} s + \frac{\lambda}{2} \int_{-\infty}^{X_s} \left(\sqrt{\cosh^2 x - a} - \cosh x \right) dx$$
$$+ \frac{\lambda}{4} \int_0^s \frac{a \tanh}{\sqrt{\cosh^2 X - a}} \cdot \frac{\sqrt{\cosh^2 X - a} - \cosh X}{\sqrt{\cosh^2 X - a} + \cosh X} ds.$$

Note, that

$$\frac{1}{2} \int_{\mathbb{R}} \left(\sqrt{\cosh^2 x - a} - \cosh x \right) dx = -\int_0^{\pi/2} \frac{a}{1 + \sqrt{1 - a \sin^2 y}} dy$$
$$= (1 - a)K(a) - E(a) = \mathcal{H}(a),$$

where this last equality can be seen by differentiating both sides with respect to a and checking equality at a = 0. It is not hard to check that

$$\left| \frac{a \tanh}{\sqrt{\cosh^2 X - a}} \frac{\sqrt{\cosh^2 X - a} - \cosh X}{\sqrt{\cosh^2 X - a} + \cosh X} \right| \le \left| \frac{a \tanh x}{\sqrt{\cosh^2 x - a}} \right| \le |a| \land \sqrt{|a|}, \quad \text{for } a < 1$$

uniformly in x. The upper bound |a| follows from $\sqrt{\cosh^2 x - a} \ge |\sinh x|$, while the bound $\sqrt{|a|}$ requires the optimization of the function $\frac{|a|\sqrt{y-1}}{\sqrt{y}\sqrt{y-a}}$ for $y \ge 1$. This gives the bound

$$\left| G_{\tau_{\lambda}}(X) - \frac{\lambda^2 a \tau}{8} - \lambda \mathcal{H}(a) \right| \le \frac{\lambda \tau(|a| \wedge \sqrt{|a|})}{4}. \tag{25}$$

To get the exponential moment bound (18) we use $1 = Ee^{G_{\tau \wedge s/\lambda}(X)}$ from (24). We let $s \to \infty$, use Fatou's lemma and (25) to get

$$1 \ge E e^{G_{\tau}(X)} \ge E e^{\frac{\lambda^2 a}{8}\tau + \lambda \mathcal{H}(a) - \frac{\lambda(|a| \wedge \sqrt{|a|})\tau}{4}}.$$
 (26)

Rearranging the terms we get (18).

To prove the lower bound (19) we write

$$P(\lambda \tau_{\lambda} \in (t_{a} - \varepsilon, t_{a} + \varepsilon)) = P(\lambda \tau_{Y,\lambda} > t_{a} - \varepsilon) - P(\lambda \tau_{Y,\lambda} > t_{a} + \varepsilon)$$

$$= E \left[\mathbf{1}(\lambda \tau_{Y,\lambda} > t_{a} - \varepsilon) e^{-G_{\tau \wedge (t_{a} - \varepsilon)/\lambda}(Y)} \right] - E \left[\mathbf{1}(\lambda \tau_{Y,\lambda} > t_{a} + \varepsilon) e^{-G_{\tau \wedge (t_{a} + \varepsilon)/\lambda}(Y)} \right]$$

$$= E \left[\mathbf{1}(\lambda \tau_{Y,\lambda} \in (t_{a} - \varepsilon, t_{a} + \varepsilon)) e^{-G_{\tau \wedge (t_{a} + \varepsilon)/\lambda}(Y)} \right]$$

$$= E \left[\mathbf{1}(\lambda \tau_{Y,\lambda} \in (t_{a} - \varepsilon, t_{a} + \varepsilon)) e^{-G_{\tau}(Y)} \right], \tag{27}$$



where we used the fact that $e^{-G_{\tau \wedge t}(Y)}$ is martingale in the third line. Because of (25) we have

$$G_{\tau}(Y) \le \frac{\lambda^2 a \tau}{8} + \lambda \mathcal{H}(a) + \frac{\lambda |a| \tau}{4},$$
 (28)

and we can bound the last expectation as

$$E\left[\mathbf{1}(\lambda\tau_{Y,\lambda}\in(t_{a}-\varepsilon,t_{a}+\varepsilon))e^{-G_{\tau}(Y)}\right]$$

$$\geq E\left[\mathbf{1}(\lambda\tau_{Y,\lambda}\in(t_{a}-\varepsilon,t_{a}+\varepsilon))e^{-\frac{\lambda^{2}a\tau}{8}-\lambda\mathcal{H}(a)-\frac{\lambda|a|\tau}{4}}\right]$$

$$> P(\lambda\tau_{Y,\lambda}\in(t_{a}-\varepsilon,t_{a}+\varepsilon))e^{-\frac{\lambda a(t_{a}\pm\varepsilon)}{8}-\lambda\mathcal{H}(a)-\frac{\lambda|a|(t_{a}+\varepsilon)}{4}}$$

where we choose the sign of ε in $t_a \pm \varepsilon$ the same way as the sign of a.

If we can show that $\lim_{\lambda\to\infty} P(\lambda \tau_{Y,\lambda} \in (t_a - \varepsilon, t_a + \varepsilon)) = 1$ for fixed a and ε then this will complete the proof of (19). Note, that $\widetilde{Y}(t) := Y_{\lambda,a}(t/\lambda)$ satisfies the SDE

$$d\widetilde{Y} = \frac{1}{2}\sqrt{\cosh^2 Y - a}dt + \frac{1}{2\lambda}\tanh\widetilde{Y}dt + \frac{1}{\sqrt{\lambda}}dB_t, \qquad \widetilde{Y}(0) = -\infty.$$

As $\lambda \to \infty$, the strong solution of this SDE converges a.s. to the solution of the ODE

$$y' = \frac{1}{2}\sqrt{\cosh^2 y - a}, \quad y(0) = -\infty.$$

This ODE can be solved and the solutions satisfies $\int_{-\infty}^{y(t)} \frac{2}{\sqrt{\cosh^2 x - a}} dx = t$. This shows that y explodes exactly at

$$\int_{-\infty}^{\infty} \frac{2}{\sqrt{\cosh^2 x - a}} dx = 4K(a) = t_a.$$

This shows that $\lim_{\lambda\to\infty} P(\lambda \tau_{Y,\lambda} \in (t_a - \varepsilon, t_a + \varepsilon)) = 1$ for fixed a and ε and this completes the proof of the proposition.

We can use the tail estimates of τ_{λ} to estimate the tail probabilities of $\widetilde{\alpha}_{\lambda}(t)$ for a fixed t. Recall the definition of $\mathcal{I}(\cdot)$ from (11).

Lemma 10 There exist a constant c so that for $\lambda > 2$ we have

$$e^{-\lambda^2 t \mathcal{I}(q) + \lambda c(t+1)(\mathcal{I}(q) + 1)} \ge \begin{cases} P(\lceil \widetilde{\alpha}_{\lambda}(t) \rceil_{2\pi} \ge qt\lambda) & \text{if } q > 1, \\ P(\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi} \le qt\lambda) & \text{if } 0 < q < 1. \end{cases}$$
(29)

Moreover, there are absolute constants c_0 , c_1 so that if $qt\lambda$, q and $\lambda q \log q$ are all bigger than c_0 then

$$P(\lceil \widetilde{\alpha}_{\lambda}(t) \rceil_{2\pi} \ge qt\lambda) \le e^{-c_1 \lambda^2 t \, q^2 \log q}. \tag{30}$$



Proof Introduce the hitting times

$$\tau_{\lambda}^{(n)} = \inf\{t > 0 : \widetilde{\alpha}_{\lambda}(t) > 2n\pi\}. \tag{31}$$

Then by Proposition 3 the random variables $\tilde{\tau}^{(n)} = \tau^{(n)} - \tau^{(n-1)}$ are i.i.d. with the same distribution as τ_{λ} . Applying the exponential Markov inequality we get

$$P(\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi} \le qt\lambda) = P\left(\sum_{i=1}^{\lceil qt\lambda/(2\pi) \rceil} \widetilde{\tau}^{(i)} \ge t\right) \le \left(Ee^{A\tau_{\lambda}}\right)^{\lceil qt\lambda/(2\pi) \rceil} e^{-At} \quad (32)$$

with any A > 0. Suppose first that q < 1 which also implies $a = a(q) = K^{-1}(\pi/(2q)) \in (0, 1)$. By choosing

$$A = \frac{\lambda^2 a}{8} - \frac{\lambda |a|}{4} \tag{33}$$

we have A > 0 if $\lambda > 2$ and from (18) we have $Ee^{A\tau_{\lambda}} \le e^{-\lambda \mathcal{H}(a)}$. Together with (32) this gives

$$P(\widetilde{\alpha}_{\lambda}(t) \leq qt\lambda) \leq e^{-\lambda \mathcal{H}(a)\lceil qt\lambda/(2\pi)\rceil - (\frac{\lambda^{2}a}{8} + \frac{\lambda|a|}{4})t}$$

$$\leq e^{-\frac{qt\lambda^{2}}{2\pi}\mathcal{H}(a) - (\frac{\lambda^{2}a}{8} + \frac{\lambda|a|}{4})t + \lambda|\mathcal{H}(a)|} = e^{-\lambda^{2}t\mathcal{I}(q) + \lambda\left(\frac{|a|t}{4} + |\mathcal{H}(a)|\right)}$$
(34)

where we used the definitions (11) and (5).

For the q > 1 case we use the same steps. Here $a = K^{-1}(\pi/(2q)) < 0$ and A defined in (33) is negative which is exactly what we need for the exponential Markov inequality. Eventually we get

$$P(\lceil \widetilde{\alpha}_{\lambda}(t) \rceil_{2\pi} \geq q\lambda t) \leq e^{-\lambda \mathcal{H}(a) \lfloor qt\lambda/(2\pi) \rfloor - (\frac{\lambda^{2}a}{8} - \frac{\lambda |a|}{4})t}$$

$$\leq e^{-\frac{qt\lambda^{2}}{2\pi} \mathcal{H}(a) - (\frac{\lambda^{2}a}{8} - \frac{\lambda |a|}{4})t + \lambda |\mathcal{H}(a)|}$$

$$= e^{-\lambda^{2}t\mathcal{I}(q) + \lambda \left(\frac{|a|t}{4} + |\mathcal{H}(a)|\right)}.$$
(35)

By Lemma 18 in the "Appendix" there is a constant c so that

$$\mathcal{H}(a(q)) + \frac{1}{4}|a(q)|t \le c(t+1)(\mathcal{I}(q)+1),$$
 (36)

for all t, q > 0 which means that we can replace the upper bounds in (34) and (35) with $e^{-\lambda^2 t \mathcal{I}(q) + \lambda c(t+1)(\mathcal{I}(q)+1)}$. This proves the first part of Lemma 10.

For the second part we repeat the same steps as in the q > 1 case, but now use

$$A = \frac{\lambda^2 a}{8} - \frac{\lambda \sqrt{|a|}}{4}.$$



This gives

$$P(\lceil \widetilde{\alpha}_{\lambda}(t) \rceil_{2\pi} \geq q\lambda t) \leq e^{-\lambda \mathcal{H}(a) \lfloor qt\lambda/(2\pi) \rfloor - \frac{\lambda^2 a}{8} t + \frac{\lambda \sqrt{|a|}}{4} t}.$$

By Proposition 17 of the "Appendix" if q is large enough then $a = K^{-1}(\pi/(2q)) > cq^2 \log^2 q$ with some positive constant c. If $-a\lambda$ and $qt\lambda$ are big enough (which can be achieved by choosing c_0 big enough), we will have

$$\lfloor qt\lambda/(2\pi)\rfloor > \frac{9}{10}qt\lambda/(2\pi), \qquad -\frac{\lambda^2 a}{8} + \frac{\lambda\sqrt{|a|}}{4} < -\frac{11}{10} \cdot \frac{\lambda^2 a}{8}.$$

Then

$$\begin{split} -\lambda \mathcal{H}(a) \lfloor qt\lambda/(2\pi) \rfloor &- \frac{1}{8}\lambda^2 at + \frac{1}{4}\lambda \sqrt{|a|}t < -\lambda^2 t \left(\frac{9}{10} \mathcal{H}(a) \frac{q}{2\pi} + \frac{11}{10} \frac{a}{8} \right) \\ &= -\lambda^2 t \left(-\frac{7a}{80} - \frac{9E(a)}{40K(a)} + \frac{9}{40} \right) \\ &< -c_2 \lambda^2 t q^2 \log^2 q. \end{split}$$

with a positive constant c_2 , where in the last step we again used the asymptotics given in Proposition 17 together with (89). This completes the proof of (30).

4 The path deviation for the $\tilde{\alpha}_{\lambda}$ process

In this section we will prove Theorem 7. In order to show the large deviation principle we need that

$$\lim_{\lambda \to \infty} \inf \frac{1}{\lambda^2} P\left(\frac{\widetilde{\alpha}_{\lambda}(\cdot)}{\lambda} \in G\right) \ge \inf_{g \in G} \mathcal{J}_{\mathrm{Sch},T}(g), \quad \text{for any open set } G \subset C[0,T],$$

$$\lim_{\lambda \to \infty} \sup \frac{1}{\lambda^2} P\left(\frac{\widetilde{\alpha}_{\lambda}(\cdot)}{\lambda} \in K\right) \le \inf_{g \in K} \mathcal{J}_{\mathrm{Sch},T}(g), \quad \text{for any closed set } K \subset C[0,T].$$

The fact that $\mathcal{J}_{Sch,T}(g)$ is a good rate function will be proved in Proposition 14 of Sect. 6.

We will use the fact that $\mathcal{I}(x)$ is strictly convex on $(0, \infty)$ with a global minimum at $\mathcal{I}(1) = 0$, and also that there is a constant c > 0 so that

$$c^{-1} \le \frac{\mathcal{I}(x)}{x^2 \log^2 x} \le c, \quad \text{for all } x > 2.$$
 (37)

These statements will be proved in Propositions 16 and 17 of the "Appendix".

Proof of the large deviations upper bound in Theorem 7 We will follow the standard strategy for proving path level large deviations. Consider a closed subset K of C[0, T].



We need to bound $P(\frac{1}{\lambda}\widetilde{\alpha}_{\lambda}(\cdot) \in K)$. Define the δ -'fattening' of K as

$$K^{\delta} := \{ f \in C[0, T] : ||f - g|| \le \delta \text{ for some } g \in K \}.$$
 (38)

From now on $\|\cdot\|$ denotes the sup-norm on the appropriate interval.

Let π_N be the following projection of C[0, T] to piecewise linear paths:

$$(\pi_N f)(iT/N) = |f(iT/N)|_{2\pi}, \quad 0 \le i \le N$$
(39)

and $\pi_N f$ is defined linearly between these points. Then

$$P(\tilde{\alpha}_{\lambda}/\lambda \in K) \le P(\|\tilde{\alpha}_{\lambda} - \pi_N \tilde{\alpha}_{\lambda}\| \ge \delta \lambda) + P(\frac{1}{\lambda} \pi_N(\tilde{\alpha}_{\lambda}) \in K^{\delta}). \tag{40}$$

We will bound the two probabilities in (40) separately.

The first term can be rewritten as

$$P\left[\|\tilde{\alpha}_{\lambda} - \pi_{N}\tilde{\alpha}_{\lambda}\| \ge \delta\lambda\right] = P\left(\max_{t \in \left[\frac{(k-1)T}{N}, \frac{kT}{N}\right]} |\pi_{N}\tilde{\alpha}_{\lambda}(t) - \tilde{\alpha}_{\lambda}(t)| \ge \delta\lambda\right). \tag{41}$$

By Proposition 3 the process $\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi}$ is non-decreasing. Thus for any fixed k we have

$$\sup_{t \in [\frac{(k-1)T}{N}, \frac{kT}{N}]} |\pi_N \tilde{\alpha}_{\lambda}(t) - \tilde{\alpha}_{\lambda}(t)| \leq \lceil \tilde{\alpha}_{\lambda}((k+1)T/N) \rceil_{2\pi} - \lfloor \tilde{\alpha}_{\lambda}(kT/N) \rfloor_{2\pi}.$$

By Proposition 4 the term on the right is stochastically dominated by $\widetilde{\alpha}_{\lambda}(T/N) + 4\pi$ therefore

$$P(\|\widetilde{\alpha}_{\lambda} - \pi_{N}\widetilde{\alpha}_{\lambda}\| \ge \delta\lambda) \le NP(\widetilde{\alpha}_{\lambda}(T/N) + 4\pi \ge \delta\lambda) \le NP\left(\frac{1}{\lambda}\widetilde{\alpha}_{\lambda}(T/N) \ge \frac{\delta}{2}\right)$$
(42)

where the last bound holds if $\lambda > 8\pi/\delta$. Using Lemma 10 we get

$$NP\left(\frac{1}{\lambda}\widetilde{\alpha}_{\lambda}(T/N) \geq \frac{\delta}{2}\right) \leq Ne^{-(\lambda^2\frac{T}{N} + \lambda c_1(T/N + 1))\mathcal{I}\left(\frac{\delta N}{2T}\right) + \lambda c_1(T/N - 1)}$$

and this leads to

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\|\tilde{\alpha}_{\lambda} - \pi_N \tilde{\alpha}_{\lambda}\| \ge \delta \lambda) \le -\frac{T}{N} \mathcal{I}\left(\frac{\delta N}{2T}\right). \tag{43}$$

Note, that for fixed δ and T as $N \to \infty$ the right hand side converges to $-\infty$ by (37).



The second term on the right side of (40) can be bounded as

$$P\left(\pi_N(\widetilde{\alpha}_{\lambda}/\lambda) \in K^{\delta}\right) \leq P\left(\mathcal{J}_{\mathrm{Sch},T}(\pi_N(\widetilde{\alpha}_{\lambda}/\lambda)) \geq \inf_{g \in K^{\delta}} \mathcal{J}_{\mathrm{Sch},T}(g)\right).$$

We introduce

$$\Delta \widetilde{\alpha}_i = \frac{N}{\lambda T} \left(\lfloor \widetilde{\alpha}(iT/N) \rfloor_{2\pi} - \lfloor \widetilde{\alpha}((i-1)T/N) \rfloor_{2\pi} \right), \quad \text{for } 1 \le i \le N$$

and $C_{\delta} = \inf_{g \in K^{\delta}} \mathcal{J}_{Sch,T}(g)$. Then we have to bound

$$P(\mathcal{J}_{Sch,T}(\pi_N(\widetilde{\alpha}_{\lambda}/\lambda)) \ge C_{\delta}) = P\left(\sum_{i=1}^N \frac{T}{N} \mathcal{I}(\Delta \widetilde{\alpha}_i) \ge C_{\delta}\right). \tag{44}$$

We can apply Proposition 4 with $t_i = \frac{iT}{N}$, $1 \le i \le N$ to get independent random variables ξ_i with $\xi_i \stackrel{d}{=} \widetilde{\alpha}_{\lambda}(T/N)$ and

$$\frac{N}{\lambda T} \lfloor \xi_i \rfloor_{2\pi} \le \Delta \widetilde{\alpha}_i \le \frac{N}{\lambda T} \left(\lfloor \xi_i \rfloor_{2\pi} + 2\pi \right).$$

Because of the convexity of $\mathcal{I}(\cdot)$ we then have

$$\mathcal{I}\left(\Delta\widetilde{\alpha}_{i}\right) \leq \max\left(\mathcal{I}\left(\frac{N\lfloor\xi_{i}\rfloor_{2\pi}}{\lambda T}\right), \mathcal{I}\left(\frac{N\lfloor\xi_{i}\rfloor_{2\pi}}{\lambda T} + \frac{2\pi N}{\lambda T}\right)\right)$$
$$\leq \left(1 + \frac{2\pi N}{\lambda T}\right)\mathcal{I}\left(\frac{N\lfloor\xi_{i}\rfloor_{2\pi}}{\lambda T}\right) + c\frac{2\pi N}{\lambda T}$$

where we used Lemma 19 of the "Appendix" for the last bound. Fix $1/2 > \varepsilon > 0$. Using the exponential Markov inequality, the independence of ξ_i and $\xi_i \stackrel{d}{=} \widetilde{\alpha}_{\lambda}(T/N)$ we get the bound

$$P\left(\sum_{i=1}^{N} \frac{T}{N} \mathcal{I}\left(\Delta \widetilde{\alpha}_{i}\right) \geq C_{\delta}\right)$$

$$\leq \left(Ee^{(1-2\varepsilon)\lambda^{2} \frac{T}{N}\left(\left(1+\frac{2\pi N}{\lambda T}\right) \mathcal{I}\left(\frac{N\lfloor \widetilde{\alpha}_{\lambda}(T/N)\rfloor_{2\pi}}{\lambda T}\right) + c\frac{2\pi N}{\lambda T}\right)\right)^{N} e^{-(1-2\varepsilon)\lambda^{2} C_{\delta}}$$

$$\leq \left(Ee^{(1-\varepsilon)\lambda^{2} \frac{T}{N} \mathcal{I}\left(\frac{N\lfloor \widetilde{\alpha}_{\lambda}(T/N)\rfloor_{2\pi}}{\lambda T}\right)\right)^{N} e^{(1-2\varepsilon)c2\pi\lambda N - (1-2\varepsilon)\lambda^{2} C_{\delta}},$$
(45)

where the second inequality holds for fixed ε , N, T if λ is big enough. Our next step is to estimate the exponential moment $Ee^{(1-\varepsilon)\lambda^2\frac{T}{N}\mathcal{I}\left(\frac{N\lfloor\widetilde{\alpha}_{\lambda}(T/N)\rfloor_{2\pi}}{\lambda T}\right)}$ for a fixed $\varepsilon>0$. By Lemma 11 below if N, T, ε are fixed then



$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log E e^{(1-\varepsilon)\lambda^2 \frac{T}{N} \mathcal{I}\left(\frac{N\lfloor \widetilde{\alpha}_{\lambda}(T/N)\rfloor_{2\pi}}{\lambda T}\right)} \leq 0.$$

Using this with (45) we get

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\sum_{i=1}^N \frac{T}{N} \mathcal{I}\left(\Delta \widetilde{\alpha}_i\right) \ge C_\delta\right) \le -(1 - 2\varepsilon)C_\delta. \tag{46}$$

Now we let $\varepsilon \to 0$ and then $N \to \infty$. The bounds (43), (46) with (40) give

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\lambda^{-1} \widetilde{\alpha}_{\lambda}(\cdot) \in K) \le -\inf_{g \in K^{\delta}} \mathcal{J}_{\mathrm{Sch}, T}(g). \tag{47}$$

Using the fact that $\mathcal{J}_{\mathrm{Sch},T}$ is a good rate function (which is proved in Proposition 14 of Sect. 6) we get that the right hand side converges to $-\inf_{g \in K} \mathcal{J}_{\mathrm{Sch},T}(g)$ as $\delta \to 0$. (See e.g. Lemma 4.1.6 from [9].) This finishes the proof of the lower bound.

Now we will prove the missing estimate for the lower bound.

Lemma 11 Fix t > 0 and $1 > \varepsilon > 0$. Then

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log E e^{(1-\varepsilon)\lambda^2 t \mathcal{I}\left(\frac{|\tilde{\alpha}_{\lambda}(t)|_{2\pi}}{\lambda t}\right)} \le 0.$$

Proof Introduce the temporary notation $G(x) = \lambda^2 t \mathcal{I}(x)$. This is a convex function with G(1) = 0 as its minimum. Then we have

$$Ee^{(1-\varepsilon)\lambda^2t\mathcal{I}\left(\frac{|\widetilde{\alpha}_{\lambda}(t)|}{\lambda t}\right)\rfloor_{2\pi}} \leq 2 - \int_0^1 (1-\varepsilon)G'(x)e^{(1-\varepsilon)G(x)}P(\lfloor \widetilde{\alpha}_{\lambda}(t)\rfloor_{2\pi} < \lambda t x)dx + \int_1^\infty (1-\varepsilon)G'(x)e^{(1-\varepsilon)G(x)}P(\lfloor \widetilde{\alpha}_{\lambda}(t)\rfloor_{2\pi} > \lambda t x)dx.$$

Using Lemma 10 we get

$$P(\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi} < \lambda t x) \le \exp\left\{-(1 - c_1 \lambda^{-1}(1 + t^{-1}))G(x) + \lambda c_1(t+1)\right\}$$

for x < 1 and a similar bound for $P(\lfloor \widetilde{\alpha}_{\lambda}(t) \rfloor_{2\pi} > \lambda t x) \leq P(\lceil \widetilde{\alpha}_{\lambda}(t) \rceil_{2\pi} > \lambda t x)$ for x > 1. This gives us

$$\begin{split} Ee^{(1-\varepsilon)\lambda^2t\mathcal{I}\left(\frac{\tilde{\alpha}_{\lambda}(t)}{\lambda t}\right)} &\leq 2 - \int_0^1 (1-\varepsilon)G'(x)e^{((1+t^{-1})(c_1/\lambda) - \varepsilon)G(x) + \lambda c_1(t+1)} dx \\ &+ \int_1^\infty (1-\varepsilon)G'(x)e^{((1+t^{-1})(c_1/\lambda) - \varepsilon)G(x) + \lambda c_1(t+1)} dx \\ &< 2 + 4\varepsilon^{-1}e^{\lambda c_1(t+1)} \end{split}$$



where the last inequality holds if $(1 + t^{-1})c_1/\lambda < \varepsilon/2$, i.e. for large enough λ . From this the lemma follows.

Now we turn to the lower bound proof in the large deviation result of Theorem 7. As we will see, we will be able to reduce the problem to studying the probability of $\frac{1}{1}\widetilde{\alpha}_{\lambda}(t)$ being close to a straight line.

Proposition 12 Fix q > 0 and $T, \varepsilon > 0$. Then

$$\lim_{\varepsilon \to 0} \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\widetilde{\alpha}(t) \in [\lambda(qt - \varepsilon), \lambda(qt + \varepsilon)], t \in [0, T]) \ge -T\mathcal{I}(q).$$

Proof As q > 0 we may assume $\varepsilon \le qT/2$ by choosing ε small enough. Let N =Theory has $q > \epsilon$ the majoration $\varepsilon_1 = \frac{\pi \varepsilon}{2q(qT+\varepsilon)}$, which satisfies $\varepsilon_1 < \frac{\varepsilon \lambda}{2qN}$ for $\lambda > 2$. Recall the definition of $\tau_{\lambda}^{(n)}$ and $\widetilde{\tau}_{\lambda}^{(n)}$ from (31). We will prove that

$$\begin{split} P(\widetilde{\alpha}(t) &\in [\lambda(qt-\varepsilon), \lambda(qt+\varepsilon)], t \in [0,T]) \\ &\geq P\left(\lambda\widetilde{\tau}_{\lambda}^{(k)} \in \left(\frac{2\pi}{q} - \varepsilon_1, \frac{2\pi}{q} + \varepsilon_1\right), 1 \leq k \leq N\right). \end{split}$$

Roughly speaking, this will follow from the simple fact that if we are within ε/q of the line y = qt in the horizontal direction, then we are within ε in the vertical direction. If $\lambda \widetilde{\tau}_{\lambda}^{(k)} \geq \frac{2\pi}{a} - \varepsilon_1$ for $1 \leq k \leq N$ then $\lambda \tau_{\lambda}^{(k)} \geq k(\frac{2\pi}{a} - \varepsilon_1)$ and

$$\widetilde{\alpha}_{\lambda}\left(\frac{k}{\lambda}\left(\frac{2\pi}{q}-\varepsilon_{1}\right)\right) \leq 2k\pi = \lambda \frac{2\pi}{2\pi/q-\varepsilon_{1}} \cdot \frac{k}{\lambda}\left(\frac{2\pi}{q}-\varepsilon_{1}\right)$$

for $1 \le k \le N$. Together with the fact that $\lfloor \widetilde{\alpha}_{\lambda} \rfloor_{2\pi}$ is non-decreasing we get that

$$\widetilde{\alpha}_{\lambda}(t) \leq \lambda \frac{2\pi}{2\pi/q - \varepsilon_1} \cdot \left(t + \frac{1}{\lambda}(2\pi/q - \varepsilon_1)\right), \quad \text{for } t \leq \frac{N}{\lambda}\left(\frac{2\pi}{q} - \varepsilon_1\right).$$

This inequality implies $\widetilde{\alpha}_{\lambda}(t) \leq \lambda(qt+\varepsilon)$, for $t \leq T$, $\lambda \varepsilon > 4\pi$. The other direction is similar, if we have $\lambda \widetilde{\tau}_{\lambda}^{(k)} \leq \frac{2\pi}{q} + \varepsilon_1$ for $1 \leq k \leq N$ then

$$\widetilde{\alpha}_{\lambda}\left(\frac{k}{\lambda}\left(\frac{2\pi}{q} + \varepsilon_{1}\right)\right) \geq 2k\pi = \lambda \frac{2\pi}{2\pi/q + \varepsilon_{1}} \cdot \frac{k}{\lambda}\left(\frac{2\pi}{q} + \varepsilon_{1}\right)$$

which implies

$$\widetilde{\alpha}_{\lambda}(t) \ge \lambda \frac{2\pi}{2\pi/q + \varepsilon_1} \cdot \left(t - \frac{1}{\lambda}(2\pi/q + \varepsilon_1)\right), \quad \text{for } t \le \frac{N}{\lambda}\left(\frac{2\pi}{q} + \varepsilon_1\right).$$

and $\widetilde{\alpha}_{\lambda}(t) \geq \lambda(qt - \varepsilon)$ for $t \leq T$. Using the independence of $\widetilde{\tau}_{\lambda}^{(k)}$ we get the bound

$$P(\widetilde{\alpha}(t) \in [\lambda(qt - \varepsilon), \lambda(qt + \varepsilon)], t \in [0, T]) \ge P\left(\lambda \widetilde{\tau}_{\lambda}^{(k)} \in (\frac{2\pi}{q} - \varepsilon_1, \frac{2\pi}{q} + \varepsilon_1)\right)^{N}.$$
(48)



By the lower bound (19) we have

$$\begin{split} \log P(\widetilde{\alpha}(t) &\in [\lambda(qt-\varepsilon), \lambda(qt+\varepsilon)], t \in [0,T]) \\ &\geq \frac{\lambda(qT+2\varepsilon)}{2\pi} \left(-\frac{2\pi\lambda}{q} \mathcal{I}(q) - \frac{\lambda|a|}{8} (\varepsilon_1 + 4(2\pi/q + \varepsilon_1)) + \log A(\varepsilon_1,\lambda,a) \right). \end{split}$$

Recalling $\varepsilon_1 = \frac{\pi \varepsilon}{2q(qT+\varepsilon)}$ we get

$$\lim_{\varepsilon \to 0} \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\widetilde{\alpha}(t) \in [\lambda(qt - \varepsilon), \lambda(qt + \varepsilon)], t \in [0, T]) \ge -T\mathcal{I}(q).$$

Proof of the lower bound in Theorem 7 Let G be an open subset of C[0, T]. We would like to show that

$$\liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\frac{1}{\lambda} \widetilde{\alpha}_{\lambda}(\cdot) \in G\right) \ge -\inf_{g \in G} \int_0^T \mathcal{I}\left(g'(t)\right) dt. \tag{49}$$

For this it is enough to prove that for any $g \in G$ with $\int_0^T \mathcal{I}\left(g'(t)\right) dt < \infty$ and $\delta > 0$ we have

$$\liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\frac{1}{\lambda} \widetilde{\alpha}_{\lambda}(\cdot) \in G\right) \ge -\int_0^T \mathcal{I}\left(g'(t)\right) dt - \delta. \tag{50}$$

We can approximate g with a piecewise linear function \tilde{g} in the sup-norm so that we have $|\int_0^T \mathcal{I}\left(g_n'(t)\right)dt - \int_0^T \mathcal{I}\left(g'(t)\right)dt| < \delta$. Because of this we may assume that g is piecewise linear, moreover, we may assume that there are no horizontal segments in g. Suppose that g is linear with slope q_i on the interval $[T_i, T_{i+1}]$ with $0 \le i \le k-1$ and $0 = T_0 < T_1 < \cdots < T_k = T$. We claim that if $\lambda > \lambda_0(\varepsilon, k)$ then

$$P\left(\left\|\frac{1}{\lambda}\widetilde{\alpha}_{\lambda}(\cdot) - g(\cdot)\right\| \leq \varepsilon\right)$$

$$\geq P\left(\left|\frac{1}{\lambda}(\widetilde{\alpha}_{\lambda}(t) - \widetilde{\alpha}_{\lambda}(T_{i})) - q_{i}(t - T_{i})\right| \leq \varepsilon/k, \text{ if } t \in [T_{i}, T_{i+1}]\right)$$

$$\geq \prod_{i=0}^{k-1} P\left(\left|\frac{1}{\lambda}\widetilde{\alpha}_{\lambda}(t) - q_{i}t\right| \leq \varepsilon/(2k), \text{ for } t \in [0, T_{i+1} - T_{i}]\right)$$
(51)

The first inequality is straightforward, to prove the second we use the coupling in the proof of Proposition 4. Recall the definition of the processes $\hat{\alpha}_i(s)$ defined on $[t_{i-1}, t_i]$. These were independent for different values of i and the process $\hat{\alpha}_i(s+t_{i-1}), s \in [0, t_i-t_{i-1}]$ had the same distribution as $\widetilde{\alpha}_{\lambda}(s), s \in [0, t_i-t_{i-1}]$. We also had

$$\hat{\alpha}_i(s) + \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi} \leq \widetilde{\alpha}_{\lambda}(s) \leq \hat{\alpha}_i(s) + \lfloor \widetilde{\alpha}_{\lambda}(t_{i-1}) \rfloor_{2\pi} + 2\pi.$$



for $s \in [t_{i-1}, t_i]$. By choosing $\lambda > \lambda_0 = 4\pi k/\varepsilon$ the inequality (51) follows by the independent increment property of the Brownian motion.

By Proposition 12 we have the bound

$$\lim_{\varepsilon \to 0} \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P \left(\| \frac{1}{\lambda} \widetilde{\alpha}_{\lambda}(\cdot) - g(\cdot) \| \le \varepsilon \right)$$

$$\geq -\sum_{i=0}^{k-1} (T_{i+1} - T_i) \mathcal{I}(q_i) = -\int_0^T \mathcal{I}\left(g'(t)\right) dt$$

from which (50) and thus the proof of the lower bound follows.

5 The path deviation for the Sine₈ process

This section contains the proof of Theorem 6. The strategy for the proof is to approximate the SDE (3) with a version where the drift is piecewise constant and then use elements of the proof of Theorem 7. Just as in the proof of Theorem 7, we need to show an upper and a lower bound to prove the large deviation principle. The fact that $\mathcal{J}_{Sine_{\mathcal{B}}}$ is a good rate function will be proved in Proposition 14 of Sect. 6.

Proof of the upper bound in Theorem 6 For the proof of the upper bound we go through a series of approximations: we essentially cut of the tail of the process, then replace the drift in the SDE with a piecewise constant version and then approximate the process with a piecewise linear version. Recall that $\alpha_{\lambda}(t)$ solves the SDE (2) and that we introduced the notation $\mathfrak{f}(t) = \frac{\beta}{4}e^{\frac{\beta}{4}t}$. Fix T > 0, the value of which will go to infinity later. The first approximating process is defined as

$$\alpha_{\lambda}^{(1)}(t) = \alpha_{\lambda}(t)\mathbf{1}(t \le T) + (\alpha_{\lambda}(T) + \lambda(e^{-\frac{\beta}{4}T} - e^{-\frac{\beta}{4}t}))\mathbf{1}(t > T),$$

this solves the SDE (2) with the noise 'turned off' at t = T. For the second process we define

$$f_N(t) = f(Ti/N), \quad t \in [Ti/N, T(i+1)/N)$$
 (52)

and consider the solution $\xi_{\lambda f_N}$ of (15) with drift λf_N and initial condition 0. Let

$$\alpha_{\lambda}^{(2)}(t) = \xi_{\lambda f_N}(t) \mathbf{1}(t \le T) + (\xi_{\lambda f_N}(T) + \lambda (e^{-\frac{\beta}{4}T} - e^{-\frac{\beta}{4}t})) \mathbf{1}(t > T).$$

Finally, let π_{MN} is the projection defined in (39) with intervals of size T/MN, that is $\pi_{MN} f$ is the piecewise linear path that satisfies

$$(\pi_{MN} f)(Ti/(MN)) = |f(Ti/(MN))|_{2\pi},$$

and is linear between these values. Define

$$\alpha_{\lambda}^{(3)}(t) = \pi_{MN} \xi_{\lambda \mathfrak{f}_N}(t) \mathbf{1}(t \leq T) + (\pi_{MN} \xi_{\lambda \mathfrak{f}_N}(T) + \lambda (e^{-\frac{\beta}{4}T} - e^{-\frac{\beta}{4}t})) \mathbf{1}(t > T).$$



Then for any closed set $K \subset C[0, \infty)$ we have that

$$P\left(\frac{\alpha_{\lambda}}{\lambda} \in K\right) \leq P\left(\frac{\alpha_{\lambda}^{(3)}}{\lambda} \in K^{3\delta}\right) + P(\|\alpha_{\lambda}^{(1)} - \alpha_{\lambda}\|_{\infty} \geq \delta\lambda) + P(\|\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}\|_{\infty} \geq \delta\lambda) + P(\|\alpha_{\lambda}^{(3)} - \alpha_{\lambda}^{(2)}\|_{\infty} \geq \delta\lambda),$$

$$(53)$$

where $K^{3\delta}$ is defined similarly to (38), as the 3δ -fattening of K. We will begin with the main term. Let

$$\mathcal{J}_N(g) = \int_0^\infty \mathfrak{f}_N^2(t) \mathcal{I}\left(\frac{g'(t)}{\mathfrak{f}_N(t)}\right) dt,$$

and define (similarly to the $\tilde{\alpha}_{\lambda}$ case in the proof of Theorem 7)

$$\Delta \alpha_i = \frac{MN}{\lambda f_N(\frac{Ti}{MN})T} \left(\lfloor \alpha^{(3)}(Ti/(MN)) \rfloor_{2\pi} - \lfloor \alpha^{(3)}(T(i-1)/(MN)) \rfloor_{2\pi} \right),$$

for $1 \le i \le MN$. Then,

$$P\left(\frac{\alpha_{\lambda}^{(3)}}{\lambda} \in K^{3\delta}\right) \leq P\left(\mathcal{J}_{N}\left(\frac{\alpha_{\lambda}^{(3)}}{\lambda}\right) \geq \inf_{g \in K^{3\delta}} \mathcal{J}_{N}(g)\right)$$

$$= P\left(\sum_{i=1}^{MN} \frac{T(\mathfrak{f}_{N}(Ti/(MN)))^{2}}{MN} \mathcal{I}(\Delta\alpha_{i}) \geq \inf_{g \in K^{3\delta}} \mathcal{J}_{N}(g)\right).$$

Take $\widehat{\alpha}_i$ to solve (12) but with the Brownian motion B(t+Ti/(MN)) - B(Ti/(MN)) and $\lambda_i = \lambda f_N(Ti/(MN))$. Then using the same arguments as in the bound (45) we get

$$P\left(\frac{\alpha_{\lambda}^{(3)}}{\lambda} \in K^{3\delta}\right)$$

$$\leq e^{-(1-\varepsilon)\lambda^{2}C_{\delta,N}} \prod_{i=1}^{MN} \left(Ee^{(1-\varepsilon)\lambda_{i}^{2} \frac{T}{MN} \left((1 + \frac{2\pi MN}{\lambda T}) \mathcal{I}\left(\frac{|\widehat{a}_{i}(T/(MN))|_{2\pi}}{\lambda_{i}(T/MN)}\right) + c_{2} \frac{2\pi MN}{\lambda_{i}T}\right)}\right)$$

where $C_{\delta,N} = \inf_{g \in K^{3\delta}} \mathcal{J}_N(g)$. Using the bound proved in Lemma 11 we get that

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\frac{\alpha_{\lambda}^{(3)}}{\lambda} \in K^{3\delta}\right) \le -(1 - \varepsilon)C_{\delta, N} \tag{54}$$

We now turn to the first error term. Using the fact that $\lfloor \alpha_{\lambda} \rfloor_{2\pi}$ is non-decreasing (which follows from (i) of Proposition 5) we get that

$$\|\alpha_{\lambda}^{(1)} - \alpha_{\lambda}\| \le \alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) + \lambda e^{-\frac{\beta}{4}T},$$

where $\alpha_{\lambda}(\infty)$ is the limit of $\alpha_{\lambda}(t)$ as $t \to \infty$. Choose T large enough so that $e^{-\frac{\beta}{4}T} \le \delta/2$. Then

$$P(\|\alpha_{\lambda}^{(1)} - \alpha_{\lambda}\| \ge \delta\lambda) \le P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) \ge \delta\lambda/2)$$

We will deal with this tail probability in Proposition 13 below. In particular, we will show that there is a constant $c_1 > 0$ so that

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\|\alpha_{\lambda}^{(1)} - \alpha_{\lambda}\| \ge \delta\lambda) \le -c_1 T \delta^2.$$
 (55)

For the second error term we first note that $\|\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}\| = \sup_{t \in [0,T]} |\alpha_{\lambda}^{(2)}(t) - \alpha_{\lambda}^{(1)}(t)|$. Using the coupling of Proposition 5 we can show that on [0,T] the process $\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}$ will have the same distribution as the solution of the SDE (15) with initial condition 0 and drift $\lambda(\mathfrak{f}_N - \mathfrak{f}) \geq 0$. Moreover, this process will be non-negative (because the drift is non-negative), and since $\lambda(\mathfrak{f}_N(t) - \mathfrak{f}(t)) \leq \lambda \frac{\beta T}{4N}$ for $t \in [0,T]$, it will be bounded by the solution of the SDE (15) with a constant drift $\lambda \frac{\beta T}{4N}$. Because of this $\|\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}\|$ is stochastically bounded by $\sup_{t \in [0,T]} \widetilde{\alpha}_{\lambda} \frac{\beta T}{4N}(t) \leq \widetilde{\alpha}_{\lambda} \frac{\beta T}{4N}(T) + 2\pi$ with $\widetilde{\alpha}_{\lambda}$ from (12, using the fact that $[\widetilde{\alpha}_{\lambda}(t)]_{2\pi}$ is non-decreasing. Thus for $\delta\lambda > 4\pi$ we have

$$P(\|\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}\| \ge \delta \lambda) \le P\left(\tilde{\alpha}_{\lambda \frac{\beta T}{4N}}(T) \ge \frac{1}{2}\delta \lambda\right).$$

If N and T are fixed then if λ is big enough then we can apply Lemma 10 for the right hand side with $\tilde{\lambda} = \lambda \frac{\beta T}{4N}$, t = T and $q = \frac{\frac{1}{2}\delta\lambda}{T\lambda\frac{\beta T}{4N}} = \frac{2\delta N}{\beta T^2}$. This leads to

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\|\alpha_{\lambda}^{(2)} - \alpha_{\lambda}^{(1)}\| \ge \frac{1}{2}\delta\lambda) \le -\frac{\beta^2 T^3}{4^2 N^2} \mathcal{I}\left(\frac{2\delta N}{\beta T^2}\right). \tag{56}$$

For the third error term we first note that

$$\|\alpha_{\lambda}^{(3)} - \alpha_{\lambda}^{(2)}\| \le \sup_{t \in [0,T]} |\alpha_{\lambda}^{(3)}(t) - \alpha_{\lambda}^{(2)}(t)| \le \max_{i} \sup_{t \in [Ti/N, T(i+1)/N]} |\alpha_{\lambda}^{(3)}(t) - \alpha_{\lambda}^{(2)}(t)|,$$

and thus

$$P(\|\alpha_{\lambda}^{(3)} - \alpha_{\lambda}^{(2)}\| \ge \delta\lambda) \le \sum_{i=0}^{N-1} P\left(\sup_{t \in [Ti/N, T(i+1)/N]} |\alpha_{\lambda}^{(3)}(t) - \alpha_{\lambda}^{(2)}(t)| \ge \delta\lambda\right).$$

In the interval [Ti/N, T(i+1)/N] the process $\alpha_{\lambda}^{(2)}$ solves the SDE (12) with constant drift $\lambda f_N(Ti/N)$. Here we can use the same steps that we used in the proof of Theorem 7 between (41) and (42) to get



$$P(\|\alpha_{\lambda}^{(3)} - \alpha_{\lambda}^{(2)}\| \ge \delta\lambda) \le \sum_{i=1}^{N-1} MP\left(\tilde{\alpha}_{\lambda f_N(Ti/N)}(T/(MN)) \ge \delta\lambda/2\right)$$
$$\le MNP\left(\tilde{\alpha}_{\frac{\beta}{4}\lambda}(T/(MN)) \ge \delta\lambda/2\right)$$

for λ big enough compared to δ^{-1} . For large enough λ we can apply Lemma 10 for the right hand side with $\tilde{\lambda} = \frac{\beta}{4}\lambda$, t = T/(MN) and $q = \frac{2\delta MN}{\beta T}$ to get

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\|\alpha_{\lambda}^{(3)} - \alpha_{\lambda}^{(2)}\| \ge \delta\lambda) \le -\frac{\beta^2}{4^2} \frac{T}{MN} \mathcal{I}\left(\frac{2\delta MN}{\beta T}\right). \tag{57}$$

Now taking (53) with the bounds (54), (55), (56) and (57) we get

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^{2}} \log P\left(\frac{\alpha_{\lambda}}{\lambda} \in K\right) \\
\leq \max \left\{ -(1-\varepsilon)C_{\delta,N}, -c_{1}T\delta^{2}, -\frac{\beta^{2}T^{3}}{4^{2}N^{2}}\mathcal{I}\left(\frac{2\delta N}{\beta T^{2}}\right), -\frac{\beta^{2}}{4^{2}}\frac{T}{MN}\mathcal{I}\left(\frac{2\delta MN}{\beta T}\right) \right\}. \tag{58}$$

Taking N to ∞ the last two terms go to $-\infty$ (using the bounds (37)) while the first term converges to $(1 - \varepsilon)C_{\delta}^{T}$ with

$$C_{\delta}^{T} = \inf_{g \in K^{3\delta}} \int_{0}^{T} \mathfrak{f}^{2}(t) \mathcal{I}\left(g'(t)/\mathfrak{f}(t)\right) dt$$

Letting now $T \to \infty$ and then $\varepsilon \to 0$ we get

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\frac{\alpha_{\lambda}}{\lambda} \in K\right) \le -\inf_{g \in K^{3\delta}} \mathcal{J}_{\operatorname{Sine}_{\beta}}(g).$$

Finally taking $\delta \to 0$ and using the fact that $\mathcal{J}_{Sine_{\beta}}$ is a good rate function gives the result

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\frac{\alpha_{\lambda}}{\lambda} \in K\right) \le -\inf_{g \in K} \mathcal{J}_{\operatorname{Sine}_{\beta}}(g).$$

This completes the proof of the lower bound.

We now prove the tail bound for the proof of the lower bound.

Proposition 13 Fix T, $\delta > 0$, then there is a constant c > 0 so that

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) \ge \delta \lambda) \le -cT\delta^2.$$
 (59)



Proof of Proposition 13 Take v = 1/8, and set $T_k = \frac{k(k+1)}{2}\theta T$ where the value of $\theta > 0$ will be specified later. Then we can break up the probability in question as

$$P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) \ge \delta\lambda) \le \sum_{k=1}^{\lfloor 2\sqrt{\lambda} \rfloor} P(\alpha_{\lambda}(T_{k+1}) - \alpha_{\lambda}(T_{k}) \ge \frac{\delta}{4}\lambda\nu k^{-(1+\nu)}) + P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T_{\lfloor 2\sqrt{\lambda} \rfloor + 1}) \ge \delta\lambda/2).$$
 (60)

Note, that for any fixed s>0 the process $\widehat{\alpha}_{s,\lambda}(t)=\alpha_{\lambda}(s+t)$ satisfies the SDE (3) with $\widehat{\lambda}=\lambda e^{-\frac{\beta}{4}s}$ with initial condition $\alpha_{\lambda}(s)$. Using the coupling techniques of Propositions 4 and 5 one can show that $\widehat{\alpha}_{s,\lambda}(t)-\widehat{\alpha}_{s,\lambda}(0)=\alpha_{\lambda}(s+t)-\alpha_{\lambda}(s)$ is stochastically dominated by $\widetilde{\alpha}_{\lambda\mathfrak{f}(s)}+2\pi$. This (together with $T_{k+1}-T_k=\theta(k+1)T$) gives

$$P(\alpha_{\lambda}(T_{k+1}) - \alpha_{\lambda}(T_k) \ge \frac{\delta}{4} \lambda \nu k^{-(1+\nu)}) \le P(\widetilde{\alpha}_{\lambda \mathfrak{f}(T_k)}(\theta(k+1)T) \ge \frac{\delta}{4} \lambda \nu k^{-(1+\nu)} - 2\pi)$$

$$\le P(\widetilde{\alpha}_{\lambda \mathfrak{f}(T_k)}(\theta(k+1)T) \ge \frac{\delta}{8} \lambda \nu k^{-(1+\nu)})$$

where the last bound follows for big enough λ from $k \leq 2\sqrt{\lambda}$. We can use bound (30) of Lemma 10 for the probability on the right with $\tilde{\lambda} = \lambda \mathfrak{f}(T_k)$, $t = \theta(k+1)T$ and $q = \frac{\delta \nu k^{-(1+\nu)}}{8\theta T(k+1)\mathfrak{f}(T_k)}$, since with these choices $qt\tilde{\lambda}$, q and $\tilde{\lambda}q\log q$ are all big, if we choose $\theta > 0$ small enough and then λ big enough. This leads to

$$\begin{split} &P(\widetilde{\alpha}_{\lambda\mathfrak{f}(T_{k})}(T) \geq \frac{\delta}{4}\lambda\nu k^{-(1+\nu)}) \\ &\leq \exp\left(-c_{1}\frac{\delta^{2}}{8^{2}}\lambda^{2}\nu^{2}k^{-2(1+\nu)}\theta^{-1}(k+1)^{-1}T^{-1}\log^{2}\left(\frac{\delta\nu k^{-(1+\nu)}}{8\theta(k+1)T\mathfrak{f}(T_{k})}\right)\right) \\ &\leq \exp\left(-c_{2}\delta^{2}\lambda^{2}k^{-3-2\nu}T^{-1}\left(c_{3} + \frac{\beta}{4}T\frac{k(k+1)}{2}\right)^{2}\right) \leq \exp\left(-c_{4}\lambda^{2}\delta^{2}Tk^{1-2\nu}\right), \end{split}$$

with a positive constant c_4 , which in turn implies (for large enough λ)

$$\sum_{k=1}^{\lfloor 2\sqrt{\lambda}\rfloor} P\left(\alpha_{\lambda}(T_{k+1}) - \alpha_{\lambda}(T_k) \ge \frac{\delta}{4}\lambda\nu k^{-(1+\nu)}\right) \le 2\exp\left(-c_4\lambda^2\delta^2T\right). \tag{61}$$

Lastly we bound the remaining term using Proposition 5:

$$P(\alpha(\infty) - \alpha(T\lambda) \ge \delta\lambda/2) = P(\xi_{\lambda\mathfrak{f}(\lambda T)}(\infty) \ge \lfloor\delta\lambda/2\rfloor) \le 2\left(e^{-\frac{\beta}{4}\lambda T}\right)^{\lfloor\delta\lambda/2\rfloor},$$

which together with (60) and (61) gives us the necessary upper bound for (59). \square *Proof of the lower bound in Theorem 6* We will show that if $g \in C[0, \infty)$ with $\mathcal{J}_{Sine_{\beta}}(g) < \infty$ then

$$\lim_{\varepsilon \to 0} \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(\|\lambda^{-1}\alpha_{\lambda}(\cdot) - g(\cdot)\| \le \varepsilon) \ge -\mathcal{J}_{\operatorname{Sine}_{\beta}}(g). \tag{62}$$



From this the lower bound will follow.

In Proposition 14 of the "Appendix" we will prove that if $\mathcal{J}_{\mathrm{Sine}_{\beta}}(g) < \infty$ then $g(\infty) = \lim_{t \to \infty} g(t) < \infty$ exists. Let $\varepsilon > 0$ and choose T > 0 so that

$$g(\infty) - g(T) \le \varepsilon/2$$
, and $e^{-\frac{\beta}{4}T} \le \varepsilon/4$. (63)

From the first assumption in (63) and the Markov property we have

$$P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \leq \varepsilon, t \geq 0)$$

$$\geq P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \leq \varepsilon/2, t \in [0, T], |\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T)| \leq \lambda\varepsilon/4)$$

$$\geq P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \leq \varepsilon/2, t \in [0, T])$$

$$\times \sup_{x} P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) \leq \lambda\varepsilon/4 |\alpha_{\lambda}(T) = x). \tag{64}$$

Using the same line of reasoning as in the proof of Proposition 13 (see after (60)) we get that with $\lambda_T = \lambda e^{-\frac{\beta}{4}T}$ we have

$$P(\alpha_{\lambda}(\infty) - \alpha_{\lambda}(T) \le \lambda \varepsilon / 4 | \alpha_{\lambda}(T) = x)$$

$$\ge P(\alpha_{\lambda_{T}}(\infty) \le \lambda \varepsilon / 4 - 2\pi) \ge P(\alpha_{\lambda_{T}}(\infty) \le \lambda \varepsilon / 8),$$

where the second inequality follows if λ is big enough compared to ε . Now we can use part (iii) of Proposition 5 with $f(t) = \lambda_T f(t)$, k = 1 and $a = \lambda \varepsilon/8$ to get

$$P(\alpha_{\lambda_T}(\infty) \le \lambda \varepsilon/8) = 1 - P(\alpha_{\lambda_T}(\infty) > \lambda \varepsilon/8) \ge 1 - 2\frac{8\lambda_T}{2\pi\lambda\varepsilon} \ge 1 - \frac{2}{\pi},$$

where the last step follows from the second assumption of (63).

Using this with (64) we get that

$$\begin{split} & \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(|\lambda^{-1} \alpha_{\lambda}(t) - g(t)| \le \varepsilon, t \ge 0) \\ & \ge \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(|\lambda^{-1} \alpha_{\lambda}(t) - g(t)| \le \varepsilon/2, t \in [0, T]), \end{split}$$

and it is enough to estimate the right hand side. We do this by introducing the process $\xi_N(t)$ on [0, T] which is a solution of the SDE (15) with initial condition 0 and the piecewise constant drift function λf_N where f_N is defined as in (52). From Proposition 5 we have that $\alpha_{\lambda}(t) \leq \xi_N(t)$ and $\widehat{\xi}_N(t) = \xi_N(t) - \alpha_{\lambda}(t)$ satisfies SDE (16) with initial condition 0 and drift $\lambda(f_N(t) - f(t))$. We have

$$P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \le \varepsilon/2, t \in [0, T]) \ge P(|\lambda^{-1}\xi_{N}(t) - g(t)| \le \varepsilon/4, t \in [0, T])$$

$$-P\left(\sup_{t \in [0, T]} |\xi_{N}(t) - \alpha_{\lambda}(t)| \ge \lambda\varepsilon/4\right). \tag{65}$$



The second term on the right may be bounded in the same manner as (56). this gives us

$$\limsup_{\lambda \to \infty} \frac{1}{\lambda^2} \log P \left(\sup_{t \in [0,T]} |\xi_N(t) - \alpha_\lambda(t)| \ge \lambda \varepsilon / 4 \right) \le -\frac{\beta^2 T^3}{4^2 N^2} \mathcal{I} \left(\frac{\varepsilon N}{\beta T^2} \right).$$

Note, that as $N \to \infty$ the right hand side converges to $-\infty$.

The only thing left is to estimate the first term on the right of (65). Introduce the notation $t_k = \frac{Tk}{N}$. We start with the bound

$$\log P(|\lambda^{-1}\xi_N(t) - g(t)| \le \varepsilon/4, t \in [0, T])$$

$$\ge P(|\lambda^{-1}(\xi_N(s + t_k) - \xi_N(t_k)) - (g(s + t_k) - g(t_k))| \le \varepsilon/(4N), s \in [0, T/N]).$$

For any fixed k the process $\xi_N(s+t_k)$, $s \in [0, T/N]$ satisfies the SDE (16) with initial condition $\xi_N(t_k)$ and a constant drift $\lambda f_N(t_k)$. Using the coupling in the proof of Proposition 4 we can construct independent processes $\widehat{\alpha}_k(t)$, $t \in [0, T/N]$ so that

$$\widehat{\alpha}_k(s) - 2\pi \le \xi_N(s + t_k) - \xi_N(t_k) \le \widehat{\alpha}_k(t) + 2\pi, \quad s \in [0, T/N]$$

and $\widehat{\alpha}_k(t)$, $t \in [0, T/N]$ has the same distribution as $\widetilde{\alpha}_{\lambda f_N(t_k)}(t)$, $t \in [0, T/N]$. From this it immediately follows that

$$\begin{split} & \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(|\lambda^{-1} \xi_N(t) - g(t)| \le \varepsilon/4, t \in [0, T]) \\ & \ge \sum_{k=0}^{N-1} \liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P\left(\left| \frac{\tilde{\alpha}_{\lambda f_N(t_k)}(s)}{\lambda} - (g(s+t_k) - g(t_k)) \right| < \frac{\varepsilon}{8N}, s \in \left[0, \frac{T}{N}\right] \right). \end{split}$$

From our path level large deviation lower bound on $\tilde{\alpha}$ we get

$$\lim_{\lambda \to \infty} \inf_{\lambda^{2}} \log P(|\lambda^{-1} \widetilde{\alpha}_{\lambda f_{N}(t_{k})}(s) - (g(s + t_{k}) - g(t_{k}))| \le \varepsilon/(4N), s \in [0, T/N])$$

$$\ge - \inf_{\substack{|\tilde{g}(s) - g(s)| < \varepsilon/(8N) \\ s \in [t_{k}, t_{k+1}]}} f_{N}(t_{k})^{2} \int_{0}^{T/N} \mathcal{I}(f_{N}(t_{k})^{-1} \tilde{g}'(t_{k} + s)) ds.$$

This yields the estimate

$$\lim_{\lambda \to \infty} \inf_{\lambda^{2}} \frac{1}{\lambda^{2}} \log P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \le \varepsilon/4, t \in [0, T])$$

$$\ge - \inf_{\substack{|\tilde{g}(s) - g(s)| < \varepsilon/(8N), \\ s \in [0, T]}} \int_{0}^{T} \mathfrak{f}_{N}(s)^{2} \mathcal{I}(\mathfrak{f}_{N}^{-1}\tilde{g}'(s)) ds.$$



Letting $N \to \infty$ the lower bound converges to $-\int_0^T f(s)^2 \mathcal{I}(f^{-1}(s)g'(s))ds$ which (together with our previous estimates) shows that

$$\liminf_{\lambda \to \infty} \frac{1}{\lambda^2} \log P(|\lambda^{-1}\alpha_{\lambda}(t) - g(t)| \le \varepsilon, t \ge 0) \ge -\int_0^T \mathfrak{f}(s)^2 \mathcal{I}(\mathfrak{f}^{-1}(s)g'(s)) ds.$$

Letting $\varepsilon \to 0$ we also have $T = T_{\varepsilon} \to \infty$ which yields the bound (62) and concludes the proof of the lower bound in the large deviation principle.

6 \mathcal{J}_{Sch} and $\mathcal{J}_{Sine_{\theta}}$ are good rate functions

In this section we will show that $\mathcal{J}_{Sch,T}$ and $\mathcal{J}_{Sine_{\beta}}$ are good rate functions. Our main tools are the bound (37) and the estimate

$$\mathcal{I}(x) > c_1(x-1)^2$$
, if $x > 0$ (66)

both of which will be proved in Proposition (17) of the "Appendix".

Proposition 14 The functions $\mathcal{J}_{Sine_{\beta}}(\cdot)$ and $\mathcal{J}_{Sch,T}(\cdot)$ are both good rate functions on the spaces $C[0,\infty)$ and C[0,T] respectively. Moreover, if $g \in C[0,\infty)$ and $\mathcal{J}_{Sine_{\beta}}(g) < \infty$ then $\lim_{t \to \infty} g(t)$ is finite.

Proof Fix T>0 and $r\geq 0$. In order to prove that $K_r=\{g: \mathcal{J}_{\mathrm{Sch},T}(\cdot)\leq r\}$ is compact we first show the equicontinuity of this set. Suppose that $g\in K_r$. Then g(0)=0 and $g'(x)\geq 0$ exists a.e. in [0,T]. We have for $0\leq x\leq y\leq T$

$$|g(x) - g(y) - (x - y)| = \left| \int_{x}^{y} (g'(s) - 1) ds \right| \le (y - x)^{1/2} \sqrt{\int_{x}^{y} (g'(s) - 1)^{2} ds}$$

$$\le c(y - x)^{1/2} \sqrt{\int_{x}^{y} \mathcal{I}(g'(s)) ds} \le c(y - x)^{1/2} r^{1/2}$$

where we used (66) in the second step. This shows that K_r is equicontinuous. Using Tonelli's semicontinuity theorem (e.g. Theorem 3.5, [5]) the compactness of K_r now follows.

The proof for $\mathcal{J}_{\mathrm{Sine}_{\beta}}(\cdot)$ is bit more involved. Fix $\beta > 0$. It is convenient to transform the interval $[0, \infty)$ into [0, 1) using the function $y = 1 - e^{-\beta t/4}$. Then for a $g \in C[0, \infty)$ with $\mathcal{J}_{\mathrm{Sine}_{\beta}}(g) < \infty$ we have

$$\mathcal{J}_{\operatorname{Sine}_{\beta}}(g) = \int_{0}^{\infty} \mathfrak{f}^{2}(t) \mathcal{I}(g'(t) \mathfrak{f}^{-1}(t)) dt = \frac{\beta}{4} \int_{0}^{1} (1 - y) \mathcal{I}(\tilde{g}'(y)) dy$$



where $\tilde{g}(y) = g(-\frac{4}{\beta}\log(1-y))$, $\tilde{g} \in C[0,1)$. Consider the functional $\tilde{\mathcal{J}}_{Sine}(\cdot)$ on C[0,1) defined as

$$\tilde{\mathcal{J}}_{\text{Sine}}(g) = \frac{\beta}{4} \int_0^1 (1 - t) \mathcal{I}(g'(t)) dt \tag{67}$$

if g'(t) exists and non-negative for a.e. $0 \le t < 1$, and as ∞ otherwise. Clearly, if we show that $\tilde{\mathcal{J}}_{Sine}(\cdot)$ is a good rate function on C[0,1) then the same will hold for $\mathcal{J}_{Sine_{\beta}}$. We first show that if $g \in C[0,1)$ and $\tilde{\mathcal{J}}_{Sine}(g) < \infty$ then $\lim_{y \to 1^{-}} g(y)$ is finite, i.e. we can consider $\tilde{\mathcal{J}}_{Sine}(\cdot)$ on C[0,1]. We have

$$\lim_{y \to 1^{-}} g(y) = \int_{0}^{1} g'(y) dy \le 2 + \int_{0}^{1} g'(y) \mathbf{1}(g'(y) \ge 2) dy.$$

We will prove that

if
$$h(y) \ge 0$$
, and $\int_0^1 (1 - y)h(y)^2 \log^2(h(y) + e)dy < \infty$,
then $\int_0^1 h(y)dy < \infty$. (68)

Using this with $h(y) = g'(y)\mathbf{1}(g'(y) \ge 2)$ together with the bound in (37) we get the boundedness of $\int_0^1 g'(y)dy$ and the existence on $\lim_{y\to 1^-} g(y)$.

Let $\Phi(x) = x^2 \log^2(|x| + e)$, this is a strictly convex, even function with $\lim_{x\to 0} \frac{\Phi(x)}{x} = 0$, $\lim_{x\to \infty} \frac{\Phi(x)}{x} = \infty$ (i.e. Φ is a 'nice Young function'). Introduce the complementary function

$$\Psi(x) = \Phi^*(x) = \sup_{y \ge 0} \{y|x| - \Phi(y)\} = \int_0^{|x|} (\Phi')^{(-1)}(y) dy$$

where $(\Phi')^{(-1)}$ is the inverse of the strictly increasing function Φ' on $[0, \infty)$. Assume that

$$A = \int_0^1 (1 - y)\Phi(h(y))dy < \infty \tag{69}$$

and let μ the measure on [0, 1] with $d\mu = \frac{1}{A}(1-x)dx$. Consider the Orlicz spaces

$$\begin{split} L_{\mu}^{\Phi} &= \left\{ f : \text{there is an } a > 0 \text{ with } \int_{[0,1]} \Phi(af) d\mu < \infty \right\}, \\ L_{\mu}^{\Psi} &= \left\{ f : \text{there is an } a > 0 \text{ with } \int_{[0,1]} \Psi(af) d\mu < \infty \right\} \end{split}$$



with the Luxemburg-norms defined as

$$||f||_{\Phi} = \inf \left\{ b > 0 : \int_{[0,1]} \Phi(b^{-1}f) d\mu \le 1 \right\},$$

$$||f||_{\Psi} = \inf \left\{ b > 0 : \int_{[0,1]} \Psi(b^{-1}f) d\mu \le 1 \right\}.$$
(70)

(See e.g. [21] for more on Orlicz spaces.) Note, that by our assumption (69) we have $\|h\|_{\Phi} \leq 1$. By the generalized Hölder inequality for Orlicz spaces (c.f. Theorem 3 in Chapter III of [21]), for any $f \in L^{\Psi}_{\mu}$ one has

$$||fh||_1 \le 2||f||_{\Psi} ||h||_{\Phi} \le 2||f||_{\Psi} \tag{71}$$

where $\|\cdot\|_1$ is the L^1 norm on [0,1] with reference measure μ . Choose $f(x)=\frac{1}{1-x}$. If we show that $\|f\|_{\Psi}<\infty$ then this would imply

$$\infty > 2\|f\|_{\Psi} \ge \|fh\|_1 = \frac{1}{A} \int_0^1 \frac{1}{1-x} h(x)(1-x) dx = \frac{1}{A} \int_0^1 h(x) dx,$$

and the statement (68) would follow. It is not hard to check, that there is a c>0 so that

$$\Psi(x) \le c \frac{x^2}{\log^2(x+e)}, \quad \text{for } x \ge 0.$$
 (72)

Since the integral $\int_0^1 (1-x) \frac{(1-x)^{-2}}{\log^2((1-x)^{-1}+e)} dx$ is finite, this implies that $\|\frac{1}{1-x}\|_{\Psi}$ is finite and thus $\int_0^1 h(y) dy < \infty$. This completes the proof that if $\tilde{\mathcal{J}}_{Sine}(g) < \infty$ then $\lim_{y \to 1^-} g(y)$ is finite, and also shows the last statement of the proposition.

Next we will prove the equicontinuity of the set $K_r = \{f : \tilde{\mathcal{J}}_{Sine}(f) \leq r\}$, we will show that if $g \in K_r$ then for $\varepsilon < \varepsilon_0$ we have

$$|g(a+\varepsilon) - g(a)| \le C(\log \varepsilon^{-1})^{-1/3} \quad \text{for any } a \in [0, 1-\varepsilon].$$

Here ε_0 , C only depend on r.

We first assume $a \leq 1 - \sqrt{\varepsilon}$. Then

$$\begin{split} |g(a+\varepsilon)-g(a)-\varepsilon| &= \left|\int_a^{a+\varepsilon} (g'(y)-1)dy\right| \\ &\leq \left(\int_a^{a+\varepsilon} \frac{1}{1-y}dy\right)^{1/2} \left(\int_a^{a+\varepsilon} (1-y)(g'(y)-1)^2dy\right)^{1/2} \\ &\leq Cr^{1/2} \left(\log\left(1+\frac{\varepsilon}{1-a-\varepsilon}\right)\right)^{1/2} \leq Cr^{1/2}\varepsilon^{1/4} \end{split}$$



Where we used $1 - a > \sqrt{\varepsilon}$, the bound (66) and the fact that ε can be chosen to be small enough.

Next we assume that $a>1-\sqrt{\varepsilon}$. Because of the monotonicity of g it is enough to bound $|g(1)-g(1-\sqrt{\varepsilon})|$. Setting $f(x)=\frac{1}{1-x}$ and $h(x)=g'(x)\mathbf{1}(g'(x)\geq 2)$ we have

$$g(1) - g(1 - \sqrt{\varepsilon}) \le 2\sqrt{\varepsilon} + \int_{1 - \sqrt{\varepsilon}}^{1} h(x)dx.$$
 (74)

Since $\tilde{\mathcal{J}}_{Sine}(g) \leq r$, we can assume that (69) holds with some finite A > 0. We will now follow the previous argument using Orlicz spaces. We use the same definitions for Ψ , Φ , μ but for the norms $\|\cdot\|_{\Psi}$, $\|\cdot\|_{\Phi}$ defined in (70) we use the interval $[1-\sqrt{\varepsilon},1]$ instead of [0,1].

Using inequality (74) and (71) we get the bound

$$g(1) - g(1 - \sqrt{\varepsilon}) \le 2\sqrt{\varepsilon} + A \int_{1 - \sqrt{\varepsilon}}^{1} f(x)h(x)d\mu(x) \le 2\sqrt{\varepsilon} + A\|f\|_{\Psi}.$$

To estimate $||f||_{\Psi}$ we will prove that with $b=(\log \varepsilon^{-1})^{-1/3}$ there is a constant ε_0 depending on A so that

$$\int_{1-\sqrt{\varepsilon}}^1 \Psi(b^{-1}f(x)) d\mu(x) = A^{-1} \int_{1-\sqrt{\varepsilon}}^1 (1-x) \Psi(b^{-1}(1-x)^{-1}) dx < 1, \quad \text{for } \varepsilon < \varepsilon_0.$$

This will imply that for such ε we have $||f||_{\Psi} \le b$. Using (72) we get

$$\begin{split} A^{-1} \int_{1-\sqrt{\varepsilon}}^{1} (1-x) \Psi(b^{-1}(1-x)^{-1}) dx &\leq c b^{-2} A^{-1} \int_{0}^{\sqrt{\varepsilon}} x^{-1} \frac{1}{\log^2(2x^{-1}b^{-1})} dx \\ &\leq \frac{c (\log \varepsilon^{-1})^{2/3}}{A \log \left(2 (\log \varepsilon^{-1})^{-1/3} \varepsilon^{-1/2}\right)}. \end{split}$$

Since the right hand side converges to 0 as $\varepsilon \to 0$ we get that $||f||_{\Psi} \le b$ for small enough ε which in turn leads to the upper bound (73). This completes the proof of the equicontinuity of the set K_r and the compactness follows again by Tonelli's theorem.

7 From the path to the endpoint

In this section we will complete the proofs of Theorems 1 and 2.

Proof of Theorem 2 Consider the continuous map $F:C[0,T]\to\mathbb{R}$ given by $F(g)=g(T)/(2\pi)$. By the contraction principle (see e.g. [9]) the random variables $\frac{1}{\lambda}\frac{\alpha_{\lambda}(T)}{2\pi}$ satisfy a large deviation principle with scale function λ^2 and good rate function J defined as



$$J(\rho) = \min \left\{ \int_0^T \mathcal{I}(g'(t))dt : g'(t) \ge 0, \ g(T) = 2\pi\rho \right\}. \tag{75}$$

We will now solve this variational problem. If g provides the minimum then we can assume that g' is monotone decreasing. To see this define \tilde{g} with $\tilde{g}(0) = 0$ and $\tilde{g}'(t) = \sup\{x : m(g'(s) \ge x) \ge t\}$ where m indicates Lesbegue measure. Then $\tilde{g}(T) = g(T), \mathcal{J}_{Sch,T}(\tilde{g}) = \mathcal{J}_{Sch,T}(g), \text{ and } \tilde{g}'(t) \text{ is decreasing. If } g' > 0 \text{ on } [0,a]$ and g'=0 on (a,1] then by the classical variational method we get that $\mathcal{I}'(g'(x))$ is constant on [0,a]. This means that $g'(x)=\frac{2\pi\rho}{a}$ on [0,a] and g'(x)=0 on (a,T]and our variational problem is reduced to finding the minimum of

$$f(a) = a\mathcal{I}\left(\frac{2\pi\rho}{a}\right) + (T-a)\mathcal{I}(0), \quad \text{on} \quad 0 \le a \le T.$$

But we have

$$f'(a) = \mathcal{I}\left(\frac{2\pi\rho}{a}\right) - \mathcal{I}(0) - \mathcal{I}'\left(\frac{2\pi\rho}{a}\right)\frac{2\pi\rho}{a} < 0,$$

since \mathcal{I} is strictly convex, which means that the minimum is at a = T. Thus

$$J(\rho) = \min \left\{ \int_0^T \mathcal{I}(g'(t))dt, \ g'(t) \ge 0, \ g(T) = 2\pi\rho \right\} = T\mathcal{I}(2\pi\rho/T) \quad (76)$$

is the large deviation rate function for $\frac{1}{\lambda} \frac{\alpha_{\lambda}(T)}{2\pi}$. Now recall that the counting function of Sch_{τ} is given by

$$\widetilde{N}_{\tau}(\lambda) = \#\{\nu : 0 \le \nu \le \lambda, \phi_{\lambda/\tau}(\tau) \in 2\pi\mathbb{Z}\}\$$

where ϕ_{λ} is the solution of (10). Note, that $\phi_{\lambda}(t) - \phi_0(t)$ has the same distribution as $\xi_{f,0}(t)$ with constant $f = \lambda$, which in turn has the same distribution as $\widetilde{\alpha}_{\lambda}(t)$. Using the coupling methods of Proposition 5 we can show that $\phi_{\lambda}(t)$ is increasing in λ for any fixed t (see [19] for a detailed proof of this fact). From this it follows that

$$\left| \tilde{N}_{\tau}(\lambda) - \frac{1}{2\pi} \left(\phi_{\lambda/\tau}(\tau) - \phi_0(\tau) \right) \right| \leq 1.$$

This means that in order to get a large deviation principle for $\frac{1}{\lambda}\tilde{N}_{\tau}(\lambda)$ it is enough to prove one for $\frac{1}{\lambda}\frac{\phi_{\lambda/\tau}(\tau)-\phi_0(\tau)}{2\pi}$. But this has the same distribution as $\frac{1}{\lambda}\frac{\tilde{\alpha}_{\lambda/\tau}(\tau)}{2\pi}$, and a simple rescaling of (76) completes the proof of the theorem.

Proof of Theorem 1 Theorem 6 shows that $\frac{1}{2}\alpha_{\lambda(\cdot)}$ satisfies a path level large deviation principle. By applying the time change $y = 1 - e^{-\frac{\beta}{4}t}$, we get that $t \to \frac{1}{\lambda}\alpha_{\lambda}(1 - e^{-\frac{\beta}{4}t})$ satisfies a path level LDP on C[0,1) with the modified rate function $\tilde{\mathcal{J}}_{Sine}$ given in (67). In Proposition 14 we showed that if $\tilde{\mathcal{J}}_{Sine}(g) < \infty$ then the limit as $t \to 1^$ exists and so the LDP actually holds on C[0, 1]. Using the contraction principle with



the functional $F(g)=\frac{1}{2\pi}g(1)$, we get that $\frac{1}{\lambda}\frac{\alpha_{\lambda}(\infty)}{2\pi}$ satisfies a large deviation principle with speed function λ^2 and a good rate function

$$\begin{split} J^{\beta}(\rho) &= \min \left\{ \tilde{\mathcal{J}}_{\text{Sine}}(g) : g(1) = 2\pi \rho \right\} \\ &= \min \left\{ \frac{\beta}{4} \int_{0}^{1} (1 - t) \mathcal{I}(g'(t)) dt : g(0) = 0, g'(t) \ge 0, g(1) = 2\pi \rho \right\}. \end{split}$$

The counting function $N_{\beta}(\lambda)$ of Sine_{β} is given by $\frac{\alpha_{\lambda}(\infty)}{2\pi}$, so Theorem 1 will follow if we can show that the solution of this variational problem is given by $\beta I_{\text{Sine}}(\rho)$ as defined in the theorem.

The function $\tilde{\mathcal{J}}_{\text{Sine}}$ is a good rate function, so for any $\rho \geq 0$ the minimum is achieved at some $g_{\rho} \in C[0, 1]$. Clearly, when $\rho = \frac{1}{2\pi}$ then the minimum is zero, as the g(t) = t function shows. (We will not denote the dependence of ρ in $g = g_{\rho}$ from this point.)

We may assume that for the minimizer the derivative g' will not take values from both $(1,\infty)$ and [0,1) because otherwise we could construct a function \hat{g} with the same boundary condition $\hat{g}(1)=2\pi\rho$, but with $\tilde{\mathcal{J}}_{\mathrm{Sine}}(g)>\tilde{\mathcal{J}}_{\mathrm{Sine}}(\hat{g})$. The construction is as follows. Assume $\rho<1/(2\pi)$ and that $A=\{t:g'(t)>1\}$ has positive measure. Since $\int_0^1 (g'(t)-1)dt=2\pi\rho-1<0$ and $\int_A (g'(t)-1)dt>0$, by the intermediate value theorem we can find $B\subset[0,1]\backslash A$ so that $\int_{A\cup B} (g'(t)-1)dt=0$. Define \hat{g} with $\hat{g}(0)=0$, $\hat{g}'(t)=g'(t)$ if $t\notin A\cup B$ and $\hat{g}'(t)=1$ otherwise. Then $g(1)=\int_0^1 g'(t)dt=\int_0^1 \hat{g}'(t)dt=\hat{g}(1)$, but clearly $\tilde{\mathcal{J}}_{\mathrm{Sine}}(g)>\tilde{\mathcal{J}}_{\mathrm{Sine}}(\hat{g})$. A similar construction works for $\rho>1/(2\pi)$. Thus we may assume that $g'(t)\leq 1$ for all t if $\rho<1/(2\pi)$, and $g'(t)\geq 1$ for all t if $\rho>1/(2\pi)$.

First assume that $\rho > \frac{1}{2\pi}$. Then $g'(t) \ge 1$ for all t and we can use the classical variational method (see e.g. [5]) to conclude that $(1-t)\mathcal{I}'(g'(t))$ is constant in t. Thus the optimizer is given by a function g_{ρ} which satisfies

$$g_{\rho}(0) = 0, \qquad \mathcal{I}'(g'(t)) = \frac{c_{\rho}}{1-t}, \qquad \int_{0}^{1} (\mathcal{I}')^{(-1)} \left(\frac{c_{\rho}}{1-t}\right) dt = 2\pi\rho,$$
 (77)

for some constant c_{ρ} and the solution of the variational problem is

$$J^{\beta}(\rho) = \frac{\beta}{4} \int_0^1 (1-t) \mathcal{I}\left((\mathcal{I}')^{(-1)} \left(\frac{c_{\rho}}{1-t}\right)\right) dt. \tag{78}$$

In Proposition 15 below we will show that this is equal to $\beta I_{Sine}(\rho)$ as defined in Theorem 1.

Now assume that $\rho < \frac{1}{2\pi}$, here we can assume that the minimizer satisfies $g'(t) \le 1$. As in the case of $\operatorname{Sch}_{\tau}$ we may assume g' is decreasing, this can be shown using the same construction as found in the paragraph directly following equation (75). Suppose that g' is zero for $t \in [a,1]$ and g'(t)>0 in [0,a]. Then on [0,a] the classical variational method shows that $(1-t)\mathcal{I}'(g'(t))$ must be constant. Thus the optimizer must be of the following form:



$$g'(t) = \begin{cases} (\mathcal{I}')^{(-1)} \left(\frac{c_{\rho,a}}{1-t}\right), & 0 \le t \le a \\ 0, & a < t \le 1, \end{cases}$$

$$(79)$$

for some constant c_{ρ} which satisfies

$$2\pi\rho = \int_0^a (\mathcal{I}')^{(-1)} \left(\frac{c_{\rho,a}}{1-t}\right) dt. \tag{80}$$

By Propositions 16 and 17 of the "Appendix" the function $\mathcal{T}'(x)$ is strictly increasing on $(0,\infty)$ with a limit of $-\frac{1}{2\pi}$ at x=0. Thus $c_{\rho,a}$ in (79) cannot be smaller than $-\frac{1-a}{2\pi}$. Our next claim is that the optimizer has a continuous derivative at t=a, which will identify $c_{\rho,a}$ as $c=-\frac{1-a}{2\pi}$. Assume the opposite, i.e. that $c_{\rho,a}>-\frac{1-a}{2\pi}$ and g'(a)>0. Let $\eta_\delta(x)=\mathbf{1}_{(a,a+\delta)}-\mathbf{1}_{(a-\delta,a)}$. If δ,ε are small enough then $g'-\varepsilon\eta_\delta\geq 0$ in [0,1] and $\tilde{g}(t)=\int_0^t (g'(s)-\varepsilon\eta_\delta(s))ds$ satisfies the same boundary conditions as g. Since g is a minimizer, the derivative of $h(\varepsilon)=\tilde{\mathcal{J}}_{\mathrm{Sine}}(g+\varepsilon\eta_\delta)$ at $\varepsilon=0$ cannot be negative. We can compute the derivative as

$$h'(0) = \frac{\beta}{4} \int_0^1 (1-t) \mathcal{I}'(g(t)) \eta_{\delta}(t) dt$$

= $-\int_{a-\delta}^a (1-t) \frac{c_{\rho,a}}{1-t} dt + \int_a^{a+\delta} (1-t) \left(-\frac{1}{2\pi}\right) dt.$

This is equal to $\delta(-c_{\rho,a}-\frac{1-a}{2\pi})+\frac{\delta^2}{4\pi}$ which is negative if δ is small enough (by our assumption that $c_{\rho,a}>-\frac{1-a}{2\pi}$). The contradiction shows that we must have $c=-\frac{1-a}{2\pi}$. Thus the optimizer is given by

$$g'(t) = \begin{cases} (\mathcal{I}')^{(-1)} \left(\frac{a-1}{2\pi(1-t)} \right), & 0 \le t \le a \\ 0, & a < t \le 1. \end{cases}$$
 (81)

for some 0 < a < 1 with

$$2\pi\rho = \int_0^a (\mathcal{I}')^{(-1)} \left(\frac{a-1}{2\pi(1-t)}\right) dt.$$
 (82)

and the solution of the variational problem in the $2\pi\rho$ < 1 case is given by

$$J^{\beta}(\rho) = \frac{\beta}{4} \int_0^a (1-t)\mathcal{I}\left((\mathcal{I}')^{(-1)} \left(\frac{a-1}{2\pi(1-t)}\right)\right) dt + \frac{\beta}{64}(1-a)^2.$$
 (83)

In Proposition 15 below we will show that this is equal to $\beta I_{\text{Sine}}(\rho)$.

Proposition 15 The rate function for the Sine_{β} process is given by

$$\beta I_{Sine}(\rho) = \frac{\beta}{8} \left[\frac{\nu}{8} + \rho \mathcal{H}(\nu) \right]$$



where $v = \gamma^{-1}(\rho)$, and γ is the strictly increasing function given in (7).

Proof We have to show that $J^{\beta}(\rho)$ defined by (77) and (78) for $\rho > 1/(2\pi)$ and by (82) and (83) for $\rho < 1/(2\pi)$ is equal to βI_{Sine} given above.

We begin with the case where $\rho > \frac{1}{2\pi}$. In this case the minimizer $g = g_{\rho}$ is given by (77). One easily checks that

$$\frac{d}{dt} \left(\frac{\beta}{8} \left(-(1-t)^2 \mathcal{I}(g'(t)) + c_{\rho}(1-t)g'(t) + c_{\rho}g(t) \right) \right) = \frac{\beta}{4} (1-t)\mathcal{I}(g'(t)). \tag{84}$$

From this we get

$$J^{\beta}(\rho) = \frac{\beta}{4} \int_{0}^{1} (1 - t) \mathcal{I}(g'(t)) dt = \frac{\beta}{8} \left[\mathcal{I}(g'(0)) - c_{\rho}g'(0) + 2\pi\rho c_{\rho} \right]$$

where we used g(0) = 0, $g(1) = 2\pi\rho$, and the limits

$$\lim_{t \to 1^{-}} (1-t)^{2} \mathcal{I}(g'(t)) = \lim_{x \to \infty} \frac{c_{\rho}^{2} \mathcal{I}(x)}{\mathcal{I}'(x)^{2}} = 0, \quad \lim_{t \to 1^{-}} (1-t)g'(t) = \lim_{x \to \infty} \frac{x}{c_{\rho} \mathcal{I}'(x)} = 0$$

which follow from the asymptotics (90) and (91) to be proven in Proposition 17.

Now for the case where $\rho < \frac{1}{2\pi}$ we have that g_{ρ} is given by (81). Using the notation $c = c_{\rho} = \frac{a-1}{2\pi}$, the identity (84) gives

$$J^{\beta}(\rho) = \frac{\beta}{4} \int_{0}^{2\pi c_{\rho}+1} (1-t) \mathcal{I}(g'(t)) dt + \frac{\beta}{8} (2\pi c_{\rho})^{2} \mathcal{I}(0)$$
$$= \frac{\beta}{8} \left[\mathcal{I}(g'(0)) - c_{\rho} g'(0) + 2\pi \rho c_{\rho} \right],$$

where we used g(0)=0, $g(a)=2\pi\rho$, and g'(a)=0. Note, that $c_{\rho}>0$ if $\rho>1/(2\pi)$ and $-\frac{1}{2\pi}\leq c_{\rho}<0$ if $\rho<1/(2\pi)$. Introducing $\nu=K^{(-1)}\left(\frac{\pi}{2(\mathcal{I}')^{(-1)}(c_{\rho})}\right)$, we get for both $\rho<1/(2\pi)$ and $\rho>1/(2\pi)$ that

$$J^{\beta}(\rho) = \frac{\beta}{8} \left(\frac{\nu}{8} + \rho \mathcal{H}(\nu) \right)$$

which agrees with (6), we just have to show that $\nu = \gamma^{(-1)}(\rho)$. Note, that $\nu = \nu(\rho) < 0$ if $2\pi\rho > 1$ and $0 < \nu < 1$ if $2\pi\rho < 1$.

Recall from (77) and (82) that

$$\rho = \frac{1}{2\pi} \int_0^1 (\mathcal{I}')^{(-1)} \left(\frac{c_\rho}{1-t}\right) dt, \quad \text{if } \rho > \frac{1}{2\pi}, \quad \text{and}$$

$$\rho = \frac{1}{2\pi} \int_0^{2\pi c + 1} (\mathcal{I}')^{(-1)} \left(\frac{c_\rho}{1-t}\right) dt, \quad \text{if } \rho < \frac{1}{2\pi}.$$



Applying the change of variables to both integrals with a new variable x satisfying $\frac{\pi}{2K(x)} = (\mathcal{I}')^{(-1)}(\frac{c_\rho}{1-t})$, we get that ρ depends on $\nu = K^{(-1)}\left(\frac{\pi}{2(\mathcal{I}')^{(-1)}(c_\rho)}\right)$ exactly via (7) which finishes the proof. Note, that the finiteness of the integrals in (7) follow from the asymptotics of K(x) and E(x) near 1 and $-\infty$ (see the proof of Proposition 17).

Acknowledgments B. Valkó was partially supported by the National Science Foundation CAREER award DMS-1053280. The authors would like to thank P. J. Forrester for helpful comments and additional references.

Appendix A: All about \mathcal{I}

In this section we prove the needed estimates about the function \mathcal{I} .

Proposition 16 The function $\mathcal{I}(x)$ is strictly convex and continuous on $(0, \infty)$. It has an absolute minimum at x = 1 where it is equal to 0.

Proof K is strictly increasing on $(-\infty, 1)$ which shows that $\mathcal{I}(x)$ is well-defined (and differentiable) on $(0, \infty)$. By differentiating (11) and using the identities

$$K'(x) = \frac{E(x) - (1 - x)K(x)}{2(1 - x)x}, \qquad E'(x) = \frac{E(x) - K(x)}{2x}$$

we can compute that

$$\mathcal{I}'(x) = \frac{1}{x}\mathcal{I}(x) - \frac{1}{8x}K^{-1}\left(\frac{\pi}{2x}\right),\tag{85}$$

and

$$\mathcal{I}''(x) = \frac{\pi}{16x^3} \frac{1}{K'(K^{-1}(\frac{\pi}{2x}))}.$$
 (86)

Observe that K'(y) > 0 for y < 1 which gives $\mathcal{I}''(x) > 0$ for x > 0 and the strict convexity of \mathcal{I} .

Using $K(0) = E(0) = \frac{\pi}{2}$ we get $\mathcal{I}(1) = \mathcal{I}'(1) = 0$ which (by the strict convexity) proves the second half of the proposition.

Proposition 17 We have $\lim_{x\to 0^+} \mathcal{I}(x) = \frac{1}{8}$ and $\lim_{x\to 0^+} \mathcal{I}'(x) = -\frac{1}{2\pi}$. There is a constant $c_1 > 0$ so that

$$\mathcal{I}(x) \ge c_1(x-1)^2$$
, for all x , and (87)

$$\frac{\mathcal{H}(-x)}{\sqrt{x}\log x}$$
, $\frac{\mathcal{I}(x)}{x^2\log^2 x}$ and $\frac{-K^{-1}(1/x)}{x^2\log^2 x}$ are bounded away from 0 and ∞ for $x > 2$.

(88)

Proof The following asymptotics can be readily derived from the definitions of elliptic integrals (or by the existing more sophisticated expansions c.f. [6,15]). There is a



constant c > 0 so that

$$\left| K(-a) - \frac{1}{2\sqrt{a}} \log(16a) \right| \le \frac{c}{a^{3/2}} \log(a),$$

$$\left| E(-a) - \sqrt{a} \right| \le \frac{c}{a^{1/2}} \log(a), \quad \text{for } a > 2.$$
(89)

From this it is easy to check that

$$\lim_{x \to \infty} \frac{\mathcal{H}(-x)}{\sqrt{x} \log x} = \frac{1}{2}, \qquad \lim_{x \to \infty} \frac{-K^{-1}(1/x)}{x^2 \log^2 x} = 1, \text{ and}$$

$$\lim_{x \to \infty} \frac{\mathcal{I}(x)}{x^2 \log^2 x} = \frac{1}{2\pi^2}.$$
(90)

This gives (88). Note, that together with (85) this also gives

$$\lim_{x \to \infty} \frac{\mathcal{I}'(x)}{x \log^2 x} = \frac{1}{\pi^2}.$$
 (91)

Using the functional identities

$$E(z) = \sqrt{1 - z}E\left(\frac{z}{z - 1}\right), \qquad K(z) = \frac{1}{\sqrt{1 - z}}K\left(\frac{z}{z - 1}\right), \qquad z \in (0, 1)$$

the asymptotics of (89) can be transformed into

$$K(a) \sim -\frac{1}{2}\log(1-a), \quad E(a) \sim 1, \quad \text{as } a \to 1^{-1}$$

with explicit error bounds for 2/3 < a < 1. From this we can obtain $\lim_{x \to 0^+} \mathcal{I}(x) = \frac{1}{8}$ and $\lim_{x \to 0^+} \mathcal{I}'(x) = -\frac{1}{2\pi}$. Using (88) with the continuity of \mathcal{I} and the fact that $\mathcal{I}(1) = 0$ is an absolute minimum with $\mathcal{I}''(1) > 0$ gives (87).

The following two lemmas help to consolidate error terms that appear in the proofs Theorems 1 and 2.

Lemma 18 There exists an absolute constant c such that for any t, q > 0 we have

$$|\mathcal{H}(a)| + |a|t/2 \le c(t+1)(\mathcal{I}(q)+1)$$
 (92)

where $a = a(q) = K^{-1}(\pi/(2q))$.

Proof Using (89) with the definition (11) we get that there is a constant c_2 so that

$$c_2^{-1}a(q) \le \mathcal{I}(q) \le c_2 a(q), \quad \text{if } q > 2,$$
 (93)

and the same bounds also give

$$|\mathcal{H}(a(q))| \le c_3 \sqrt{|a(q)|} \log |a(q)| \tag{94}$$



for some constant c_3 in the same region. This shows the existence of a constant A with

$$|\mathcal{H}(a)| \le A \mathcal{I}(q)$$
, and $a(q) \le A \mathcal{I}(q)$, for $q > 2$.

Since for 0 < q < 2 both a(q) and $\mathcal{H}(a(q))$ are bounded the lemma follows.

Lemma 19 For any $0 \le \varepsilon < 1/2$ there exists an absolute constant c, so that

$$\mathcal{I}(x+\varepsilon) < (1+\varepsilon)I(x) + c\varepsilon \tag{95}$$

Proof Since \mathcal{I} is convex, we have

$$\mathcal{I}(x+\varepsilon) < \mathcal{I}(x) + \varepsilon \mathcal{I}'(x+\epsilon).$$

Since $\mathcal{I}(x)$ is decreasing on $[0, \pi/2]$, the bound (95) follows immediately for $x \in [0, \pi/2 - \varepsilon]$ with any $c \ge 0$. Using (85) we get

$$\mathcal{I}(x+\varepsilon) \le \mathcal{I}(x) + \frac{\varepsilon}{x+\varepsilon} \left(\mathcal{I}(x+\varepsilon) - \frac{1}{8}K^{-1} \left(\frac{\pi}{2(x+\varepsilon)} \right) \right).$$
 (96)

From (90) it follows that there exists an $x_0 > 0$ such that

$$\frac{\varepsilon}{x+\varepsilon} \left(\mathcal{I}(x+\varepsilon) - \frac{1}{8} K^{-1} \left(\frac{\pi}{2(x+\varepsilon)} \right) \right) \le \varepsilon \mathcal{I}(x), \quad \text{if } x \ge x_0$$

uniformly in $\varepsilon \in [0, 1/2]$. Therefore, for $x > x_0$ we have that $\mathcal{I}(x + \varepsilon) \le (1 + \varepsilon)\mathcal{I}(x)$. We can assume $x_0 > \pi/2$. By choosing

$$c = \sup_{x \in [\pi/2, x_0 + 1/2]} \mathcal{I}'(x) = \mathcal{I}'(x_0 + 1/2)$$

we get $\mathcal{I}(x+\varepsilon) \leq \mathcal{I}(x) + c\varepsilon$ on $[\pi/2 - \varepsilon, x_0]$ with any $0 \leq \varepsilon < 1/2$ and the lemma follows.

Appendix B: Properties of I_{Sine}

In the final section of the appendix we describe the behavior of the function $I_{\text{Sine}}(\rho)$ near $\rho = \frac{1}{2\pi}$ and $\rho \to \infty$.

Proposition 20 The functions $\gamma(\nu)$ and $I_{Sine}(\rho)$ satisfy the following.

1. The function $\gamma(v)$ defined in (7) is continuous and strictly decreasing. It satisfies the differential equation $4x(1-x)\gamma''(x) = \gamma(x)$ on $(-\infty,0)$ and on (0,1) with boundary behavior $\lim_{x\to 0^{\pm}} \gamma(x) = \frac{1}{2\pi}$, $\gamma(1) = 0$ and $\lim_{x\to -\infty} \frac{\gamma(x)}{\sqrt{|x|}} = \frac{1}{4}$. These limits and the differential equation identify $\gamma(x)$ uniquely on $(-\infty,0) \cup (0,1]$.



2. We have $I_{Sine}(0) = \frac{1}{64}$, $I_{Sine}(\frac{1}{2\pi}) = 0$, and $I''_{Sine}(x) > 0$ for $x \neq \frac{1}{2\pi}$. Moreover, we have the following limits:

$$\lim_{x \to 0} \frac{I_{\operatorname{Sine}}(\frac{1}{2\pi} + x)}{\frac{x^2}{\log(1/|x|)}} = \frac{\pi^2}{4}, \quad and \quad \lim_{\rho \to \infty} \frac{I_{\operatorname{Sine}}(\rho)}{\rho^2 \log \rho} = \frac{1}{2}.$$

Proof Recall the function $\gamma(\nu)$ given in (7). Using the asymptotics (90) proved in Proposition 17 it is easy to see that $\gamma(\nu)$ is well-defined and positive in $(-\infty, 0) \cup (0, 1)$ with $\lim_{\nu \to 1^-} \gamma(\nu) = 0 = \gamma(1)$ and $\lim_{x \to -\infty} \frac{\gamma(x)}{\sqrt{|x|}} = \frac{1}{4}$. We also get that $\gamma(x)\mathcal{H}(x)$ blows up as $\nu \to 0^-$ or 0^+ .

Differentiating (7) and using the definition (5) lead to

$$\mathcal{H}(x)\gamma'(x) = \frac{1}{8} + \mathcal{H}'(x)\gamma(x), \qquad \frac{\gamma''(x)}{\gamma(x)} = \frac{\mathcal{H}''(x)}{\mathcal{H}(x)} = \frac{1}{4x(1-x)}, \tag{97}$$

for $x \in (-\infty, 0) \cup (0, 1)$. We have $\mathcal{H}'(x) = -\frac{K(x)}{2} < 0$ for x < 1 and $\mathcal{H}(0) = 0$. Thus from the second identity we get that $\gamma'(\nu)$ is strictly decreasing in $(-\infty, 0)$ and strictly increasing in (0, 1). From the asymptotics (90) of Proposition 17 it is not hard to check that $\lim_{\nu \to -\infty} \gamma'(\nu) = 0$ and $\lim_{\nu \to 1^-} \gamma'(\nu) = -\frac{1}{8}$. This, together with the previous statement, proves that $\gamma(\nu)$ is decreasing on $(-\infty, 0)$ and also on (0, 1].

Since $(\gamma(x)\mathcal{H}(x)^{-1})' = \frac{1}{8}\mathcal{H}(x)^{-2}$, L'Hospital's rule gives

$$\lim_{\nu \to 0} \gamma(\nu) = \lim_{\nu \to 0} \frac{\gamma(\nu)\mathcal{H}(\nu)^{-1}}{\mathcal{H}(\nu)^{-1}} = -\frac{1}{8}H'(0) = \frac{1}{2\pi} = \gamma(0).$$

Then from (97) it follows that

$$\lim_{\nu \to 0} \gamma''(\nu)\nu = \frac{1}{8\pi}, \quad \lim_{\nu \to 0} \frac{\gamma'(\nu)}{\log |\nu|} = \frac{1}{8\pi},$$

and also that

$$\lim_{x \to 0} \frac{\gamma^{(-1)}(\frac{1}{2\pi} + x)}{8\pi \frac{x}{\log|x|}} = 1. \tag{98}$$

This shows that $\gamma(\nu)$ is continuous and strictly decreasing on $(-\infty, 1]$. We have $\gamma^{(-1)}(1) = 0$ and $I_{\text{Sine}}(0) = \frac{1}{64}$.

Note, the fact that $\gamma(x)$ solves $4x(1-x)\gamma''(x) = \gamma(x)$ on $(-\infty,0) \cup (0,1)$ and has the proven asymptotics at $-\infty,0$ and 1 uniquely identifies it. The equation $4x(1-x)\gamma''(x) = y(x)$ has two linearly independent solutions on both $(-\infty,0)$ and (0,1). The function $\mathcal{H}(x)$ also solves the equation (on both intervals), but with $\mathcal{H}(0) = 0$, $\lim_{x \to 1^-} \mathcal{H}(x) = -1$ and $\lim_{x \to -\infty} \frac{\mathcal{H}(x)}{\sqrt{|x| \log |x|}} = \frac{1}{2}$. This shows that any solution on $(-\infty,0)$ or (0,1) can be expressed as $c_1\gamma + c_2\mathcal{H}$ with some constants c_1,c_2 , and the values of the constants are determined by the behavior of the solution at the end of the interval.



Using (97) together with (6) we can also compute that

$$\begin{split} I_{\mathrm{Sine}}'(\rho) &= \frac{1}{8} \left[\frac{1}{8\gamma'(\nu)} + \mathcal{H}(\nu) + \frac{\gamma(\nu)\mathcal{H}'(\nu)}{\gamma'(\nu)} \right] = \frac{1}{4}\mathcal{H}(\nu), \text{ and} \\ I_{\mathrm{Sine}}''(\rho) &= \frac{1}{4} \frac{\mathcal{H}'(\nu)}{\gamma'(\nu)} = -\frac{1}{8} \frac{K(\nu)}{\gamma'(\nu)}, \end{split}$$

where ν is short for $\gamma^{-1}(\rho)$. From this $I_{\text{Sine}}(\frac{1}{2\pi})=0$ follows, together with $I_{\text{Sine}}(x)>0$ for $x\neq\frac{1}{2\pi}$. The asymptotics of $I_{\text{Sine}}(\frac{1}{2\pi}+x)$ as $x\to 0$ can be obtained from the definition (6), the asymptotics (98), and the fact that $\mathcal{H}(0)=0$, $\mathcal{H}'(0)=-\frac{\pi}{4}$.

Lastly we can look at the asymptotics of $I_{\text{Sine}}(\rho)$ as $\rho \to \infty$. Recalling again (90) we get

$$\mathcal{H}(-x) \sim \frac{1}{2}\sqrt{x}\log x$$
, $\gamma(-x) \sim \frac{\sqrt{x}}{4}$, $\gamma^{(-1)}(x) \sim 16x^2$ as $x \to \infty$,

from which $I_{\text{Sine}}(\rho) \sim \frac{1}{2}\rho^2 \log \rho$ follows for $\rho \to \infty$.

References

- Abramowitz, M., Stegun, I.A. (eds.): Handbook of mathematical functions with formulas, graphs, and mathematical tables. Wiley, (1984)
- Anderson, G., Guionnet, A., Zeitouni, O.: Introduction to random matrices. Cambridge University Press, (2009)
- Basor, E.L., Tracy, C.A., Widom, H.: Asymptotics of level-spacing distributions for random matrices. Phys. Rev. Lett. 69(1), 5–8 (1992)
- Ben Arous, G., Guionnet, A.: Large deviations for Wigner's law and Voiculescu's non-commutative entropy. Probab. Theory Related Fields 108(4), 517–542 (1997)
- Buttazzo, G., Giaquinta, M., Hildebrant, S.: Oxford lecture series in mathematics and its applications. One-dimensional Variational Problems. Clarendon Press. Oxford (1998)
- Carlson, B.C., Gustafson, J.L.: Asymptotic expansion of the first elliptic integral. SIAM J. Math Anal. 16(5), 1072–1092 (1985)
- Deift, P., Its, A., Krasovsky, I., Zhou, X.: The Widom-Dyson constant for the gap probability in random matrix theory. J. Comput. Appl. Math. 202(1), 26–47 (2007)
- Deift, P., Its, A., Zhou, X.: A Riemann-Hilbert approach to asymptotic problems arising in the theory of random matrices and also in the theory of integrable statistical mechanics. Ann. Math. 146, 149–235 (1997)
- 9. Dembo, A., Zeitouni, O.: Large deviations techniques and applications, vol. 38. Springer, (2010)
- Dyson, F.J.: The Coulomb fluid and the fifth Painlevé transcendent. In: Liu, C.S., Yau, S.-T. (eds.) Chen Ning Yang, pp. 131–146. Int. Press, Cambridge (1995)
- Ehrhardt, T.: Dyson's constant in the asymptotics of the Fredholm determinant of the sine kernel. Comm. Math. Phys. 262(2), 317–341 (2006)
- Fogler, M.M., Shklovskii, B.I.: Probability of an eigenvalue number fluctuation in an interval of a random matrix spectrum. Phys. Rev. Lett. 74(17), 3312–3315 (1995)
- Forrester, P.J.: Log-gases and random matrices, volume 34 of London Mathematical Society Monographs Series. Princeton University Press, Princeton, NJ (2010)
- Forrester, P.J., Witte, N.S.: Asymptotic forms for hard and soft edge general î conditional gap probabilities. Nuclear Physics B 859(3), 321–340 (2012)
- Gustafson, J.L.: Asymptotic formulas for elliptic integrals. PhD thesis, Iowa State University, Ames, Iowa (1982)
- Karatzas, I., Ruf, J.: Distribution of the time to explosion for one-dimensional diffusions. http://arxiv. org/abs/1303.5899 (2013)



 Karatzas, I., Shreve, S.E.: Brownian motion and stochastic calculus, volume 113 of Graduate Texts in Mathematics, 2nd edn. Springer-Verlag, New York (1991)

- Krasovsky, I.V.: Gap probability in the spectrum of random matrices and asymptotics of polynomials orthogonal on an arc of the unit circle. Int. Math. Res. Not. 25, 1249–1272 (2004)
- Kritchevski, E., Valkó, B., Virág, B.: The scaling limit of the critical one-dimensional random Schrödinger operator. Comm. Math. Phys. 314(3), 775–806 (2012)
- Mehta, M.L.: Random matrices, volume 142 of Pure and Applied Mathematics (Amsterdam), 3rd edn. Elsevier/Academic Press, Amsterdam (2004)
- Rao, M.M., Ren, Z.D.: Theory of Orlicz spaces, volume 146 of Monographs and Textbooks in Pure and Applied Mathematics. Marcel Dekker Inc, New York (1991)
- Tracy, C.A., Widom, H.: Introduction to random matrices. In: Geometric and quantum aspects of integrable systems (Scheveningen, 1992), volume 424 of Lecture Notes in Physics, pp. 103–130. Springer, Berlin (1993)
- Valkó, B., Virág, B.: Continuum limits of random matrices and the Brownian carousel. Inventiones Math. 177, 463–508 (2009)
- 24. Valkó, B., Virág, B.: Large gaps between random eigenvalues. Ann. Probab. 38(3), 1263-1279 (2010)
- 25. Widom, H.: The asymptotics of a continuous analogue of orthogonal polynomials. J. Approx. Theory 77, 51–64 (1996)

