



# Central Spectral Gaps of the Almost Mathieu Operator

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**Abstract:** We consider the spectrum of the almost Mathieu operator  $H_\alpha$  with frequency  $\alpha$  and in the case of the critical coupling. Let an irrational  $\alpha$  be such that  $|\alpha - p_n/q_n| < cq_n^{-\varkappa}$ , where  $p_n/q_n, n = 1, 2, \dots$  are the convergents to  $\alpha$ , and  $c, \varkappa$  are positive absolute constants,  $\varkappa < 56$ . Assuming certain conditions on the parity of the coefficients of the continued fraction of  $\alpha$ , we show that the central gaps of  $H_{p_n/q_n}, n = 1, 2, \dots$ , are inherited as spectral gaps of  $H_\alpha$  of length at least  $c'q_n^{-\varkappa/2}, c' > 0$ .

## 1. Introduction

Let  $H_{\alpha,\theta}$  with  $\alpha, \theta \in (0, 1]$  be the self-adjoint operator acting on  $l^2(\mathbb{Z})$  as follows:

$$(H_{\alpha,\theta}\phi)(n) = \phi(n-1) + \phi(n+1) + 2\cos 2\pi(\alpha n + \theta)\phi(n), \quad n = \dots, -1, 0, 1, \dots \quad (1.1)$$

This operator is known as the almost Mathieu, Harper, or Azbel-Hofstadter operator. It is a one-dimensional discrete periodic (for  $\alpha$  rational) or quasiperiodic (for  $\alpha$  irrational) Schrödinger operator which models an electron on the 2-dimensional square lattice in a perpendicular magnetic field. Analysis of the spectrum of  $H_{\alpha,\theta}$  (and its natural generalization when the prefactor 2 of cosine, the coupling, is replaced by an arbitrary real number  $\lambda$ ) has been a subject of many investigations. In the present paper, we are concerned with the structure of the spectrum of  $H_{\alpha,\theta}$  as a set. Denote by  $a_j \in \mathbb{Z}_+, j = 1, 2, \dots$  (infinite sequence if  $\alpha$  is irrational) the coefficients of the continued fraction of  $\alpha$ :

$$\alpha = [a_1, a_2, \dots] = \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

If  $\alpha = p/q$  is rational, where  $p, q$  are coprime, i.e.  $(p, q) = 1$ , positive integers, there exists  $n$  such that

$$p/q = [a_1, a_2, \dots, a_n] = \frac{1}{a_1 + \frac{1}{a_2 + \dots + \frac{1}{a_n}}}$$

We denote by  $S(\alpha)$  the union of the spectra of  $H_{\alpha,\theta}$  over all  $\theta \in (0, 1]$ . (Note, however, that if  $\alpha$  is irrational, the spectrum of  $H_{\alpha,\theta}$  does not depend on  $\theta$ ). If  $\alpha = p/q$ ,  $S(p/q)$  consists of  $q$  bands separated by gaps. As shown by van Mouche [22] and by Choi, Elliott, and Yui [7], all the gaps (with the exception of the centermost gap when  $q$  is even) are open. Much effort was expended to prove the conjectures of [1,4] that if  $\alpha$  is irrational, the spectrum is a Cantor set. B ellissard and Simon proved in [6] that the spectrum of the generalized operator mentioned above is a Cantor set for an (unspecified) dense set of pairs  $(\alpha, \lambda)$  in  $\mathbb{R}^2$ . Helffer and Sj strand [11] proved the Cantor structure and provided an analysis of gaps in the case when all the coefficients  $a_j$ 's of  $\alpha$  are sufficiently large. Choi, Elliott, and Yui [7] showed that in the case of  $\alpha = p/q$ , each open gap is at least of width  $8^{-q}$  (this bound was improved in [3] to  $e^{-\varepsilon q}$  with any  $\varepsilon > 0$  for  $q$  sufficiently large) which, together with a continuity result implies that all admissible gaps are open (in particular, the spectrum is a Cantor set) if  $\alpha$  is a Liouvilian number whose convergents  $p/q$  satisfy  $|\alpha - p/q| < e^{-Cq}$ . Last [16] showed that  $S(\alpha)$  has Lebesgue measure zero (and hence, since  $S(\alpha)$  is closed and known not to contain isolated points, a Cantor set) for all  $\alpha = [a_1, a_2, \dots]$  such that the sequence  $\{a_j\}_{j=1}^\infty$  is unbounded. The set of such  $\alpha$ 's has full measure 1. On the other hand, it was shown by Puig [23] that in the generalized case  $\lambda \neq \pm 2, 0$ , the spectrum is a Cantor set for  $\alpha$  satisfying a Diophantine condition. Finally, Avila and Krikorian [2] completed the proof that the spectrum for  $\lambda = 2$  has zero measure, and hence a Cantor set, for all irrational  $\alpha$ 's; moreover, the proof of the fact that the spectrum is a Cantor set for all real  $\lambda \neq 0$  and irrational  $\alpha$  was completed by Avila and Jitomirskaya in [3]. The measure of the spectrum for any irrational  $\alpha$  and real  $\lambda$  is  $|4 - 2|\lambda||$ : in the case  $\lambda \neq \pm 2$ , proved for a.e.  $\alpha$  also in [16] and for all irrationals in [12]. Also available are bounds on the measure of the union of all gaps, see [8, 14, 17]. Furthermore, see [19] for a recent work on the Hausdorff dimension of the spectrum, and [20], on the question of whether all admissible gaps are open.

In order to have a quantitative description of the spectrum, one would like to know if the exponential  $e^{-\varepsilon q}$  estimates for the sizes of the individual gaps can be improved at least for some of the gaps.

In this paper we provide a power-law estimate  $Cq^{-\kappa}$ ,  $\kappa < 28$ , for the widths of central gaps of  $S(p/q)$ , i.e., the gaps around the centermost band (Theorem 3 below), on a parity condition for the coefficients  $a_k$  in  $p/q = [a_2, a_2, \dots, a_n]$ .

From this result we deduce that  $S(\alpha)$  has an infinite number of power-law bounded gaps for any irrational  $\alpha = [a_1, a_2, \dots]$  admitting a power-law approximation by its convergents  $p_n/q_n = [a_1, a_2, \dots, a_n]$  and with a parity condition on  $a_j$ 's (Theorem 4 below). These gaps are inherited from the central ones of  $S(p_n/q_n)$ ,  $n = 1, 2, \dots$

First, let  $\alpha = p/q$ ,  $(p, q) = 1$ . A standard object used for the analysis of  $H_{\alpha,\theta}$  is the discriminant

$$\sigma(E) = -\text{tr} \left\{ \begin{pmatrix} E - 2 \cos(2\pi p/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} E - 2 \cos(2\pi 2p/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \dots \begin{pmatrix} E - 2 \cos(2\pi qp/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \right\}, \tag{1.2}$$

a polynomial of degree  $q$  in  $E$  with the property that  $S(p/q)$  is the image of  $[-4, 4]$  under the inverse of the mapping  $\sigma(E)$ . The fact that  $S(p/q)$  consists of  $q$  bands separated by

$q - 1$  open gaps (except for the centermost empty gap for  $q$  even) means that all the zeros of  $\sigma(E)$  are simple, in all the maxima the value of  $\sigma(E)$  is strictly larger than 4, while in all the minima, strictly less than  $-4$  (except for  $E = 0$  for  $q$  even, where  $|\sigma(0)| = 4$  and the derivative  $\sigma'(0) = 0$ ). Note an important fact that  $\sigma(E) = (-1)^q \sigma(-E)$ , and hence  $S(p/q)$  is symmetric w.r.t.  $E = 0$ .

In what follows, we assume that  $q$  is odd. The case of even  $q$  can be considered similarly. Let us number the bands from left to right, from  $j = -(q - 1)/2$  to  $j = (q - 1)/2$ . Let  $\lambda_j$  denote the centers of the bands, i.e.  $\sigma(\lambda_j) = 0$ . Note that, by the symmetry of  $\sigma(E)$ ,  $\lambda_0 = 0$ . Let  $\mu_j$  and  $\eta_j$  denote the edges of the bands, i.e.  $|\sigma(\mu_j)| = |\sigma(\eta_j)| = 4$ , assigned as follows. If  $q = 4k + 3, k = 0, 1, \dots$ , we set  $\sigma(\mu_j) = 4, \sigma(\eta_j) = -4$  for all  $j$ . (In this case the derivative  $\sigma'(0) > 0$ , as follows from the fact that  $\sigma(E) < 0$  for all  $E$  sufficiently large.) If  $q = 4k + 1, k = 0, 1, \dots$ , we set  $\sigma(\mu_j) = -4, \sigma(\eta_j) = 4$  for all  $j$ . (In this case the derivative  $\sigma'(0) < 0$ .) Thus, in both cases, the bands are  $B_j = [\eta_j, \mu_j]$  for  $|j|$  even, and  $B_j = [\mu_j, \eta_j]$  for  $|j|$  odd.

Let  $w_j = \mu_j - \lambda_j, w'_j = \lambda_j - \eta_j$  for  $|j|$  even, and  $w_j = \eta_j - \lambda_j, w'_j = \lambda_j - \mu_j$  for  $|j|$  odd. Thus, the width of the  $j$ 's band is always  $w_j + w'_j$ . By the symmetry, for the centermost band  $B_0 = [\eta_0, \mu_0], w_0 = w'_0$ , and in general  $w_j = w'_{-j}$ .

For any real  $\alpha$ , denote the gaps of  $S(\alpha)$  by  $G_j(\alpha)$  and their length by  $\Delta_j(\alpha)$ . For  $\alpha = p/q$ , we order them in the natural way, namely,

$$G_j = (\mu_j, \mu_{j+1}), \quad \Delta_j = \mu_{j+1} - \mu_j, \quad \text{for } |j| \text{ even}, \tag{1.3}$$

$$G_j = (\eta_j, \eta_{j+1}), \quad \Delta_j = \eta_{j+1} - \eta_j, \quad \text{for } |j| \text{ odd}. \tag{1.4}$$

By the symmetry,  $\Delta_j = \Delta_{-j-1}$  for  $0 \leq j < (q - 1)/2$ .

In Sect. 2, we prove

**Lemma 1** (Comparison of the widths for the gaps and bands). *Let  $q \geq 3$  be odd. There hold the inequalities*

$$\Delta_0 > \left(\frac{w_0}{4}\right)^2, \quad \Delta_j > \frac{w_j^2}{4C_0^{2(j+1)}}, \quad 1 \leq j < \frac{q-1}{2}, \tag{1.5}$$

$$\Delta_j > \left(\frac{w_0}{8}\right)^{2j}, \quad 1 \leq j < \frac{q-1}{2}, \tag{1.6}$$

where  $C_0 = 1 + 2e/(\sqrt{5} - 1) = 5.398\dots$

*Remark.* The inequalities of Lemma 1 are better for small  $j$ , i.e., for central gaps and bands, which is the case we need below. For large  $j$ , note the following estimate which one can deduce using the technique of Last [16]:  $\Delta_j > \min\{w_j^2, w'_{j+1}{}^2\}/(4q), 0 \leq j < (q - 1)/2$ .

The inequality (1.6) gives us a lower bound for the width of the  $j$ 's gap provided an estimate for the width of the 0's band can be established. Such an estimate is given by

**Lemma 2** (Bound for the width of the centermost band). *Let  $q \geq 1, p/q = p_n/q_n = [a_1, a_2, \dots, a_n]$ , where  $a_1$  is odd and  $a_k, 2 \leq k \leq n$  are even. Then there exist absolute constants  $1 < C_1 < 14$  and  $1 < C_2 < e^{10}$  such that for the derivative of  $\sigma(E)$  at zero*

$$|\sigma'(0)| < C_2 q^{C_1}, \tag{1.7}$$

and half the width of the centermost band of  $S(p/q)$

$$w_0 \geq \frac{4}{|\sigma'(0)|} > 4C_2^{-1}q^{-C_1}. \tag{1.8}$$

If, in addition,  $q_{k+1} \geq q_k^\nu$ , for some  $\nu > 1$  and all  $1 \leq k \leq n - 1$ , then for any  $\varepsilon > 0$  there exists  $Q = Q(\varepsilon, \nu)$  such that if  $q > Q$ ,

$$|\sigma'(0)| < q^{5+\gamma_0+\varepsilon}, \quad w_0 > 4q^{-(5+\gamma_0+\varepsilon)}, \tag{1.9}$$

where  $\gamma_0$  is Euler's constant.

*Remark.* The bounds on  $C_1, C_2$  can be somewhat improved.

This lemma is proved in Sect. 3. The inequalities (1.5), (1.6), and especially (1.7) are the main technical results of this paper.

Combination of Lemmata 1 and 2 immediately yields

**Theorem 3** (Bound for the widths of the gaps). *Let  $q \geq 3$ ,  $p/q = [a_1, a_2, \dots, a_n]$ , where  $a_1$  is odd and  $a_k, 2 \leq k \leq n$ , are even. Then, with  $C_k, k = 1, 2$ , from Lemma 2, the width of the  $j$ 's gap of  $S(p/q)$  is*

$$\Delta_0 > \left(\frac{1}{C_2q^{C_1}}\right)^2, \quad \Delta_j > \left(\frac{1}{2C_2q^{C_1}}\right)^{2j}, \quad 1 \leq j < \frac{q-1}{2}. \tag{1.10}$$

*Remark.* The improvements for large  $q$  on the additional condition  $q_{k+1} > Cq_k^\nu$  are obvious from (1.9).

A consequence of this is the following theorem proved in Sect. 4.

**Theorem 4** *There exists an absolute  $C_3 > 0$  such that the following holds. Let  $\alpha = [a_1, a_2, \dots] \in (0, 1)$  be an irrational such that  $a_1$  is odd,  $a_k, k \geq 2$ , are even, and such that*

$$\left| \alpha - \frac{p_n}{q_n} \right| < \frac{1}{C_3q_n^\alpha}, \quad \alpha = 4C_1, \tag{1.11}$$

for all  $p_n/q_n = [a_1, a_2, \dots, a_n], n = 1, 2, \dots$ , where  $C_1$  is the constant from Lemma 2. Then

- (a) *The interior of the centermost band  $B_0$  of  $S(p_n/q_n)$  contains the centermost band and the closures of the gaps  $G_0, G_{-1}$  of  $S(p_{n+1}/q_{n+1}), n = 1, 2, \dots$*
- (b) *There exist distinct gaps  $G_{n,j}(\alpha), n = 1, 2, \dots, j = 1, 2$ , of  $S(\alpha)$ , such that the intersections  $G_{n,1}(\alpha) \cap G_{-1}(p_n/q_n), G_{n,2}(\alpha) \cap G_0(p_n/q_n), n = 1, 2, \dots$  are non-empty and the length of the gap  $G_{n,j}(\alpha)$*

$$\begin{aligned} \Delta_{n,j}(\alpha) &= |G_{n,j}(\alpha)| \geq |G_{n,j}(\alpha) \cap G_{j-2}(p_n/q_n)| \\ &> \frac{1}{C_4q_n^{\alpha/2}}, \quad n = 1, 2, \dots, \quad j = 1, 2, \end{aligned} \tag{1.12}$$

for some absolute  $C_4 > 0$ , where  $|A|$  denotes the Lebesgue measure of  $A$ .

- (c) *Let  $\varepsilon > 0$ , replace  $C_3$  by 2, and set  $\alpha = 4(5 + \gamma_0 + \varepsilon)$  in (1.11). Then there exists  $n_0 = n_0(\varepsilon)$  such that (a) and (b) hold for all  $n = n_0, n_0 + 1, \dots$  (instead of  $n = 1, 2, \dots$ ) with  $C_4$  replaced by 2, and with  $\alpha/2$  in (1.12) replaced by  $2(5 + \gamma_0) + \varepsilon$ .*

*Remarks.* (1) The statements (a), (b) of the theorem hold a fortiori for  $\varkappa = 4 \cdot 14 = 56$  and for any larger  $\varkappa$ . It is easy to provide explicit examples of irrationals satisfying the conditions of Theorem 4: take  $\varkappa = 56$ , any odd  $a_1$ , and even  $a_{n+1}$  such that  $a_{n+1} > C_3 q_n^{\varkappa-2}$ ,  $n \geq 1$ . Indeed, in this case,

$$\left| \alpha - \frac{p_n}{q_n} \right| < \frac{1}{q_n q_{n+1}} < \frac{1}{a_{n+1} q_n^2} < \frac{1}{C_3 q_n^\varkappa}.$$

- (2) Note that the parity condition on  $a_j$ 's implies, in particular, that all  $q_n$ 's are odd. This condition can be relaxed in all our statements. For example, we can allow a finite number of  $a_j$ 's to be odd at the expense of excluding some  $G(p_n/q_n)$ 's from the statement of Theorem 4 and worsening the bound on  $C_1$ . Note that in Lemma 2 we need  $q$  to be odd in order to use the estimate (1.7) on  $\sigma'(0)$  to obtain (1.8). One could obtain a bound on  $w_0$  for even  $q$  by providing an estimate on the second derivative  $\sigma''(0)$  in this case: for  $q$  even  $\sigma'(0) = 0$ . The parity condition we assume in this paper allows the best estimates and simplest proofs.
- (3) In Theorem 4, we only use Theorem 3 for  $j = 0$ , i.e., for the 2 centermost gaps. One can extend the result of Theorem 4, with appropriate changes, to more than 2 (at least a finite number) of central gaps of  $S(p_n/q_n)$ .
- (4) We can take  $C_3 = 4^2 60^2 C_2^4$ ,  $C_4 = 2C_2^2$ , in terms of the constant  $C_2$  from Lemma 2.
- (5) The statement (a) of the theorem holds already for  $\varkappa = 2C_1$ .

**2. Proof of Lemma 1**

Assume that  $q = 4k + 3$ ,  $k = 0, 1, \dots$  (A proof in the case  $q = 4k + 1$  is almost identical.) Let

$$s = \frac{q - 1}{2}.$$

In our notation, we can write

$$\begin{aligned} \sigma(E) &= \prod_{k=-s}^s (E - \lambda_k), & \sigma(E) - 4 &= \prod_{k=-s}^s (E - \mu_k), \\ \sigma(E) + 4 &= \prod_{k=-s}^s (E - \eta_k). \end{aligned} \tag{2.1}$$

Setting in the last 2 equations  $E = \lambda_j$ , we obtain the useful identities

$$4 = \prod_{k=-s}^s |\lambda_j - \mu_k|, \quad 4 = \prod_{k=-s}^s |\lambda_j - \eta_k|, \quad -s \leq j \leq s. \tag{2.2}$$

Fix  $0 \leq j \leq s$  (by the symmetry of the spectrum, it is sufficient to consider only nonnegative  $j$ ). It was shown by Choi, Elliott, and Yui [7] that

$$\prod_{k \neq j} |\mu_j - \mu_k| \geq 1, \quad \prod_{k \neq j} |\eta_j - \eta_k| \geq 1. \tag{2.3}$$

For simplicity of notation, we assume from now on that  $j < s - 1$ : the extension to  $j = s - 1$  is obvious. Let  $j \geq 0$  be even. By the first inequality in (2.3), we can write

$$\begin{aligned}
 1 &\leq |\sigma'(\mu_j)| = \prod_{k \neq j} |\mu_j - \mu_k| \\
 &= |\mu_j - \mu_{j+1}| \frac{\prod_{k=-s}^s |\lambda_j - \mu_k|}{|\lambda_j - \mu_j| |\lambda_j - \mu_{j+1}|} \prod_{k=-s}^{j-1} \left| 1 + \frac{\mu_j - \lambda_j}{\lambda_j - \mu_k} \right| \prod_{k=j+2}^s \left| 1 - \frac{\mu_j - \lambda_j}{\mu_k - \lambda_j} \right|.
 \end{aligned}
 \tag{2.4}$$

According to our notation,  $\mu_{j+1} - \mu_j = \Delta_j$ ,  $\mu_j - \lambda_j = w_j$ ,  $\mu_{j+1} - \lambda_j = w_j + \Delta_j$ . Recalling the first identity in (2.2) and rearranging the last product in (2.4), we continue (2.4) as follows

$$= \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\mu_k - \mu_j} \right|} < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\mu_k - \lambda_j} \right|},
 \tag{2.5}$$

because  $\mu_j > \lambda_j$ .

Now note that, by the symmetry of the spectrum,

$$|\mu_{j+\ell} - \lambda_j| < |\mu_{-j-\ell-1} - \lambda_j|, \quad \ell = 2, 3, \dots
 \tag{2.6}$$

Therefore, the r.h.s. of (2.5) is

$$< \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \prod_{k=-j-2}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right| \frac{1}{1 + \frac{w_j}{\mu_s - \lambda_j}}.
 \tag{2.7}$$

In the case  $j = 0$ , we now use the symmetry

$$w_0 = w'_0 < \lambda_0 - \mu_{-k}, \quad k \geq 1,
 \tag{2.8}$$

to obtain from (2.7)

$$1 < \frac{16\Delta_0}{w_0(w_0 + \Delta_0)} < \frac{16\Delta_0}{w_0^2},$$

which gives the first inequality in (1.5).

In general, however, we need to compare  $w_j$  and  $w'_j$  to estimate (2.7). According to equations (3.11), (3.12) of Last [16],

$$w_j, w'_j < e\ell_j, \quad \ell_j = \frac{4}{|\sigma'(\lambda_j)|},
 \tag{2.9}$$

and further, by equations (3.27), (3.28) of [16],

$$\frac{\sqrt{5} - 1}{2} \ell_j < w_j, w'_j.
 \tag{2.10}$$

(In fact, more is shown in [16]: for each pair of widths  $w_j, w'_j$ , at least one of them is larger than  $\ell_j$ .)

Therefore,

$$w_j < c_1 w'_j, \quad c_1 = \frac{2e}{\sqrt{5}-1}, \quad 0 \leq j \leq s. \tag{2.11}$$

Furthermore, it is obvious that

$$\frac{w'_j}{\lambda_j - \mu_k} < 1, \quad k = -j - 2, \dots, j - 1. \tag{2.12}$$

Therefore, we have for the product in (2.7):

$$\prod_{k=-j-2}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right| < (1 + c_1)^{2(j+1)}, \tag{2.13}$$

and since  $\mu_s - \lambda_j > 0$ , (2.7) finally gives

$$1 < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} (1 + c_1)^{2(j+1)}, \tag{2.14}$$

from which the inequality (1.5) with even  $j$  easily follows.

*Remark.* Last's equation (2.9) together with the Last–Wilkinson formula [16, 18]

$$\sum_{j=-s}^s |\sigma'(\lambda_j)|^{-1} = 1/q \tag{2.15}$$

implies [16] that the measure of the spectrum  $S(p/q)$  is at most  $8e/q$  and that for any  $j$ ,

$$w_j < 4e/q. \tag{2.16}$$

Now consider  $j$  odd,  $0 < j < s$ . Using the second inequalities in (2.3) and (2.2), we obtain similarly to (2.5),

$$1 < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \eta_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\eta_k - \lambda_j} \right|}, \tag{2.17}$$

and since

$$|\eta_{j+\ell} - \lambda_j| < |\eta_{-j-\ell-1} - \lambda_j|, \quad \ell = 2, 3, \dots, \tag{2.18}$$

we obtain the inequality (1.5) for  $j$  odd in a similar way.

Let again  $j$  be even,  $0 < j < s - 1$ . In order to compare  $\Delta_j$  with the width of the centermost band and, thus, obtain (1.6), we write instead of (2.4) the following:

$$1 \leq |\sigma'(\mu_j)| = \prod_{k \neq j} |\mu_j - \mu_k| \tag{2.19}$$

$$= |\mu_j - \mu_{j+1}| \frac{\prod_{k=-s}^s |\lambda_0 - \mu_k|}{|\lambda_0 - \mu_j| |\lambda_0 - \mu_{j+1}|} \prod_{k=-s}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \prod_{k=j+2}^s \left| 1 - \frac{\mu_j - \lambda_0}{\mu_k - \lambda_0} \right|. \tag{2.20}$$

Proceeding in a similar way as before, and using the inequalities

$$|\mu_{j+\ell} - \lambda_0| < |\mu_{-j-\ell-1} - \lambda_0|, \quad \ell = 2, 3, \dots, \quad (2.21)$$

we obtain

$$\begin{aligned} 1 &< \frac{4\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \prod_{k=-j-2}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \\ &< \frac{16\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \prod_{k=-j}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \\ &= \frac{16\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \left| \frac{\mu_j - \mu_0}{\lambda_0 - \mu_0} \right| \left| \frac{\mu_j - \mu_1}{\lambda_0 - \mu_1} \right| \prod_{\substack{k=-j \\ k \neq 0,1}}^{j-1} \left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right| \\ &< \frac{16\Delta_j}{|\lambda_0 - \mu_0||\lambda_0 - \mu_1|} \prod_{\substack{k=-j \\ k \neq 0,1}}^{j-1} \left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right|, \end{aligned} \quad (2.22)$$

and since (note that  $S(\alpha) \in [-4, 4]$ )

$$\left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right| < \frac{8}{w_0},$$

we obtain

$$1 < \frac{\Delta_j}{4} \left( \frac{8}{w_0} \right)^{2j}, \quad (2.23)$$

which gives an inequality slightly better than (1.6) for  $j$  even. Finally, we establish (1.6) for  $j$  odd by starting (instead of (2.19)) with the inequality  $1 \leq |\sigma'(\eta_{j+1})| = \prod_{k \neq j+1} |\eta_{j+1} - \eta_k|$  and arguing similarly.

*Remark.* Using the Last estimate (2.9)

$$\frac{4e}{w_j} > |\sigma'(\lambda_j)| = \prod_{k \neq j} |\lambda_j - \lambda_k|, \quad (2.24)$$

one can establish, in a way similar to the argument above, inequalities of the type

$$\Delta_j + w'_{j+1} > \frac{w_j}{C^j}, \quad (2.25)$$

with some absolute constant  $C > 0$ .

□



### 3. Proof of Lemma 2

As noted in a remark following Theorem 4, the parity conditions imposed on  $p/q$  in Lemma 2 imply, in particular, that  $q$  is odd. It follows from the symmetry of the discriminant  $\sigma(E) = -\sigma(-E)$  in this case that the maximum of the absolute value of the derivative  $\sigma'(E)$  in the  $j = 0$  band is at  $E = 0$ . Therefore,

$$w_0 \geq \frac{4}{|\sigma'(0)|} \tag{3.1}$$

(with the equality only for  $q = 1$ ), and hence, in order to prove Lemma 2, it remains to obtain the inequalities (1.7) and (1.9).

If  $q = 1$ , we have  $\sigma(E) = -E$ , and the result is trivial. Assume now that  $q$  is any (even or odd) integer larger than 1. We start with the following representation of  $\sigma(E)$  in terms of a  $q \times q$  Jacobi matrix with the zero main diagonal:

$$\sigma(E) = \det(\widehat{H} - EI), \tag{3.2}$$

where  $I$  is the identity matrix, and  $\widehat{H}$  is a  $q \times q$  matrix  $\widehat{H}_{jk}$ ,  $j, k = 1, \dots, q$ , where

$$\widehat{H}_{j,j+1} = \widehat{H}_{j+1,j} = 2 \sin\left(\pi \frac{p}{q} j\right), \quad j = 1, \dots, q - 1, \tag{3.3}$$

and the rest of the matrix elements are zero. For a proof, see e.g. the appendix of [15]. (This is related to a matrix representation for the almost Mathieu operator corresponding to the chiral gauge of the magnetic field potential, noticed by several authors [13, 21, 25].) The absence of the main diagonal in  $\widehat{H}$  allows us to obtain a simple expression for the derivative  $\sigma'(E)$  at  $E = 0$ . If  $q$  is even, it is easily seen that  $\sigma'(E) = 0$ . If  $q$  is odd, we denote  $s = (q - 1)/2$  and immediately obtain from (3.2) (henceforth we set  $\prod_{j=a}^b \equiv 1$  and  $\sum_{j=a}^b \equiv 0$  if  $a > b$ ):

$$\sigma'(0) = (-1)^s \sum_{k=0}^s \left[ \prod_{j=1}^k 2 \sin \frac{\pi p}{q} (2j - 1) \prod_{j=k+1}^s 2 \sin \frac{\pi p}{q} 2j \right]^2. \tag{3.4}$$

From now on, we assume that  $q \geq 3$  is odd unless stated otherwise.

*Remark.* Using the identity  $\prod_{j=1}^{(q-1)/2} 2 \sin \frac{\pi p}{q} 2j = q$ , we can represent (3.4) in the form

$$\sigma'(0) = (-1)^s q \left( 1 + \sum_{k=1}^s \prod_{j=1}^k \frac{\sin^2 \frac{\pi p}{q} (2j - 1)}{\sin^2 \frac{\pi p}{q} 2j} \right), \tag{3.5}$$

which exhibits the fact that  $|\sigma'(0)| > q$ . This is in accordance with the Last–Wilkinson formula (2.15).

Thus we have

$$|\sigma'(0)| = \sum_{k=0}^s \exp\{L_k\}, \tag{3.6}$$

where

$$\frac{1}{2}L_k = \sum_{j=1}^k \ln \left| 2 \sin \frac{\pi p}{q} (2j - 1) \right| + \sum_{j=k+1}^s \ln \left| 2 \sin \frac{\pi p}{q} 2j \right| \tag{3.7}$$

$$= \sum_{j=-s+k+1}^k \ln \left| 2 \sin \frac{\pi p}{q} (2j - 1) \right| = \sum_{j=1}^s \ln \left| 2 \sin \frac{\pi p}{q} 2(j + k) \right|. \tag{3.8}$$

Here we changed the summation variable  $j = s + j'$  in the second sum in (3.7) to obtain the first equation in (3.8), and then changed the variable  $j = k - s + j'$  to obtain the final equation in (3.8).

We will now analyze  $L_k$ . Using the Fourier expansion, we can write

$$L_k = -2 \sum_{j=1}^s \sum_{n=1}^{\infty} \frac{1}{n} \cos 4n \frac{\pi p}{q} (j + k). \tag{3.9}$$

Representing  $n$  in the form  $n = qm + \ell$ , where  $\ell = 1, 2, \dots, q - 1$  for  $m = 0$ , and  $\ell = 0, 1, \dots, q - 1$  for  $m = 1, 2, \dots$ , we have

$$L_k = - \sum_{m=1}^{\infty} \left( \frac{q - 1}{qm} - \frac{1}{q} \sum_{\ell=1}^{q-1} \frac{1}{m + \ell/q} F(\ell, k) \right) + S_k, \tag{3.10}$$

where

$$S_k = \sum_{\ell=1}^{q-1} \frac{1}{\ell} F(\ell, k), \tag{3.11}$$

and

$$\begin{aligned} F(\ell, k) &= -2 \sum_{j=1}^s \cos 4\ell \frac{\pi p}{q} (j + k) \\ &= -2\Re \sum_{j=1}^s \exp\{4\pi i \frac{p}{q} (j + k)\ell\} = \frac{\cos \pi \frac{p}{q} \ell(4k + 1)}{\cos \pi \frac{p}{q} \ell}. \end{aligned} \tag{3.12}$$

Note that since  $1 < j + k \leq q - 1$ ,  $(p, q) = 1$ , and  $q$  is odd, we have that  $(2j + 2k)p \not\equiv 0 \pmod{q}$ .

For  $0 \leq k \leq s$ ,

$$\sum_{\ell=1}^{q-1} F(\ell, k) = -2\Re \left( \sum_{j=1}^s \sum_{\ell=1}^{q-1} \exp\{4\pi i \frac{p}{q} (j + k)\ell\} + 1 - 1 \right) = q - 1. \tag{3.13}$$

We will use this fact later on.

Recall that

$$\sum_{m=1}^M \frac{1}{m} = \ln M + \gamma_0 + o(1), \quad M \rightarrow \infty, \tag{3.14}$$

where  $\gamma_0 = 0.5772\dots$  is Euler’s constant. Recall also Euler’s  $\psi$ -function

$$\psi(x + 1) = \frac{\Gamma'(x + 1)}{\Gamma(x + 1)} = -\gamma_0 - \sum_{m=1}^{\infty} \left( \frac{1}{m + x} - \frac{1}{m} \right), \quad x \geq 0. \tag{3.15}$$

The function  $\psi(x)$  continues to a meromorphic function in the complex plane with first-order poles at nonpositive integers  $x = 0, -1, -2, \dots$ . The function  $\psi(x)$  satisfies the equation

$$\psi(x + 1) = \psi(x) + \frac{1}{x}.$$

Expressions (3.14), (3.15) imply

$$\begin{aligned} \sum_{m=1}^M \frac{1}{m + x} &= \sum_{m=1}^M \left( \frac{1}{m + x} - \frac{1}{m} \right) + \sum_{m=1}^M \frac{1}{m} \\ &= \ln M - \psi(x + 1) + o(1), \quad M \rightarrow \infty, \end{aligned} \tag{3.16}$$

uniformly for  $x \in [0, 1]$ , in particular. We now rewrite (3.10) in the form

$$L_k = - \lim_{M \rightarrow \infty} \left[ \frac{q - 1}{q} \sum_{m=1}^M \frac{1}{m} - \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \sum_{m=1}^M \frac{1}{m + \ell/q} \right] + S_k. \tag{3.17}$$

Substituting here (3.14), (3.16), and using (3.13), we finally obtain

$$L_k = -\frac{q - 1}{q} \gamma_0 - \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \psi(1 + \ell/q) + S_k. \tag{3.18}$$

We will now provide an upper bound for the absolute values of sums in this expression. First, note that the derivative  $\psi'(x) \geq 0, x \in [1, 2]$ , and  $\psi(1) = -\gamma_0, \psi(2) = \psi(1) + 1 = 1 - \gamma_0$ . Therefore,

$$\max_{x \in [1, 2]} |\psi(x)| = \max\{|\psi(1)|, |\psi(2)|\} = \gamma_0.$$

Thus, for the first sum in the r.h.s. of (3.18) we have

$$\begin{aligned} \left| \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \psi(1 + \ell/q) \right| &\leq \frac{\gamma_0}{q} \sum_{\ell=1}^{q-1} |F(\ell, k)| = \frac{\gamma_0}{q} \sum_{\ell=1}^{q-1} \frac{1}{|\cos(\pi p \ell/q)|} \\ &= \frac{4\gamma_0}{q} \sum_{m=0}^{s-1} \frac{1}{|1 - e^{i(2m+1)\pi/q}|} < \gamma_0 \sum_{m=0}^{s-1} \frac{1}{m + 1/2} \\ &< \gamma_0 \left( 2 + \int_0^{s-1} \frac{dx}{x + 1/2} \right) < \gamma_0 (\ln q + 2), \quad q \geq 3. \end{aligned} \tag{3.19}$$

We will need a more subtle estimate for

$$S_k = \sum_{m=1}^{q-1} \frac{1}{m} F(m, k) = \sum_{m=1}^{q-1} \frac{1}{m} \frac{\cos \pi \frac{p}{q} m (4k + 1)}{\cos \pi \frac{p}{q} m} \tag{3.20}$$

in the r.h.s. of (3.18). We follow a method of Hardy and Littlewood [9] (see also [24]). It relies on a recursive application of a suitably constructed contour integral.

For  $q \geq 3$  odd,  $(p, q) = 1$ , let

$$I(p/q, \gamma) = -2 \int_{\Gamma_q} \frac{e^{(1+p/q)z}}{(1 + e^{z p/q})(1 - e^z)} \frac{e^{-\gamma z}}{z} dz, \quad \frac{1}{2} \frac{p}{q} \leq \gamma \leq 1 + \frac{1}{2} \frac{p}{q}, \tag{3.21}$$

where the contour  $\Gamma_q$  are the 2 direct lines parallel to the real axis given by: (1)  $\pi i/2 + x$ ,  $x \in \mathbb{R}$ , oriented from  $-\infty$  to  $+\infty$ ; (2)  $2\pi i(q - 1/4) + x$ ,  $x \in \mathbb{R}$ , oriented from  $+\infty$  to  $-\infty$ . Note that the choice of  $\gamma$  in (3.21) ensures that the integral converges both at  $+\infty$  and  $-\infty$ .

Now again for  $q \geq 3$  odd,  $(p, q) = 1$ , let

$$S(p/q, \gamma) = \sum_{m=1}^{q-1} \frac{e^{\pi i \frac{p}{q} m - 2\pi i \gamma m}}{m \cos(\pi \frac{p}{q} m)}. \tag{3.22}$$

Note that the sum (3.20)

$$S_k = \Re S(p/q, -2kp/q), \tag{3.23}$$

and that the denominators in (3.22) are nonzero.

We will also need the following auxiliary sum,  $(p, q) = 1$ ,

$$T(p/q, \gamma, \delta) = 2 \sum_{n=1}^q \frac{(-1)^\delta e^{2\pi i \frac{p}{q} (n - \frac{1}{2})} e^{-2\pi i \gamma (n - \frac{1}{2})}}{1 - (-1)^\delta e^{2\pi i \frac{p}{q} (n - \frac{1}{2})} n - \frac{1}{2}}, \quad \delta = 0, 1, \tag{3.24}$$

where we assume that  $p$  is odd if  $\delta = 0$  and that  $p$  and  $q$  have opposite parities if  $\delta = 1$ . These conditions imply that  $p(2n - 1) \neq q(2m - \delta)$ ,  $m, n \in \mathbb{Z}$ , and therefore the denominators in (3.24) are nonzero.

With this notation we have

**Lemma 5.** *Let  $q$  be odd,  $(p, q) = 1$ ,  $p > 0$ ,  $q > 1$ . Let  $p/q$  have the following continued fraction:*

$$\frac{p}{q} = \frac{1}{a + \frac{p'}{q'}}, \quad p' \geq 0, \quad q' > p'. \tag{3.25}$$

Then

$$I(p/q, \gamma) = S(p/q, \gamma) - (-1)^{\varepsilon'} T(p'/q', \gamma', a \bmod 2), \quad \gamma' = \frac{q}{p} \gamma \pmod{1}, \tag{3.26}$$

where  $\varepsilon' = 0$  if  $\gamma' = \frac{q}{p} \gamma \pmod{2}$ , and  $\varepsilon' = 1$  otherwise.

Moreover, there holds the bound

$$|I(p/q, \gamma)| < 4 \ln \frac{q}{p} + \frac{5}{e\pi p} + \beta, \quad \beta = 4(e^{-1} + \operatorname{arcsinh}(4/\pi)) = 5.719 \dots \tag{3.27}$$

*Remarks.*(1) One can take any  $\gamma'$  satisfying the congruence in (3.26).

(2) The bound in (3.27) can be somewhat decreased by improving (3.33), (3.34) below. Similar can be achieved in (3.38) below.

*Proof.* Consider the integral  $I_\Gamma(p/q, \gamma)$ , which has the same integrand as in (3.21), but with integration over some contour  $\Gamma$ . Let  $\Gamma_{q,\xi}$  be the following quadrangle traversed in the positive direction:  $\Gamma_{q,\xi} = \cup_{j=1}^4 \Gamma_{q,\xi}^{(j)}$ ,  $\xi > 0$ , where  $\Gamma_{q,\xi}^{(1)} = [\pi i/2 - \xi, \pi i/2 + \xi]$ ,  $\Gamma_{q,\xi}^{(2)} = [\pi i/2 + \xi, 2\pi i(q - 1/4) + \xi]$ ,  $\Gamma_{q,\xi}^{(3)} = [2\pi i(q - 1/4) - \xi, 2\pi i(q - 1/4) + \xi]$ ,  $\Gamma_{q,\xi}^{(4)} = [2\pi i(q - 1/4) - \xi, \pi i/2 - \xi]$ .

Recalling the conditions on  $\gamma$  in (3.21), we first note that on the vertical segments, for some constants  $C$  which may depend on  $p, q$ ,

$$|I_{\Gamma_{q,\xi}^{(2)}}(p/q, \gamma)| \leq C \frac{e^{-\gamma\xi}}{\xi} \leq C \frac{e^{-\frac{p}{2q}\xi}}{\xi} \rightarrow 0, \quad \text{as } \xi \rightarrow \infty, \tag{3.28}$$

$$|I_{\Gamma_{q,\xi}^{(4)}}(p/q, \gamma)| \leq C \frac{e^{-(1+p/q-\gamma)\xi}}{\xi} \leq C \frac{e^{-\frac{p}{2q}\xi}}{\xi} \rightarrow 0, \quad \text{as } \xi \rightarrow \infty. \tag{3.29}$$

and we conclude that

$$I(p/q, \gamma) = \lim_{\xi \rightarrow \infty} I_{\Gamma_{q,\xi}}(p/q, \gamma). \tag{3.30}$$

On the other hand,  $I_{\Gamma_{q,\xi}}(p/q, \gamma)$  is given by the sum of residues inside the contour. Clearly, the integrand has poles there at the points:

$$z_m = 2\pi im, \quad m = 1, \dots, q - 1 \tag{3.31}$$

$$\tilde{z}_n = 2\pi i \frac{q}{p}(n - 1/2), \quad n = 1, \dots, p. \tag{3.32}$$

Note that for all these  $m, n$ ,  $z_m \neq \tilde{z}_n$  because  $m \neq \frac{q}{p} \frac{2n-1}{2}$  as  $q$  is odd. Hence we conclude that all the poles inside  $\Gamma_{q,\xi}$  are simple. Computing the residues and using the facts that

$$\frac{q}{p} = a + \frac{p'}{q'}, \quad q' = p,$$

we obtain (3.26) by (3.30). Note that the conditions on  $p', q'$  in  $T(p'/q')$  are fulfilled since  $q$  is odd.

Now, in order to obtain the inequality (3.27), we evaluate the integral along the contour  $\Gamma_q$ . On the lower part of it,

$$\begin{aligned} |I_{\frac{\pi i}{2} + \mathbb{R}}(p/q, \gamma)| &\leq 2 \int_0^\infty \frac{e^{-\gamma x} + e^{-(1+\frac{p}{q}-\gamma)x}}{(1 + 2e^{-\frac{p}{q}x} \cos \frac{\pi p}{2q} + e^{-2\frac{p}{q}x})^{1/2} (1 + e^{-2x})^{1/2} (x^2 + \frac{\pi^2}{4})^{1/2}} dx \\ &< 4 \int_0^\infty \frac{e^{-\frac{p}{2q}x}}{(x^2 + \frac{\pi^2}{4})^{1/2}} dx = 4 \int_0^\infty \frac{e^{-u}}{(u^2 + (\frac{\pi p}{4q})^2)^{1/2}} dx. \end{aligned} \tag{3.33}$$

Separating the final integral into 2 parts, one along  $(0, 1)$  and another along  $(1, \infty)$ , we can continue (3.33) as follows:

$$\begin{aligned}
 &< 4 \left( \int_0^1 \frac{du}{(u^2 + (\frac{\pi p}{4q})^2)^{1/2}} + \frac{1}{(1 + (\frac{\pi p}{4q})^2)^{1/2}} \int_1^\infty e^{-u} du \right) \\
 &= 4 \left( -\ln \frac{\pi p}{4q} + \ln \left[ 1 + \sqrt{1 + \left(\frac{\pi p}{4q}\right)^2} \right] + \frac{e^{-1}}{(1 + (\frac{\pi p}{4q})^2)^{1/2}} \right) \\
 &< 4 \left( \ln \frac{q}{p} + \ln \left[ \frac{4}{\pi} + \sqrt{\left(\frac{4}{\pi}\right)^2 + 1} \right] + e^{-1} \right). \tag{3.34}
 \end{aligned}$$

Similarly, we obtain (recall that  $q > 1$ )

$$\begin{aligned}
 |I_{2\pi i(q-\frac{1}{4})+\mathbb{R}}(p/q, \gamma)| &\leq 4 \int_0^\infty \frac{e^{-u}}{(u^2 + (\pi(p - \frac{p}{4q}))^2)^{1/2}} dx \\
 &< \frac{4e^{-1}}{\pi(p - \frac{p}{4q})} < \frac{4e^{-1}}{\pi p} \frac{1}{1 - \frac{1}{4q}} < \frac{5}{e\pi p}. \tag{3.35}
 \end{aligned}$$

The sum of (3.34) and (3.35) gives (3.27), and thus we finish the proof of Lemma 5.  $\square$

We will need another similar lemma.

**Lemma 6.** *Let  $(p, q) = 1$ ,  $p > 0$ ,  $q > 1$ ,*

$$\begin{aligned}
 J(p/q, \gamma, \delta) &= 2 \int_{\Gamma_q} \frac{(-1)^\delta e^{(1+p/q)z}}{(1 - (-1)^\delta e^{zp/q})(1 + e^z)} \frac{e^{-\gamma z}}{z} dz, \\
 \frac{1}{2} \frac{p}{q} &\leq \gamma \leq 1 + \frac{1}{2} \frac{p}{q}, \tag{3.36}
 \end{aligned}$$

where  $\delta = \{0, 1\}$ , and  $\Gamma_q$  is the same contour as in (3.21). Assume that  $p$  is odd if  $\delta = 0$ , and either  $(p - \text{even}, q - \text{odd})$ , or  $(p - \text{odd}, q - \text{even})$  if  $\delta = 1$ . Let  $p/q$  have the continued fraction (3.25). Then

$$\begin{aligned}
 &J(p/q, \gamma, \delta) \\
 &= T(p/q, \gamma, \delta) + \begin{cases} -S(p'/q', \gamma'), & \text{if } \delta = 0 \\ (-1)^{\varepsilon'} T(p'/q', \gamma', a + 1 \bmod 2), & \text{if } \delta = 1 \end{cases}, \quad \gamma' = \frac{q}{p} \gamma \pmod{1}, \tag{3.37}
 \end{aligned}$$

where  $\varepsilon' = 0$  if  $\gamma' = \frac{q}{p} \gamma \pmod{2}$ , and  $\varepsilon' = 1$  otherwise.

Moreover, there holds the bound with  $\beta$  from (3.27)

$$\begin{aligned}
 |J(p/q, \gamma, \delta)| &< A_\delta \left( 4 \ln \frac{q}{p} + \frac{5}{e\pi p} + \beta \right), \\
 A_\delta &= \begin{cases} (1 - \cos^2 \pi \frac{p}{q})^{-1/2}, & \text{if } \delta = 0 \\ 1, & \text{if } \delta = 1 \end{cases} \tag{3.38}
 \end{aligned}$$

*Proof.* We argue as in the proof of Lemma 5. The poles of the integrand in (3.36) inside  $\Gamma_{q,\xi}(p/q, \gamma)$  are:

$$z_m = 2\pi i(m - 1/2), \quad m = 1, \dots, q, \tag{3.39}$$

$$\tilde{z}_n = 2\pi i \frac{q}{p}(n - \delta/2), \tag{3.40}$$

and  $n = 1, \dots, p - 1$  if  $\delta = 0$ , while  $n = 1, \dots, p$  if  $\delta = 1$ . Our assumptions on the parity of  $p, q$  immediately imply that all  $z_n \neq \tilde{z}_m$  and hence all the poles inside  $\Gamma_{q,\xi}(p/q, \gamma)$  are simple. Computing the residues we obtain (3.37).

Denote by  $J_\Gamma(p/q, \gamma, \delta)$  the integral which has the same integrand as in (3.36), but with integration over some contour  $\Gamma$ . As in the previous proof, consider now an estimate for the integral along the lower part of  $\Gamma_q$ :

$$\begin{aligned} |J_{\frac{\pi i}{2} + \mathbb{R}}(p/q, \gamma)| &\leq 2 \int_0^\infty \frac{e^{-\gamma x} + e^{-(1+\frac{p}{q}-\gamma)x}}{(1 - 2(-1)^\delta e^{-\frac{p}{q}x} \cos \frac{\pi p}{2q} + e^{-2\frac{p}{q}x})^{1/2} (1 + e^{-2x})^{1/2} (x^2 + \frac{\pi^2}{4})^{1/2}} dx \\ &< 4A_\delta \int_0^\infty \frac{e^{-\frac{p}{2q}x}}{(x^2 + \frac{\pi^2}{4})^{1/2}} dx, \end{aligned} \tag{3.41}$$

where  $A_\delta$  is given in (3.38). The rest of the argument is very similar to that in the proof of Lemma 5. □

**Lemma 7** (recurrence). *Let  $q > 1, \frac{p}{q} = \frac{p_n}{q_n} = [a_1, \dots, a_n]$  and denote  $t_j = [a_j, \dots, a_n], j = 1, \dots, n$ . Assume that  $a_1$  is odd, and  $a_j$  are even for  $j = 2, \dots, n$ . Fix  $0 \leq k \leq (q - 1)/2$ . Then*

$$S(p/q, \gamma_1) = I(t_1, \gamma_1) + \sum_{j=2}^n (-1)^{\varepsilon_j} J(t_j, \gamma_j, 1) + 2(-1)^{\varepsilon_n} e^{-i\pi\gamma_n a_n} \tag{3.42}$$

where  $\varepsilon_j = \{0, 1\}$  and

$$\gamma_1 = -2kt_1 + k_1, \quad \gamma_j = k_{j-1}t_j + k_j, \quad j = 2, \dots, n, \tag{3.43}$$

with the sequence  $k_j \in \mathbb{Z}, j = 1, \dots, n$ , chosen so that  $\frac{1}{2}t_j \leq \gamma_j \leq 1 + \frac{1}{2}t_j$ .

Furthermore, for this  $p/q$ , there holds the following bound for the sum (3.20):

$$|S_k| = \left| \sum_{m=1}^{q-1} \frac{1}{m} \frac{\cos \pi \frac{p}{q} m(4k+1)}{\cos \pi \frac{p}{q} m} \right| < \left( 4 + \frac{\beta}{\ln 2} \right) \ln q + 9. \tag{3.44}$$

If, in addition,  $q_{k+1} \geq q_k^\nu$ , for some  $\nu > 1$  and all  $1 \leq k \leq n - 1$ , then for any  $\varepsilon > 0$  there exist  $Q = Q(\varepsilon, \nu)$  such that if  $q > Q$ ,

$$|S_k| < (4 + \varepsilon) \ln q. \tag{3.45}$$

*Proof.* First, note that  $t_1 = p/q$  and  $t_j = 1/(a_j + t_{j+1})$ ,  $j = 1, \dots, n - 1$ ,  $t_n = 1/a_n$ . Now note a simple fact that the conditions  $a_1 - \text{odd}$ ,  $a_2, \dots, a_n - \text{even}$  ensure that  $q$  is odd and all the fractions  $t_j$  for  $j = 2, \dots, n$  are either  $\frac{\text{odd}}{\text{even}}$  or  $\frac{\text{even}}{\text{odd}}$ . We now choose  $k_1$  so that  $\gamma_1 = -2kt_1 + k_1$  satisfies  $\frac{1}{2}t_1 \leq \gamma_1 \leq 1 + \frac{1}{2}t_1$ . Applying Lemma 5 to  $S(t_1, \gamma_1)$ , we obtain

$$S(t_1, \gamma_1) = I(t_1, \gamma_1) + (-1)^\varepsilon T(t_2, \gamma_2, a_1 \bmod 2) \tag{3.46}$$

We can and will choose  $\gamma_2$  so that  $\frac{1}{2}t_2 \leq \gamma_2 \leq 1 + \frac{1}{2}t_2$  by picking the appropriate  $k'_2$  and setting

$$\begin{aligned} \gamma_2 &= \frac{q}{p}\gamma_1 + k'_2 \\ &= \frac{q}{p}(-2kt_1 + k_1) + k'_2 \\ &= -2k + \frac{q}{p}k_1 + k'_2 = -2k + (a_1 + t_2)k_1 + k'_2 = t_2k_1 + k_2, \end{aligned} \tag{3.47}$$

where  $k_2 = -2k + a_1k_1 + k'_2$ . The constant  $\varepsilon$  is determined by  $\gamma_2$  as described in Lemma 5.

Since  $a_1 \bmod 2 = 1$ , we can now apply Lemma 6 with  $\delta = 1$  to  $T$  in the r.h.s. of (3.46). This gives

$$S(t_1, \gamma_1) = I(t_1, \gamma_1) + (-1)^\varepsilon [J(t_2, \gamma_2, 1) - (-1)^{\varepsilon'} T(t_3, \gamma_3, a_2 + 1 \bmod 2)]. \tag{3.48}$$

with

$$\gamma_3 = \frac{1}{t_2}\gamma_2 + k'_3 = k_1 + (a_3 + t_3)k_2 + k'_3 = t_3k_2 + k_3.$$

chosen so that  $\frac{1}{2}t_3 \leq \gamma_3 \leq 1 + \frac{1}{2}t_3$ .

Note that according to our assumption  $a_j + 1 \bmod 2 = 1$ ,  $j = 2, \dots, n$ , so that we can continue applying Lemma 6 with  $\delta = 1$  in (3.48) recursively. At the final step, revisiting the residue calculations for Lemma 6 gives:

$$T(t_n, \gamma_n, 1) = J(t_n, \gamma_n, 1) + 2e^{-i\pi\gamma_n a_n}, \tag{3.49}$$

which proves (3.42).

Now using the bounds (3.27) and (3.38) for  $\delta = 1$ , we write

$$\begin{aligned} |S_k| &< |S(p/q, \gamma_1)| \leq |I(t_1, \gamma_1)| + \sum_{j=2}^n |J(t_j, \gamma_j, 1)| + 2 \\ &< 4 \ln \frac{1}{t_1 t_2 \cdots t_n} + \frac{5}{e\pi} \sum_{j=1}^n \frac{1}{p'_j} + n\beta + 2, \end{aligned} \tag{3.50}$$

where  $t_j = [a_j, \dots, a_n] = \frac{p'_j}{q'_j}$ .

Recall that we denote  $q = q_n$ ,  $p = p_n$ . Observe that the following recurrence [10] with  $t_{n+1} = 0$

$$q_n + t_{n+1}q_{n-1} = (a_n + t_{n+1})q_{n-1} + q_{n-2} = \frac{1}{t_n}(q_{n-1} + t_n q_{n-2})$$



gives

$$\frac{1}{t_1 t_2 \cdots t_n} = q_n. \tag{3.51}$$

Note that this equation holds for any continued fraction.

Furthermore, the recurrence  $t_{j-1} = 1/(a_{j-1} + t_j)$  implies

$$p'_{j-1} = q'_j, \quad q'_{j-1} = a_{j-1} p'_{j-1} + p'_j,$$

and so

$$p'_{j-2} = a_{j-1} p'_{j-1} + p'_j,$$

which, with the initial conditions  $p'_n = 1$ ,  $p'_{n-1} = a_n$ , and our assumption that all  $a_j$ ,  $j = 2, \dots, n$ , are even, gives

$$p'_j \geq a_{j+1} \cdots a_n \geq 2^{n-j}, \quad j = 1, \dots, n. \tag{3.52}$$

Thus we have

$$\sum_{j=1}^n \frac{1}{p'_j} < 2. \tag{3.53}$$

Finally, since

$$q_n = a_n q_{n-1} + q_{n-2} \geq a_n q_{n-1} \geq \cdots \geq a_n a_{n-1} \cdots a_1 \geq 2^{n-1},$$

we have

$$n \leq \frac{\ln q_n}{\ln 2} + 1. \tag{3.54}$$

(In fact, as is well known, a slightly worse bound on  $n$  holds for any continued fraction).

Using (3.51), (3.53), (3.54) in (3.50), we obtain (3.44).

To obtain (3.45) note first the following.

**Lemma 8.** *Let  $\nu > 1$ ,  $n \geq 2$ ,  $p_k/q_k = [a_1, a_2, \dots, a_k]$ , and such that  $q_{k+1} \geq q_k^\nu$ ,  $1 \leq k \leq n - 1$ ,  $q_2 \geq 3$ . Then*

$$n \leq \frac{1}{\ln \nu} \ln \frac{\ln q_n}{\ln 3} + 2$$

*Proof.* We have

$$q_n \geq q_{n-1}^\nu \geq q_{n-2}^{\nu^2} \geq \cdots \geq q_2^{\nu^{n-2}} \geq 3^{\nu^{n-2}}, \quad n \geq 2,$$

from which the result follows.  $\square$

Now using Lemma 8 instead of (3.54) in (3.50) and choosing  $Q$  sufficiently large, we obtain (3.45) if  $q > Q$ , and finish the proof of Lemma 7.  $\square$

Bringing together (3.6), (3.18), (3.19), and (3.44), yields the bound

$$|\sigma'(0)| < e^{-\frac{2}{3}\gamma_0} \frac{2}{3} q \cdot q^{\gamma_0+4+\beta/\ln 2} e^{2\gamma_0+9} < q^{\gamma_0+5+\beta/\ln 2} \frac{2}{3} e^{9+4\gamma_0/3}, \quad q \geq 3. \tag{3.55}$$

Since  $\gamma_0 + 5 + \beta/\ln 2 = 13.8 \dots < 14$ ,  $\frac{2}{3} e^{9+4\gamma_0/3} < e^{10}$ , and  $\sigma(E) = -E$  if  $q = 1$ , we obtain (1.7). Finally, using (3.45) instead of (3.44), we obtain (1.9). This finishes the proof of Lemma 2.  $\square$

### 4. Proof of Theorem 4

Note first that since for any irrational  $\alpha$

$$\frac{1}{2q_n q_{n+1}} < \frac{1}{q_n(q_n + q_{n+1})} < \left| \alpha - \frac{p_n}{q_n} \right|,$$

we obtain that for any  $\alpha$  satisfying the conditions of Theorem 4,

$$\frac{C_3}{2} q_n^{x-1} < q_{n+1}, \quad n = 1, 2, \dots \tag{4.1}$$

We will denote  $\mu_j(p/q)$ ,  $w_j(p/q)$ , etc, the values  $\mu_j$ ,  $w_j$ , etc, for the spectrum  $S(p/q)$ . Fix  $n \geq 1$ . Let  $E_2, E_0$  be the right edges of the centermost bands in  $S(p_n/q_n)$ ,  $S(p_{n+1}/q_{n+1})$ , respectively. By Lemma 2,

$$E_2 = \mu_0(p_n/q_n) = w_0(p_n/q_n) > \frac{4}{C_2 q_n^{C_1}}. \tag{4.2}$$

On the other hand by (2.16) and (4.1), we have

$$E_0 = \mu_0(p_{n+1}/q_{n+1}) = w_0(p_{n+1}/q_{n+1}) < \frac{4e}{q_{n+1}} < \frac{8e}{C_3 q_n^{x-1}}. \tag{4.3}$$

We will now show that  $G_0(p_{n+1}/q_{n+1}) \subset (E_0, E_2 - \epsilon)$  for a suitably chosen  $C_3$ .

Recall a continuity property found by Avron, Van Mouche, Simon [5]: if  $E \in S(\beta)$ , there is  $E' \in S(\beta')$  such that

$$|E - E'| < C|\beta - \beta'|^{1/2}. \tag{4.4}$$

In [5], the authors give a good bound on  $C$  requiring that  $|\beta - \beta'|$  be sufficiently small. As the reader can verify, a trivial modification of the proof in [5] allows us to fix  $C = 60$  for the almost Mathieu operator (worse than in [5]) but without any condition on  $\beta, \beta' \in (0, 1)$ . Thus we set  $C = 60$ .

This continuity property (4.4) for  $\beta = p_n/q_n, \beta' = p_{n+1}/q_{n+1}$ , together with the identity

$$\left| \frac{p_n}{q_n} - \frac{p_{n+1}}{q_{n+1}} \right| = \frac{1}{q_n q_{n+1}}$$

and the bound (4.1) implies that there exists  $E' \in S(p_{n+1}/q_{n+1})$  such that

$$E' \in \left( \frac{E_2}{2} - \frac{C}{\sqrt{q_n q_{n+1}}}, \frac{E_2}{2} + \frac{C}{\sqrt{q_n q_{n+1}}} \right) \subset \left( \frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^x / 2}}, \frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^x / 2}} \right). \tag{4.5}$$

Using (4.2) and recalling that<sup>1</sup>  $x = 4C_1$ , we see that

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^x / 2}} > \frac{2}{C_2 q_n^{C_1}} - \frac{C}{\sqrt{C_3 / 2} q_n^{2C_1}} = \frac{2}{C_2 q_n^{C_1}} \left( 1 - \frac{CC_2}{\sqrt{2C_3} q_n^{C_1}} \right),$$

<sup>1</sup> Note that here just  $x = 2C_1$  would do, cf a remark following Theorem 4.

and setting now

$$C_3 = 4^2 C^2 C_2^4 = 4^2 60^2 C_2^4, \tag{4.6}$$

we have

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^x / 2}} > \frac{q_n^{-C_1}}{C_2}. \tag{4.7}$$

On the other hand, using (4.3), we obtain that

$$E_0 < \frac{8e}{C_3 q_n^{4C_1-1}} = \frac{e}{2C^2 C_2^4 q_n^{4C_1-1}} < \frac{q_n^{-C_1}}{C_2}.$$

Inequality (4.7) also shows that

$$\frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^x / 2}} < E_2. \tag{4.8}$$

Thus,

$$E' \in (E_0, E_2),$$

which implies that

$$G_0(p_{n+1}/q_{n+1}) \subset (E_0, E_2 - \epsilon),$$

for some  $\epsilon > 0$ . The corresponding result for  $G_{-1}$  follows by the symmetry of the spectra. This proves the statement (a) of Theorem 4.

Now by the continuity (4.4) with  $\beta = \alpha$ ,  $\beta' = p_n/q_n$ , and Theorem 3, we conclude that, for all  $n = 1, 2, \dots$ , there exists a gap  $G_{n,2}(\alpha)$  of  $S(\alpha)$  such that  $G_{n,2}(\alpha) \cap G_0(p_n/q_n) \neq \emptyset$  and of length

$$\Delta_{n,2}(\alpha) > \Delta_0(p_n/q_n) - 2C|\alpha - p_n/q_n|^{1/2} > \frac{1}{C_2^2 q_n^{2C_1}} - \frac{2C}{C_3^{1/2} q_n^{x/2}} = \frac{1}{2C_2^2 q_n^{x/2}}. \tag{4.9}$$

We now verify that the gaps  $G_{n,2}(\alpha)$ ,  $G_{n+1,2}(\alpha)$  are distinct. Using the continuity once again, we obtain that there exists a point  $E'' \in S(\alpha)$  such that

$$E'' \in \left( \frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^x}}, \frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^x}} \right). \tag{4.10}$$

Now it is easy to verify, similar to the calculations above, that

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^x}} > \frac{7}{4} \frac{1}{C_2 q_n^{C_1}}$$

and

$$E_0 + \frac{2C}{\sqrt{C_3 q_n^x}} < \frac{1}{C_2 q_n^{C_1}},$$

and therefore (4.10) together with (4.8) yields

$$E'' \in \left( E_0 + \frac{2C}{\sqrt{C_3 q_{n+1}^\alpha}}, E_2 \right). \quad (4.11)$$

Thus  $G_{n+1,2}(\alpha)$  lies to the left of  $E''$ , and  $G_{n,2}(\alpha)$  to the right of  $E''$ , so that  $G_{n,2}(\alpha)$  and  $G_{n+1,2}(\alpha)$  are distinct gaps,  $n = 1, 2, \dots$ . Similar results for  $G_{n,1}(\alpha)$  follow by the symmetry. This proves the statement (b) of Theorem 4.

The proof of the statement (c) is similar and based on (1.9). It is a simple exercise.  $\square$

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